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# 4,6-Dimethyl-α-pyrone: a matrix isolation study of the photochemical generation of conjugated ketene, Dewar valence isomer and 1,3-dimethyl-cyclobutadiene

S. Breda<sup>a,b</sup>, L. Lapinski<sup>a,b</sup>, I. Reva<sup>a</sup>, R. Fausto<sup>a,\*</sup>

<sup>a</sup> Department of Chemistry, University of Coimbra, Coimbra P-3004-35, Portugal <sup>b</sup> Institute of Physics, Polish Academy of Sciences, Warsaw PL-02-668, Poland

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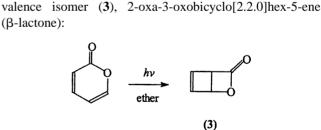
## Abstract

A combined matrix isolation and molecular orbital study of the vibrational spectra and photochemistry of 4,6-dimethyl- $\alpha$ -pyrone (DMAP) was undertaken. Two types of photoreactions: ring opening leading to conjugated ketene and valence isomerization to the Dewar form (1,5-dimethyl-2-oxa-3-oxobicyclo[2.2.0]hex-5-ene; DOOBH), occurred upon UV ( $\lambda > 315$  nm) irradiation. The latter reaction was efficient, whereas aldehyde–ketene was produced only in little amounts. In addition to the IR spectroscopic study of DMAP, the full mid-IR spectrum of the photoproduced DOOBH is reported and interpreted. Observation of 1,3-dimethyl-cyclobutadiene (DMCB), created by shorter wavelength UV irradiation ( $\lambda > 235$  nm) of DOOBH, is reported for the first time. In the matrices, DMCB forms a complex with CO<sub>2</sub>; the structure and IR absorption features of this cage confined DMCB–CO<sub>2</sub> complex are also investigated. © 2004 Elsevier B.V. All rights reserved.

Keywords: 4,6-Dimethyl-α-pyrone; Matrix isolation; DMCB-CO<sub>2</sub> complex; Photoreaction; Dewar; Ketene

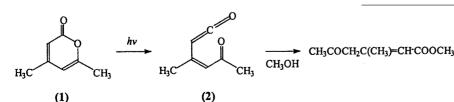
# 1. Introduction

Important, naturally occurring compounds and synthetic pesticides, such as coumarins, psoralens and chromones, that are known to be potent photosensitizers, possess pyrone moieties as their fundamental structural nucleus [1,2]. Pyrones are also very useful reagents in organic synthesis, e.g. in the Diels Alder reaction, and have been shown to exhibit interesting photochemical properties [3]. First reports on the photochemistry of  $\alpha$ -pyrones come from the early work of de Mayo [4] who observed an open-ring ester photoproduct generated from 4,6-dimethyl- $\alpha$ -pyrone (1) dissolved in methanol. A conjugated aldehyde–ketene (2) was presumed to be the primary product of the photoreaction, although this intermediate species was not directly observed:



On the other hand, irradiation of a solution of  $\alpha$ -pyrone in ether at -10 to -20 °C, carried out by Corey and Streith [5], resulted in isomerization of the compound to its Dewar

In the study of UV irradiated thin solid films layers of 4,6-dimethyl- $\alpha$ -pyrone at -190 °C, Guthrie et al. [6] were able to observe IR bands that could be attributed to



\* Corresponding author.

E-mail address: rfausto@ci.uc.pt (R. Fausto).

both photoproducts: Dewar valence isomer and conjugated aldehyde-ketene. This result indicated that both species could indeed be photoproduced from the reactant molecule.

In the present work, we report a combined matrix isolation and molecular orbital study of the vibrational spectra and photochemistry of 4,6-dimethyl- $\alpha$ -pyrone (DMAP). Both types of photoreactions, ring opening leading to conjugated ketene and valence isomerization to the Dewar form, were observed, the latter reaction being much more efficient than the first one. In addition to the photochemical and spectroscopic study of the IR spectrum of DMAP, the full mid-IR spectrum of its Dewar valence isomer is reported and interpreted. Observation of the antiaromatic 1,3-dimethyl-cyclobutadiene, created by shorter wavelength UV irradiation ( $\lambda > 235$  nm) of the DMAP Dewar isomer, is reported for the first time, and the IR absorption features of its complex with CO<sub>2</sub> are analyzed.

## 2. Experimental

DMAP was obtained from Aldrich. A sample of the compound was placed in a glass tube protected against light and connected to the chamber of the cryostat with a needle valve. Before cooling down the cryostat, the compound was degassed by the standard freeze-pump-thaw procedure and, subsequently, the vapors over the compound in the tube were evacuated several times at room temperature. This approach enabled removal of possible volatile impurities, allowing an additional purification of the compound, immediately before each experiment. In order to deposit a matrix, the vapor of DMAP was introduced into the cryostat chamber together with large excess of the host matrix gas (argon N60, from Air Liquide). The gaseous mixture was co-deposited onto a cold CsI window (T = 10 K) mounted on the tip of an APD Cryogenics DE-202A closed-cycle helium refrigerator. Care was taken to keep the guest-to-host ratio in matrices low enough to avoid association.

The matrices were irradiated through the outer KBr window of the cryostat, with filtered or unfiltered light from a 150 W xenon arc lamp (Osram XBO 150 W/CR OFR).

The infrared spectra were recorded with  $0.5 \text{ cm}^{-1}$  resolution using a Mattson (AR60) Infinity Series FTIR spectrometer equipped with a KBr beamsplitter and a DTGS detector.

### **3.** Computational details

The equilibrium geometries for all studied species were fully optimized at the DFT level of theory with the standard 6-311++G(d,p) basis set. The DFT calculations were carried out with the three-parameter density functional (B3LYP) which includes Becke's gradient exchange correction [7], the Lee et al. correlation functional [8] and the Vosko et al. [9] correlation functional. During optimizations no restriction of symmetry was imposed on the initial structure. The

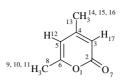


Fig. 1. Atom numbering used in the normal mode analysis of 4,6-dimethyl- $\alpha$ -pyrone.

optimized geometry for 4,6-dimethyl- $\alpha$ -pyrone converged to a structure with the planar heavy-atom backbone (Fig. 1).

The geometry optimizations were followed by IR frequency calculations to verify whether the calculated structures are local minima on the potential energy surfaces of the respective systems. For each of the studied species, a set of internal coordinates was defined and the Cartesian force constants were transformed to the internal coordinates space, allowing ordinary normal-coordinate analysis to be performed as described by Schachtschneider [10]. Internal coordinate sets defined for DMAP, its Dewar isomer and the complex between 1,3-dimethyl-cyclobutadiene and  $CO_2$  are

Table 1

Internal coordinates used in the normal mode analysis for DMAP (atom numbering as in Fig.  $1)^{a}\,$ 

$S_1 = r_{1,2}$	v(O1–C2)
$S_2 = r_{2,3}$	v(C2–C3)
$S_3 = r_{3,4}$	$\nu$ (C3=C4)
$S_4 = r_{4,5}$	$\nu(C4-C5)$
$S_5 = r_{5,6}$	v(C5=C6)
$S_6 = r_{6,1}$	v(C6–O1)
$S_7 = r_{7,2}$	v(C2=O7)
$S_8 = r_{8,6}$	$\nu(C6-C8)$
$S_9 = (3^{-1/2})(r_{9,8} + r_{10,8} + r_{11,8})$	$\nu(CH_3)^1_s$
$S_{10} = (6^{-1/2})(2r_{11,8} - r_{10,8} - r_{9,8})$	$\nu(CH_3)_{as}^{1'}$
$S_{11} = (2^{-1/2})(r_{10,8} - r_{9,8})$	$\nu(CH_3)^{1''}_{as}$
$S_{12} = r_{12,5}$	v(C5–H12)
$S_{13} = r_{13,4}$	$\nu$ (C4–C13)
$S_{14} = (3^{-1/2})(r_{14,13} + r_{15,13} + r_{16,13})$	$\nu(CH_3)_s^2$
$S_{15} = (6^{-1/2})(2r_{16,13} - r_{15,13} - r_{14,13})$	$\nu(CH_3)^{2'}_{as}$
$S_{16} = (2^{-1/2})(r_{15,13} - r_{14,13})$	$\nu(CH_3)^{2''}_{as}$
$S_{17} = r_{17,3}$	ν(C3–H17)
$S_{18} = (6^{-1/2})(\beta_{6,2,1} - \beta_{1,3,2} + \beta_{2,4,3} - \beta_{3,5,4})$	$\delta$ ring 1
$+\beta_{4,6,5} - \beta_{5,1,6}$	
$S_{19} = (2^{-1/2})(\beta_{7,1,2} - \beta_{7,3,2})$	$\delta$ (C2=O7)
$S_{20} = (2^{-1/2})(\beta_{8,1,6} - \beta_{8,5,6})$	$\delta(C6-C8)$
$S_{21} = (6^{-1/2})(\beta_{9,10,8} + \beta_{10,11,8} + \beta_{11,9,8} - \beta_{9,6,8})$	$\delta(CH_3)^1_s$
$-\beta_{10,6,8} - \beta_{11,6,8})$	×GU )1/
$S_{22} = (6^{-1/2})(2\beta_{9,10,8} - \beta_{10,11,8} - \beta_{11,9,8})$	$\delta(CH_3)^{1'}_{as}$
$S_{23} = (2^{-1/2})(\beta_{10,11,8} - \beta_{11,9,8})$	$\delta(CH_3)^{1''}_{as}$
$S_{24} = (6^{-1/2})(2\beta_{11,6,8} - \beta_{10,6,8} - \beta_{9,6,8})$	$\gamma(CH_3)^{1'}_{as}$
$S_{25} = (2^{-1/2})(\beta_{10,6,8} - \beta_{9,6,8})$	$\gamma$ (CH <sub>3</sub> ) <sup>1"</sup> <sub>as</sub>
$S_{26} = (2^{-1/2})(\beta_{12,6,5} - \beta_{12,4,5})$	δ(C5–H12)
$S_{27} = (2^{-1/2})(\beta_{13,3,4} - \beta_{13,5,4})$	$\delta$ (C4–C13)
$S_{28} = (6^{-1/2})(\beta_{14,15,13} + \beta_{15,16,13} + \beta_{16,14,13}$	$\delta(CH_3)_s^2$
$-\beta_{14,4,13} - \beta_{15,4,13} - \beta_{16,4,13}$	
$S_{29} = (6^{-1/2})(2\beta_{14,15,13} - \beta_{15,16,13} - \beta_{16,14,13})$	$\delta(CH_3)^{2'}_{as}$

Table 1 (Continued)

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$S_{30} = (2^{-1/2})(\beta_{15,16,13} - \beta_{16,14,13})$	$\delta(CH_3)^{2''}_{as}$
$S_{31} = (2^{-1/2})(\beta_{17,4,3} - \beta_{17,2,3})$	δ(C3–H17)
$S_{32} = (6^{-1/2})(2\beta_{16,4,13} - \beta_{15,4,13} - \beta_{14,4,13})$	$\gamma$ (CH <sub>3</sub> ) <sup>2'</sup> <sub>as</sub>
$S_{33} = (2^{-1/2})(\beta_{15,4,13} - \beta_{14,4,13})$	$\gamma$ (CH <sub>3</sub> ) <sup>2''</sup> <sub>as</sub>
$S_{34} = (12^{-1/2})(2\beta_{6,2,1} - \beta_{1,3,2} - \beta_{2,4,3})$	δ ring 2
$+2\beta_{3,5,4} - \beta_{4,6,5} - \beta_{5,1,6}$	
$S_{35} = (1/2)(\beta_{1,3,2} - \beta_{2,4,3} + \beta_{4,6,5} - \beta_{5,1,6})$	$\delta$ ring 3
$S_{36} = (6^{-1/2})(\tau_{6,1,2,3} - \tau_{1,2,3,4} + \tau_{2,3,4,5})$	$\gamma$ ring 1
$-\tau_{3,4,5,6} + \tau_{4,5,6,1} - \tau_{5,6,1,2}$	
$S_{37} = (12^{-1/2})(2\tau_{6,1,2,3} - \tau_{1,2,3,4} - \tau_{2,3,4,5})$	$\gamma$ ring 2
$+ 2\tau_{3,4,5,6} - \tau_{4,5,6,1} - \tau_{5,6,1,2})$	
$S_{38} = (1/2)(\tau_{1,2,3,4} - \tau_{2,3,4,5} + \tau_{4,5,6,1} - \tau_{5,6,1,2})$	$\gamma$ ring 3
$S_{39} = (6^{-1/2})(\tau_{9,8,6,1} + \tau_{9,8,6,5} + \tau_{10,8,6,1}$	Twist (CH <sub>3</sub> ) <sup>1</sup>
$+ \tau_{10,8,6,5} + \tau_{11,8,6,1} + \tau_{11,8,6,5})$	
$S_{40} = (6^{-1/2})(\tau_{14,13,4,3} + \tau_{14,13,4,5} + \tau_{15,13,4,3})$	Twist (CH <sub>3</sub> ) <sup>2</sup>
$+ \tau_{15,13,4,5} + \tau_{16,13,4,3} + \tau_{16,13,4,5})$	
$S_{41} = \gamma_{7,3,2,1}$	$\gamma$ (C2=O7)
$S_{42} = \gamma_{8,1,6,5}$	γ(C6–C8)
$S_{43} = \gamma_{12,6,5,4}$	γ(C5–H12)
$S_{44} = \gamma_{13,5,4,3}$	γ(C4–C13)
$S_{45} = \gamma_{17,4,3,2}$	γ(C3–H17)

<sup>a</sup>  $r_{i,j}$  is the distance between atoms  $A_i$  and  $A_j$ ;  $\beta_{i,j,k}$  is the angle between vectors  $A_kA_i$  and  $A_kA_j$ ;  $\tau_{i,j,k,l}$  is the dihedral angle between the plane defined by  $A_i$ ,  $A_j$ ,  $A_k$  and the plane defined by  $A_j$ ,  $A_k$ ,  $A_l$  atoms;  $\gamma_{i,j,k,l}$  is the angle between the vector  $A_kA_i$  and the plane defined by atoms  $A_j$ ,  $A_k$ ,  $A_l$ .

#### Table 2

Internal coordinates used in the normal mode analysis for DOOBH (atom numbering as in Fig. 3)<sup>a</sup>

$S_1 = r_{1,2}$	v(O1–C2)
$S_2 = r_{2,3}$	v(C2–C3)
$S_3 = r_{3,4}$	v(C3–C4)
$S_4 = r_{4,5}$	v(C4=C5)
$S_5 = r_{5,6}$	v(C5–C6)
$S_6 = r_{6,1}$	v(C6–O1)
$S_7 = r_{3,6}$	v(C3–C6)
$S_8 = r_{7,2}$	v(C2=O7)
$S_9 = r_{8,6}$	v(C6–C8)
$S_{10} = (3^{-1/2})(r_{9,8} + r_{10,8} + r_{11,8})$	$\nu(CH_3)^1_s$
$S_{11} = (6^{-1/2})(2r_{10,8} - r_{9,8} - r_{11,8})$	$\nu(CH_3)^{1'}_{as}$
$S_{12} = (2^{-1/2})(r_{9,8} - r_{11,8})$	$\nu(CH_3)^{1''}_{as}$
$S_{13} = r_{12,5}$	v(C5–H12)
$S_{14} = r_{13,4}$	v(C4–C13)
$S_{15} = (3^{-1/2})(r_{14,13} + r_{15,13} + r_{16,13})$	$\nu(CH_3)_s^2$
$S_{16} = (6^{-1/2})(2r_{14,13} - r_{15,13} - r_{16,13})$	$\nu(CH_3)^{2'}_{as}$
$S_{17} = (2^{-1/2})(r_{15,13} - r_{16,13})$	$\nu(\mathrm{CH}_3)^{2''}_{\mathrm{as}}$
$S_{18} = r_{17,3}$	v(C3–H17)
$S_{19} = (1/2)(\beta_{2,6,1} + \beta_{2,6,3} - \beta_{1,3,2} - \beta_{1,3,6})$	$\delta$ ring 1
$S_{20} = (1/2)(\beta_{6,4,3} + \beta_{6,4,5} - \beta_{3,5,6} - \beta_{3,5,4})$	$\delta$ ring 2
$S_{21} = (2^{-1/2})(\beta_{7,1,2} - \beta_{7,3,2})$	$\delta(C2=O7)$
$S_{22} = \beta_{8,3,6}$	$\delta(C6-C8)$
$S_{23} = (6^{-1/2})(\beta_{9,10,8} + \beta_{10,11,8} + \beta_{11,9,8})$	$\delta(CH_3)^1_s$
$-\beta_{9,6,8} - \beta_{10,6,8} - \beta_{11,6,8}$	
$S_{24} = (6^{-1/2})(2\beta_{11,9,8} - \beta_{10,11,8} - \beta_{9,10,8})$	$\delta(CH_3)^{1'}_{as}$
$S_{25} = (2^{-1/2})(\beta_{10,11,8} - \beta_{9,10,8})$	$\delta(CH_3)^{1''}_{as}$
$S_{26} = (6^{-1/2})(2\beta_{10,6,8} - \beta_{9,6,8} - \beta_{11,6,8})$	$\gamma$ (CH <sub>3</sub> ) <sup>1'</sup> <sub>as</sub>
$S_{27} = (2^{-1/2})(\beta_{9,6,8} - \beta_{11,6,8})$	$\gamma$ (CH <sub>3</sub> ) <sup>1"</sup> <sub>as</sub>
$S_{28} = (2^{-1/2})(\beta_{12,6,5} - \beta_{12,4,5})$	δ(C5–H12)
$S_{29} = (2^{-1/2})(\beta_{13,3,4} - \beta_{13,5,4})$	δ(C4–C13)
$S_{30} = (6^{-1/2})(\beta_{14,15,13} + \beta_{15,16,13} + \beta_{16,14,13}$	$\delta(CH_3)^2_s$
$-\beta_{14,4,13} - \beta_{15,4,13} - \beta_{16,4,13})$	- 5
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Table 2 (Continued)

$S_{31} = (6^{-1/2})(2\beta_{15,16,13} - \beta_{14,15,13} - \beta_{16,14,13})$	$\delta(CH_3)^{2'}_{as}$
$S_{32} = (2^{-1/2})(\beta_{14,15,13} - \beta_{16,14,13})$	$\delta(CH_3)^{2''}_{as}$
$S_{33} = (6^{-1/2})(2\beta_{14,4,13} - \beta_{15,4,13} - \beta_{16,4,13})$	$\gamma$ (CH <sub>3</sub> ) <sup>2'</sup> <sub>as</sub>
$S_{34} = (2^{-1/2})(\beta_{15,4,13} - \beta_{16,4,13})$	$\gamma$ (CH <sub>3</sub> ) <sup>2''</sup> <sub>as</sub>
$S_{35} = \beta_{17,6,3}$	δ(C3–H17)
$S_{36} = (1/2)(\tau_{2,1,6,3} + \tau_{6,3,2,1} - \tau_{1,6,3,2} - \tau_{3,2,1,6})$	$\tau$ ring 1
$S_{37} = (1/2)(\tau_{6,3,4,5} + \tau_{4,5,6,3} - \tau_{5,6,3,4} - \tau_{3,4,5,6})$	$\tau$ ring 2
$S_{38} = (2^{-1/2})(\tau_{8,6,3,4} + \tau_{8,6,3,2})$	γ(C6–C8)
$S_{39} = (2^{-1/2})(\tau_{17,3,6,1} + \tau_{17,3,6,5})$	γ(C3–H17)
$S_{40} = (2^{-1/2})(\tau_{2,3,6,5} - \tau_{4,3,6,1})$	$\tau$ butterfly
$S_{41} = (6^{-1/2})(\tau_{9,8,6,1} + \tau_{9,8,6,5} + \tau_{10,8,6,1}$	Twist (CH <sub>3</sub> ) <sup>1</sup>
$+ \tau_{10,8,6,5} + \tau_{11,8,6,1} + \tau_{11,8,6,5})$	
$S_{42} = (6^{-1/2})(\tau_{14,13,4,3} + \tau_{14,13,4,5} + \tau_{15,13,4,3})$	Twist (CH <sub>3</sub> ) <sup>2</sup>
$+ \tau_{15,13,4,5} + \tau_{16,13,4,3} + \tau_{16,13,4,5})$	
$S_{43} = \gamma_{7,1,2,3}$	$\gamma$ (C2=O7)
$S_{44} = \gamma_{12,6,5,4}$	γ(C5–H12)
$S_{45} = \gamma_{13,5,4,3}$	γ(C4–C13)

<sup>a</sup>  $r_{i,j}$  is the distance between atoms  $A_i$  and  $A_j$ ;  $\beta_{i,j,k}$  is the angle between vectors  $A_kA_i$  and  $A_kA_j$ ;  $\tau_{i,j,k,l}$  is the dihedral angle between the plane defined by  $A_i$ ,  $A_j$ ,  $A_k$  and the plane defined by  $A_j$ ,  $A_k$ ,  $A_l$  atoms;  $\gamma_{i,j,k,l}$  is the angle between the vector  $A_kA_i$  and the plane defined by atoms  $A_j$ ,  $A_k$ ,  $A_l$ .

given in Tables 1–3. The calculated harmonic frequencies were also used to assist the analysis of the experimental spectra and to confirm the nature of the stationary points resulting from the calculations (for all structures discussed below no imaginary frequencies were obtained, indicating they correspond to true minima).

All calculations in this work were done using the Gaussian'98 program [11].

## 4. Results and discussion

The IR spectrum of DMAP isolated in an Ar matrix is presented in Fig. 2. This spectrum is compared with the results of the theoretical simulations carried out at the DFT(B3LYP)/6-311++G(d,p) level. The frequencies and intensities of the observed bands and the corresponding calculated values are collected in Table 4.

During UV ( $\lambda > 315$  nm) irradiation of the matrix, the initial IR spectrum was systematically decreasing, whereas the spectrum of the Dewar valence isomer of DMAP, 1,5-dimethyl-2-oxa-3-oxobicyclo[2.2.0]hex-5-ene (DOOBH; Fig. 3), appeared and was continuously growing (Fig. 4). The main spectral indication of the photoproduction of DOOBH was the strong band at 1837.7 cm<sup>-1</sup> (Ar), which corresponds to the  $\nu$ C=O vibration of this species. The high frequency of this  $\nu$ C=O band is typical for carbonyl groups directly attached to a small, four-membered ring. The general rule, saying that the smaller the ring the higher is the  $\nu$ C=O frequency, is fulfilled in the case of  $\alpha$ -pyrones. Analogously, the  $\nu$ C=O band of the Dewar isomer of 3-methyl-4(3H)-pyrimidinone was observed at 1784.0 cm<sup>-1</sup> (Ar), whereas the  $\nu$ C=O band of the normal

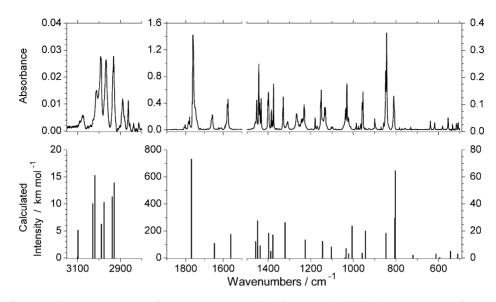


Fig. 2. Comparison of the experimental IR spectrum of DMAP monomers isolated in Ar matrix (10 K) with the spectrum of the compound calculated at the DFT(B3LYP)/6-311++G(d,p) level. Theoretical frequencies were scaled using a factor equal to 0.978 (except for the CH stretching region, where the factor 0.968 was used). Note the change of ordinate scale at  $1500 \text{ cm}^{-1}$ .

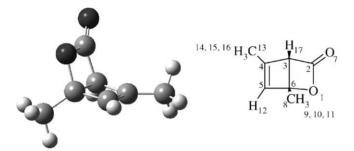


Fig. 3. Optimized structure and atom numbering of the Dewar valence isomer (DOOBH).

isomer of 3-methyl-4(3H)-pyrimidinone was observed at significantly lower frequency  $1709.0 \text{ cm}^{-1}$  (Ar) [12].

The progress of the UV ( $\lambda > 315$  nm) induced photoreaction is shown in Fig. 4. Some amount of the open-ring ketene derivative [13] of DMAP was also created, as it can be concluded from the appearance of the characteristic ketene IR band at 2128.9 cm<sup>-1</sup>. At the final stage of the observed photoreaction (after 523 min of irradiation) about 90% of the initial material was converted into DOOBH, while the total amount of the open-ring ketene could be estimated to be ca. 5% of the main photoproduct only.<sup>1</sup>

Having the IR spectra recorded at different stages of the photoreaction it was then possible to unequivocally identify the bands due to DOOBH. The experimental and calculated spectra for this main photoproduct are compared in Fig. 5, showing a good general agreement. The proposed band assignments are given in Table 5. The agreement between the

experimental and theoretical data is particularly good in the frequency range below  $1000 \text{ cm}^{-1}$ , where the most intense experimental bands at 827.4, 783.6 and 509.3 cm<sup>-1</sup> are very well reproduced by the theoretical calculations, which

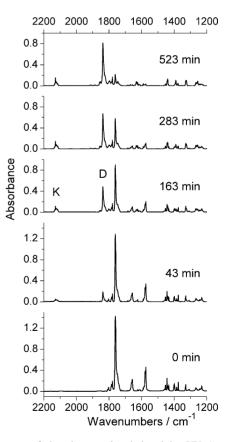


Fig. 4. Progress of the photoreaction induced by UV ( $\lambda > 315$  nm) irradiation of DMAP isolated in Ar matrix (10 K). K: open-ring conjugated aldehyde–ketene, D: Dewar valence isomer (DOOBH).

<sup>&</sup>lt;sup>1</sup> The relative amounts of DOOBH and open-ring ketene photoproducts were estimated from the ratio of the observed integral intensities of the features around 1850 and 2130 cm<sup>-1</sup>, respectively, reduced by the corresponding calculated intensities for these species.

Internal coordinates used in the normal mode analysis for the DMCB stacked to  $CO_2$  (atom numbering as in Fig. 6)<sup>a</sup>

$S_1 = (2^{-1/2})(r_{1,2} + r_{3,4})$	$\nu$ (C=C) <sub>s</sub>
$S_2 = (2^{-1/2})(r_{1,2} - r_{3,4})$	$\nu$ (C=C) <sub>as</sub>
$S_3 = (2^{-1/2})(r_{2,3} + r_{1,4})$	$\nu$ (C–C) <sub>s</sub>
$S_4 = (2^{-1/2})(r_{2,3} - r_{1,4})$	$\nu$ (C–C) <sub>as</sub>
$S_5 = (1/2)(r_{4,5} + r_{2,7})$	$\nu$ (C–H) <sub>s</sub>
$S_6 = (1/2)(r_{4,5} - r_{2,7})$	$\nu$ (C–H) <sub>as</sub>
$S_7 = (1/2)(r_{3,6} + r_{1,8})$	$\nu$ (C–CH <sub>3</sub> ) <sub>s</sub>
$S_8 = (1/2)(r_{3,6} - r_{1,8})$	$\nu$ (C–CH <sub>3</sub> ) <sub>as</sub>
$S_9 = (6^{-1/2})(r_{9,6} + r_{10,6} + r_{11,6} + r_{12,8} + r_{13,8} + r_{14,8})$	$\nu(CH_3)_s$
$S_{10} = (6^{-1/2})(r_{9,6} + r_{10,6} + r_{11,6} - r_{12,8} - r_{13,8} - r_{14,8})$	$\nu(CH_3)'_s$
$S_{11} = (12^{-1/2})(2r_{9,6} - r_{10,6} - r_{11,6} + 2r_{12,8} - r_{13,8} - r_{14,8})$	$\nu(CH_3)_{as}$
$S_{12} = (12^{-1/2})(2r_{9,6} - r_{10,6} - r_{11,6} - 2r_{12,8} + r_{13,8} + r_{14,8})$	$\nu(CH_3)'_{as}$
$S_{13} = (1/2)(r_{10.6} - r_{11.6} + r_{13.8} - r_{14.8})$	$\nu(CH_3)_{as}''$
$S_{14} = (1/2)(r_{10.6} - r_{11.6} - r_{13.8} + r_{14.8})$	$\nu(CH_3)_{as}^{\prime\prime\prime}$
$S_{15} = (1/2)(\beta_{2,4,1} + \beta_{2,4,3} - \beta_{1,3,2} - \beta_{1,3,4})$	$\delta$ ring
$S_{16} = (1/2)(\beta_{5,1,4} - \beta_{5,3,4} + \beta_{7,3,2} - \beta_{7,1,2})$	$\delta$ (C–H) <sub>s</sub>
$S_{17} = (1/2)(\beta_{6,4,3} - \beta_{6,2,3} + \beta_{8,2,1} - \beta_{8,4,1})$	$\delta(C-CH_3)_s$
$S_{18} = (1/2)(\beta_{5,1,4} - \beta_{5,3,4} - \beta_{7,3,2} + \beta_{7,1,2})$	$\delta(C-H)_{as}$
$S_{19} = (1/2)(\beta_{6,4,3} - \beta_{6,2,3} - \beta_{8,2,1} + \beta_{8,4,1})$	$\delta(C-CH_3)_{as}$
$S_{20} = (12^{-1/2})(\beta_{9,10,6} + \beta_{10,11,6} + \beta_{11,9,6} - \beta_{9,3,6} - \beta_{10,3,6} - \beta_{11,3,6})$	$\delta(CH_3)_8$
$ \begin{array}{c} \beta_{20} = (12) & \beta_{10}(0,0) + \beta_{10}(1,0) + \beta_{11}(9,0) & \beta_{9}(3,0) & \beta_{10}(3,0) \\ + \beta_{12}(13,8) + \beta_{13}(14,8) + \beta_{14}(12,8) - \beta_{12}(1,8) - \beta_{13}(1,8) - \beta_{14}(1,8) \end{array} $	0(0113)8
$S_{21} = (12^{-1/2})(\beta_{9,10,6} + \beta_{10,11,6} + \beta_{11,9,6} - \beta_{9,3,6} - \beta_{10,3,6} - \beta_{11,3,6})$	$\delta(CH_3)'_s$
$ \begin{array}{c} \sigma_{21} = (12 - \beta_{12,10,8} + \beta_{10,11,6} + \beta_{11,9,6} + \beta_{9,5,6} + \beta_{10,5,6} + \beta_{10,5,6} \\ - \beta_{12,13,8} - \beta_{13,14,8} - \beta_{14,12,8} + \beta_{12,1,8} + \beta_{13,1,8} + \beta_{14,1,8} \end{array} $	0(CH3) <sub>8</sub>
$S_{22} = (12^{-1/2})(2\beta_{9,10,6} - \beta_{10,11,6} - \beta_{11,9,6} + 2\beta_{12,13,8} - \beta_{13,14,8} - \beta_{14,12,8})$	$\delta(CH_3)_{as}$
$S_{23} = (12^{-1/2})(2\beta_{9,10,6} - \beta_{10,11,6} - \beta_{11,9,6} - 2\beta_{12,13,8} + \beta_{13,14,8} + \beta_{14,12,8})$	$\delta(CH_3)'_{as}$
$S_{24} = (1/2)(\beta_{10,11,6} - \beta_{11,9,6} + \beta_{13,14,8} - \beta_{14,12,8})$	$\delta(CH_3)_{as}''$
$S_{25} = (1/2)(\beta_{10,11,6} - \beta_{11,9,6} - \beta_{13,14,8} + \beta_{14,12,8})$ $S_{25} = (1/2)(\beta_{10,11,6} - \beta_{11,9,6} - \beta_{13,14,8} + \beta_{14,12,8})$	$\delta(CH_3)_{as}^{m}$
$S_{26} = (12^{-1/2})(2\beta_{9,3,6} - \beta_{10,3,6} - \beta_{11,3,6} + 2\beta_{12,1,8} - \beta_{13,1,8} - \beta_{14,1,8})$	$\gamma(CH_3)_{as}$
$S_{27} = (12^{-1/2})(2\beta_{9,3,6} - \beta_{10,3,6} - \beta_{11,3,6} - 2\beta_{12,1,8} + \beta_{13,1,8} + \beta_{14,1,8})$	$\gamma(CH_3)'_{as}$
$S_{28} = (1/2)(\beta_{10,3,6} - \beta_{11,3,6} + \beta_{13,1,8} - \beta_{14,1,8})$	$\gamma(CH_3)_{as}^{\prime\prime}$
$S_{29} = (1/2)(\beta_{10,3,6} - \beta_{11,3,6} + \beta_{13,1,8} - \beta_{14,1,8})$ $S_{29} = (1/2)(\beta_{10,3,6} - \beta_{11,3,6} - \beta_{13,1,8} + \beta_{14,1,8})$	$\gamma(CH_3)_{as}^{m}$
$S_{30} = (1^{-2})(\rho_{10,3,6} - \rho_{11,3,6} - \rho_{13,1,8} + \rho_{14,1,8})$ $S_{30} = (2^{-1/2})(\tau_{1,2,3,4} + \tau_{3,4,1,2})$	
$S_{30} = (2^{-1/2})(\tau_{1,2,3,4} + \tau_{2,4,1,2})$ $S_{31} = (2^{-1/2})(\tau_{4,3,6,9} + \tau_{2,1,8,12})$	au ring $ au$ (C–CH <sub>3</sub> )
$S_{31} = (2^{-1/2})(\tau_{4,3,6,9} + \tau_{2,1,8,12})$ $S_{32} = (2^{-1/2})(\tau_{4,3,6,9} - \tau_{2,1,8,12})$	· · · · · ·
$S_{32} = (2^{-1/2})(\gamma_{5,1,4,3} + \gamma_{7,3,2,1})$ $S_{33} = (2^{-1/2})(\gamma_{5,1,4,3} + \gamma_{7,3,2,1})$	$ au(C-CH_3)'$ $\gamma(C-H)_s$
$S_{33} = (2^{-1/2})(\gamma_{5,1,4,3} + \gamma_{7,3,2,1})$ $S_{34} = (2^{-1/2})(\gamma_{5,1,4,3} - \gamma_{7,3,2,1})$	• • • • •
$S_{34} = (2 - \gamma)(\gamma_{5,1,4,3} - \gamma_{7,3,2,1})$ $S_{-} = (2^{-1/2})(\gamma_{5,-1,4,3} - \gamma_{7,3,2,1})$	$\gamma$ (C–H) <sub>as</sub>
$S_{35} = (2^{-1/2})(\gamma_{6,4,3,2} + \gamma_{8,2,1,4})$	$\gamma$ (C-CH <sub>3</sub> ) <sub>s</sub>
$S_{36} = (2^{-1/2})(\gamma_{6,4,3,2} - \gamma_{8,2,1,4})$	$\gamma$ (C-CH <sub>3</sub> ) <sub>as</sub>
$S_{37} = (2^{-1/2})(r_{15,16} + r_{15,17})$	$\nu(CO_2)_s$
$S_{38} = (2^{-1/2})(r_{15,16} - r_{15,17})$	$\nu(CO_2)_{as}$
$S_{39} = \lim_{t \to 2} 1_{17,16,15,3}$	$\delta(CO_2)_s$
$S_{40} = \lim_{t \to 1} \frac{1}{17,16,15,3}$	$\delta(CO_2)_{as}$
$S_{41} = (2^{-1/2})(\beta_{4,16,1} + \beta_{2,17,3})$	CO <sub>2</sub> -ring twist
$S_{42} = (2^{-1/2})(r_{4,17} - r_{2,16})$	CO <sub>2</sub> -ring wag
$S_{43} = (2^{-1/2})(r_{3,17} + r_{1,16})$	CO <sub>2</sub> -ring symmetric stretching
$S_{44} = (2^{-1/2})(r_{3,17} - r_{1,16})$	CO <sub>2</sub> -ring asymmetric stretching
$S_{45} = (2^{-1/2})(\beta_{4,16,1} - \beta_{2,17,3})$	CO <sub>2</sub> -ring parallel shift
$\frac{1}{2}$ is the distance between storms A and A + $\theta$ is the angle between vectors A A and A	A is the dihadral angle between the plane

<sup>a</sup>  $r_{i,j}$  is the distance between atoms  $A_i$  and  $A_j$ ;  $\beta_{i,j,k}$  is the angle between vectors  $A_kA_i$  and  $A_kA_j$ ;  $\tau_{i,j,k,l}$  is the dihedral angle between the plane defined by  $A_i$ ,  $A_j$ ,  $A_k$  and the plane defined by  $A_j$ ,  $A_k$ ,  $A_l$  atoms;  $\gamma_{i,j,k,l}$  is the angle between the vector  $A_kA_i$  and the plane defined by atoms  $A_j$ ,  $A_k$ ,  $A_l$  atoms;  $\gamma_{i,j,k,l}$  is the angle between the vector  $A_kA_i$  and the plane defined by atoms  $A_j$ ,  $A_k$ ,  $A_l$ ;  $\ln 1_{i,j,k,l}$  collinear bending  $A_i$ - $A_j$ - $A_k$  distorted in the plane of  $A_iA_jA_l$ ;  $\ln 2_{i,j,k,l}$  linear bending  $A_i$ - $A_k$ - $A_j$  distorted perpendicular to the plane  $A_iA_jA_l$ .

predict these bands at 828.0, 761.9 and 503.6 cm<sup>-1</sup>, respectively. To the best of our knowledge, this is the first positive identification of a Dewar form of an  $\alpha$ -pyrone based on the analysis of the whole mid-infrared region of its spectrum.

Norrish type I,  $\alpha$ -cleavage processes leading to the open-ring ketene species are believed [14,15] to originate from the excited states with  $n\pi^*$  character, whereas formation of Dewar isomers should proceed starting from the excited  $\pi\pi^*$  state. In  $\alpha$ -pyrones, which are six-membered ring compounds with the double bond of the carbonyl

group linearly conjugated with two double bonds in the ring, the lowest excited singlet state has  $n\pi^*$  character [3]. Hence, the  $n\pi^*$ -type photochemistry should be favored for such compounds, especially when they are free of any substituents. On the other hand, methyl substituents attached to the ring lead to an effective extension of the  $\pi$ -electron system of a molecule, by hyperconjugation. As a consequence, the energy gap between the ground state and the  $\pi\pi^*$  excited state should diminish and the gap between the ground and the  $n\pi^*$  state should increase [3,16]. For compounds

Observed and calculated vibrational frequencies, intensities and potential energy distributions (PED) for DMAP monomer (frequencies ( $\nu$ ) in cm<sup>-1</sup>, theoretical intensities (I) in km mol<sup>-1</sup>)

Observed (infrared, Ar matrix, $T = 10 \text{ K}$ )		Calculated (B3LYP/6-311++G(d,p))		Symmetry	PED <sup>a</sup> (%)
ν	I <sup>b</sup>	ν <sup>c</sup>	Ι	-	
		3100.5	0.2	A'	v(C3–H17) (91.0)
3074.7	3.0	3098.0	5.1	A'	v(C5-H12) (90.9)
3012.5	6.7	3029.0	10.0	A'	$\nu(CH_3)_{as}^{1'}$ (96.8)
2991.8	14.9	3019.5	15.3	A'	$\nu(CH_3)_{as}^{2'}$ (95.8)
		2989.1	6.3	Α″	$\nu(CH_3)_{as}^{1''}$ (100.3)
2967.1	13.9	2976.6	10.3	A″	$\nu(CH_3)^{2''}_{as}$ (100.3)
2935.2		2938.1	11.3	A'	$\nu(CH_3)^1_s$ (96.8)
2932.0	11.1	2929.6	13.9	A'	$\nu(CH_3)_8^2$ (95.7)
802.4	22.5				
783.7					
778.6	60.1				
761.3	626.3	1768.8	730.6	A'	v(O7=C2) (81.2)
661.5					
657.4	106.6	1647.2	107.5	A'	$\nu$ (C6=C5) (38.8) + $\nu$ (C4=C3) (25.4)
578.3					
574.0	161.2	1559.1	174.6	A'	$\nu$ (C4=C3) (38.8) + $\nu$ (C6=C5) (22.4)
453.9	19.5	1458.8	12.1	A'	$\delta(CH_3)^{l'}_{as}(49.1) + \delta(CH_3)^{2'}_{as}(22.6)$
444.2	27.7	1450.1	27.5	A'	$\delta(CH_3)^{2'}_{as}(59.4) + \delta(CH_3)^{1'}_{as}(24.5)$
439.4	7.8	1449.0	8.6	A''	$\delta(CH_3)^{2''}_{as}$ (92.1)
433.3	11.8	1437.6	9.0	A″	$\delta(CH_3)_{as}^{1''}$ (91.2)
401.9	31.3				
399.1		1398.4	18.3	A'	$\delta(\text{CH}_3)^3_8(18.8) + \nu(\text{C5-C4})(17.6) + \delta(\text{CH}_3)^{1/}_{as}(12.7) + \delta(\text{C3-H17})(11.1)$
385.3	7.5	1387.2	5.0	A'	$\delta(CH_3)^1_8$ (89.1)
379.6	1.9				
374.7	16.7	1378.0	17.0	A'	$\delta(CH_3)^2_{s}$ (72.8)
331.7	21.5				
329.1		1320.8	26.2	A'	$\delta$ (C5–H12) (23.5) + ν(O1–C6) (15.7) + ν(C3–C2) (11.2) + $\delta$ (C3–H17) (9.8)
308.8	10.3				
266.7	24.4				
251.0	Weak				
244.3	4.9				
239.1					
231.7	30.2	1226.3	13.3	A′	$\delta$ (C3-H17) (38.9) + $\nu$ (C8-C6) (14.0) + $\delta$ (C5-H12) (11.3)
229.6					
179.6	6.0				
170.7	7.3	1144.0	10.0	•/	
151.2	41.3	1144.9	12.3	A'	$\delta$ (C5-H12) (32.1) + $\nu$ (C4-C13) (16.2) + $\delta$ (C3-H17) (15.0) - $\nu$ (O1 C6) (10.7)
134.7	32.7	1103.8	8.1	A′	$\nu$ (O1–C6) (10.7) $\nu$ (C3–C2) (42.1)
134.7	32.1	1105.8	0.1	Α	$\nu(C_{3}-C_{2})$ (42.1)
130.8	0.8				
.097.7	2.2				
071.1	2.2	1039.8	0.04	Α″	$\gamma$ (CH <sub>3</sub> ) <sup>1"</sup> <sub>as</sub> (47.5) + $\gamma$ (CH <sub>3</sub> ) <sup>2"</sup> <sub>as</sub> (30.5)
026.2	11 1			A" A"	$\gamma(CH_3)_{as}^{2''}(45.6) + \gamma(CH_3)_{as}^{2''}(30.5)$ $\gamma(CH_3)_{as}^{2''}(45.6) + \gamma(CH_3)_{as}^{1''}(30.5)$
036.3	11.1	1034.2	7.0		
030.4	41.4	1005.8	23.7	A'	$\gamma$ (CH <sub>3</sub> ) <sup>2</sup> / <sub>as</sub> (36.1) + $\nu$ (O1–C6)(19.7) + $\delta$ ring 1(9.6)
023.1	4.0	1022.2	3.1	A'	$\gamma(CH_3)^{1'}_{as}(53.0) + \gamma(CH_3)^{2'}_{as}(14.9)$
985.9 076.0	4.2				
976.0	2.4	050 5	2.5	• /	1/20 = 1/20 = 1000 = 0.000 = 0.000 = 0.000 = 0.00000 = 0.00000 = 0.00000 = 0.00000 = 0.00000 = 0.00000 = 0.00000 = 0.00000000
965.5	4.8	958.7	3.5	A′	$\delta \operatorname{ring} 1(29.6) + \nu(\text{C5-C4})(29.2) + \gamma(\text{CH}_3)^{2'}_{as}(19.2)$
958.5	30.0	0.42.2	10.0	• /	
954.5		943.3	19.9	A'	$\nu$ (C8–C6)(22.9) + $\nu$ (C13–C4) (22.6) + $\nu$ (O1–C6) (17.8)
899.7	6.7				
869.2	0.7				
857.9 848 2	2.3	846.0	10.2	A.//	$v_{1}(C_{2}, U_{1}(T_{1}), (0, 0)) + v_{1}(C_{2}(C_{2}), (1, 1, 0))$
848.2 843.6	29.3 58.4	846.9 803.0	18.3 64.5	A'' A'	$\gamma$ (C3–H17) (90.9) + $\gamma$ (C2=O7) (14.9) $\gamma$ (C2–O1) (55.0) + $\delta$ ring 1 (12.5)
	20.4	005.0	04.3	A	$\nu$ (C2–O1) (55.0) + $\delta$ ring 1 (13.5)
810.9	33.3	804.8	29.3	A″	γ(C5–H12) (100.1)

Observed (infrared, Ar matrix, $T = 10 \text{ K}$ )		Calculated (B3LYP/6-311++G(d,p))		Symmetry	PED <sup>a</sup> (%)
ν	Ip	ν <sup>c</sup>	Ι		
783.4	0.8				
757.9	0.5				
730.2	1.4	719.8	2.1	Α″	$\gamma$ (C2=O7) (60.9) + $\gamma$ ring 1 (43.8)
638.2	6.2				
617.8	5.0	611.5	2.9	Α″	$\gamma \operatorname{ring} 1 (24.7) + \gamma (C4-C13) (18.0) + \gamma \operatorname{ring} 3 (17.6) + \gamma (C6-C8) (17.0) + \gamma (C2=O7) (15.9)$
580.8	2.1	594.1	0.3	A′	$\delta \operatorname{ring} 1 (17.9) + \nu(C2-O1) (13.4) + \delta \operatorname{ring} 3 (13.0) + \nu(C13-C4) (11.5)$
554.1	7.0	543.8	5.0	A'	$\delta$ (C2=O7) (32.1) + $\delta$ ring 3 (21.3) + $\nu$ (C8–C6) (9.8)
534.4	1.8				
516.7	2.7	509.6	2.3	A'	$\delta$ ring 3 (44.0) + $\delta$ (C2=O7) (18.8)
508.0	3.4	509.1	2.9	Α″	$\gamma$ (C6–C8) (35.1) + $\gamma$ (C4–C13) (34.6) + $\gamma$ ring 2 (20.8)
472.1	2.5	464.7	0.9	A'	$\delta$ ring 2 (76.2)
		317.5	1.1	A′	$\delta$ (C4–C13) (33.5) + $\delta$ (C6–C8) (23.1) + $\delta$ (C2=O7) (20.0)
		266.4	0.6	A′	$\delta$ (C6–C8) (43.2) + $\delta$ (C4–C13) (37.8)
		201.6	2.4	Α″	$\gamma$ ring 2 (76.2) + twist (CH <sub>3</sub> ) <sup>2</sup> (10.6)
		179.8	0.8	Α″	Twist $(CH_3)^1$ (39.7) + $\gamma$ ring 3 (35.0) + $\gamma$ (C4–C13) (16.2)
		150.0	0.1	Α″	$\gamma \operatorname{ring} 1 (45.7) + \gamma \operatorname{ring} 2 (23.8) + \gamma (C6-C8) (13.8) + \operatorname{twist} (CH_3)^2 (9.8) + \gamma \operatorname{ring} 3 (9.5)$
		130.6	1.2	Α″	Twist $(CH_3)^1$ (47.9)+ $\gamma$ ring 3 (27.6) + Twist $(CH_3)^2$ (16.6)
		115.7	0.2	A″	Twist $(CH_3)^2$ (63.1)+ $\gamma$ ring 3 (23.4)

Table 4 (Continued)

<sup>a</sup> PED's lower than 10% are not included. Definition of symmetry coordinates is given in Table 1. See Fig. 1 for atom numbering.

<sup>b</sup> Relative integrated intensities.

<sup>c</sup> Theoretical positions of absorption bands above  $2900 \text{ cm}^{-1}$  were scaled down by a factor of 0.968 and below  $2000 \text{ cm}^{-1}$  by a factor 0.978.

substituted with methyl groups, this effect should promote photoreactions originating from the  $\pi\pi^*$  state, whereas the reactions typical for the  $n\pi^*$  states (prevailing for unsubstituted species) should be hindered. Strong domination of the Dewar form creation over the  $\alpha$ -bond cleavage, observed in the current work for 4,6-dimethyl- $\alpha$ -pyrone, illustrates well the influence of methyl substitution on relative effectiveness of the  $\pi\pi^*$  and  $n\pi^*$  photochemical channels.

When the UV ( $\lambda > 315$  nm) irradiated matrix was subsequently subjected to shorter wavelength UV radiation ( $\lambda >$ 

235 nm), evolution of  $CO_2$  was observed and 1,3-dimethylcyclobutadiene (DMCB) was created. The  $CO_2$  produced from DOOBH fragmentation could be easily identified spectroscopically by observation of the very intense characteristic IR absorption region around 2337.5 cm<sup>-1</sup>.

After photolysis, the photoproduced  $CO_2$  and DMCB must be confined in the same matrix cavity and they form a complex. The DFT(B3LYP)/6-311++G(d,p) calculations predict that the CO<sub>2</sub>-DMCB complex has a stacking (parallel) geometry (Fig. 6). Fig. 7 represents two potential scans

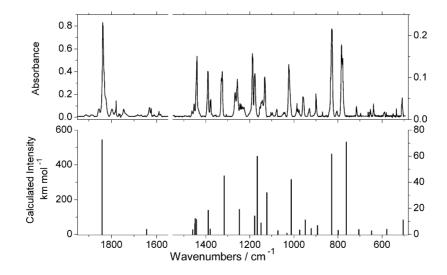


Fig. 5. Comparison of the experimental IR spectrum of the Dewar valence isomer (DOOBH), photogenerated upon UV ( $\lambda > 315$  nm) irradiation of DMAP isolated in Ar matrix (10 K), with the spectrum of DOOBH theoretically predicted at the DFT(B3LYP)/6-311++G(d,p) level. Theoretical frequencies were scaled using a factor equal to 0.978. Note the change of ordinate scale at 1500 cm<sup>-1</sup>.

Observed and calculated vibrational frequencies, intensities and potential energy distributions (PED) for DOOBH monomer (frequencies ( $\nu$ ) in cm<sup>-1</sup>, theoretical intensities (I) in km mol<sup>-1</sup>)

Observed (infrared, Ar matrix, $T = 10 \text{ K}$ )		Calculated (B3LYP/6-311++G(d,p))		PED <sup>a</sup> (%)	
ν	ľ	ν <sup>c</sup>	Ι		
3062.5	1.5	3091.8	8.3	v(C5-H12) (99.1)	
3005.7		3016.5	11.2	v(C3–H17) (99.4)	
2997.7	35.7	3011.1	13.0	$\nu(CH_3)_{as}^{l''}$ (97.4)	
2990.1		3005.9	14.1	$\nu(CH_3)_{as}^{2'}$ (97.7)	
2984.4		2998.6	12.0	$\nu(CH_3)_{as}^{l'}$ (97.1)	
2958.8	8.3	2971.3	8.7	$\nu$ (CH <sub>3</sub> ) <sup>2</sup> <sup>a''</sup> <sub>as</sub> (100.2)	
940.8	8.2	2936.4	13.9	$\nu(CH_3)_{s}^{l}$ (99.3)	
937.5					
923.8	4.7	2924.2	14.6	$\nu(CH_3)_s^2$ (97.6)	
854.1	22.3	2/2112	1 110		
837.7	444.8	1841.6	543.7	v(C2=O7) (88.1)	
631.5	41.0	1644.7	28.9	$\nu(C5=C4)$ (76.1) + $\nu(C4=C13)$ (12.9)	
.623.8	41.0	1044.7	20.7	$v(e_{3},e_{4})(v(e_{4},e_{1}))(v(e_{4},e_{1}$	
462.2	2.8	1458.7	3.5	$\delta(CH_3)_{as}^{l''}$ (84.2)	
452.8	4.8	1449.5	8.0	$\delta(CH_3)_{as}^{as}$ (64.2) $\delta(CH_3)_{as}^{ls}$ (80.5)	
				$\delta(CH_3)_{as}(00.3)$ $\delta(CH_3)_{as}^{2'}$ (84.3)	
1439.9	42.5	1447.9	12.2		
1000 6	20.7	1441.9	11.4	$\delta(CH_3)_{as'}^{2s'}$ (89.8)	
1389.6	29.7	1388.4	18.5	$\delta(CH_3)^1_s$ (83.9) + $\nu(C6-C8)$ (10.7)	
1379.2	0.7	1050 (			
1376.7	9.7	1379.6	4.1	$\delta(CH_3)^2_{s}$ (93.9)	
1328.0	41.8				
1324.5		1316.0	44.8	$\nu$ (C6–C8) (22.6) + $\nu$ (C5–C6) (18.3) + $\nu$ (C3–C6) (14.6) + $\delta$ (CH <sub>3</sub> ) <sup>1</sup> <sub>s</sub> (12.3)	
				$+\delta(C5-H12)$ (11.1)	
266.3				2/	
255.9	50.9	1247.0	19.1	$\nu$ (C3–C4) (21.1) + $\gamma$ (CH <sub>3</sub> ) <sup>2'</sup> <sub>as</sub> (11.3)	
187.7	42.4	1177.4	14.1	γ(C3–H17) (55.7)	
177.3	37.1	1165.5	59.7	$\nu$ (C2–O1) (20.8) + $\delta$ (C5–H12) (11.5) + $\delta$ (C3–H17) (10.7) + $\gamma$ (CH <sub>3</sub> ) <sup>"</sup> <sub>as</sub> (9.7)	
154.4	28.3				
143.9		1148.5	8.8	$\gamma$ (CH <sub>3</sub> ) <sup>l'</sup> <sub>as</sub> (16.3) + $\gamma$ (CH <sub>3</sub> ) <sup>l''</sup> <sub>as</sub> (15.9) + $\delta$ (C5–H12) (19.1)	
132.0	34.4	1122.1	32.0	$\delta$ (C5–H12) (17.4) + $\nu$ (C4–C13) (17.1) + $\gamma$ (C3–H17) (15.8) + $\delta$ (C3–H17) (10.3)	
077.3	3.3	1072.0	2.8	$\delta$ (C3–H17) (30.0) + $\nu$ (C6–C8) (13.5)	
1046.0	5.1	1031.1	0.9	$\gamma(CH_3)^{2''}_{as}$ (69.5) + $\gamma(C4$ -C13) (10.8)	
1040.0					
1022.3	48.2	1011.8	42.0	$\gamma(CH_3)_{as}^{l''}$ (29.7) + $\nu(C2-O1)$ (21.0) + $\gamma(CH_3)_{as}^{l'}$ (15.1)	
986.0		973.6	3.4	$\gamma(CH_3)^{2'}_{as}$ (21.3) + $\delta(C3-H17)$ (19.8) + $\delta$ ring 2 (11.3) + $\nu(C4-C13)$ (10.5)	
981.9	12.0				
976.0					
957.8	17.3	948.0	11.1	$\gamma(CH_3)_{as}^{2'}$ (25.9) + $\nu(C2-C3)$ (12.7) + $\delta$ ring 1 (11.1)	
928.9	5.4	920.8	4.6	$\nu$ (C4–C3) (17.0) + $\nu$ (C5–C6) (16.7) + $\nu$ (C2–C3) (13.6) + $\delta$ ring 2 (12.1)	
900.0	9.8	892.9	6.7	$\gamma(CH_3)_{as}^{l'}(20.5) + \nu(C5-C6) (11.6) + \nu(C3-C6) (10.3)$	
827.4	71.3	828.0	61.6	$\gamma$ (C5-H12) (76.2)	
804.8	3.2	799.0	3.1	$\gamma(C2=O7)(23.4) + \nu(C3=C6)$ (17.5)	
783.6	74.9	761.9	70.7	$\nu(01-C6)$ (53.8) + $\nu(C2-O1)$ (22.8)	
778.4	//	,01.9	, 0.7	.(01 00) (00.0) + (02 01) (22.0)	
715.1	2.9	705.2	3.8	$\gamma$ (C2=O7) (15.4) + $\delta$ (C2=O7) (13.0) + $\tau$ ring 2 (11.1)	
638.1	4.1	646.4	2.7	$\delta \operatorname{ring} 1 (21.4) + \nu(C4-C13) (14.3) + \nu(C3-C4) (11.8)$	
588.9	4.1	577.9	4.0	$\nu$ (C6–C8) (25.2) + $\delta$ ring 2 (20.0) + $\nu$ (C4–C13) (13.4)	
509.3	10.3	503.6	4.0	$\gamma(C2=O7)$ (20.4) + $\tau$ ring 2 (18.8) + $\delta(C2=O7)$ (18.6)	
443.3	8.0	438.7	6.7	$\gamma(\text{C2-C7})$ (20.4) + $\tau$ fing 2 (18.8) + $\delta(\text{C2-C7})$ (18.6) $\gamma(\text{C4-C13})$ (34.1) + $\tau$ butterfly (32.1)	
428.2	2.9	421.4	2.1	$\delta(C2=O7)$ (23.7) + $\tau$ butterfly (12.7) + $\nu(O1-C6)$ (9.8) + $\tau$ ring 2 (9.8)	
		346.5	3.6	$\delta$ (C4–C13) (30.8) + $\delta$ (C6–C8) (18.1) + $\tau$ butterfly (11.4)	
		316.8	1.9	$\gamma$ (C6–C8) (42.2)	
		228.0	0.2	$\delta$ (C6–C8) (46.8) + $\delta$ (C4–C13) (31.3)	
		213.6	0.05	Twist $(CH_3)^1$ (92.9)	
		162.6	0.06	Twist $(CH_3)^2$ (84.7)	
		133.1	1.1	$\tau$ ring 1 (29.0) + $\gamma$ (C6–C8) (21.7) + $\tau$ ring 2 (18.0) + Twist (CH <sub>3</sub> ) <sup>2</sup> (13.2)	
		125.7	3.3	$\gamma$ (C4–C13) (25.3) + $\tau$ butterfly (20.7) + $\tau$ ring 1 (14.9) + $\gamma$ (C2=O7) (13.6) +	
				ring 2 (11.2)	

<sup>a</sup> PED's lower than 10% are not included. Definition of symmetry coordinates is given in Table 2. See Fig. 3 for atom numbering.

<sup>b</sup> Relative integrated intensities.

 $^{c}$  Theoretical positions of absorption bands above 2900 cm $^{-1}$  were scaled down by a factor of 0.968 and below 2000 cm $^{-1}$  by a factor 0.978.

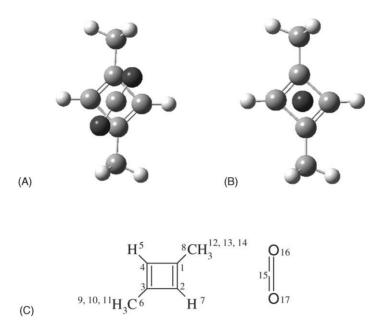


Fig. 6. Dimethylcyclobutadiene complex with CO<sub>2</sub>. (A) Stacking (parallel) dimer, corresponding to the unique theoretically predicted minimum on the PES of this system. (B) Axial configuration of the dimer (not a minimum; see text). (C) Atom numbering used in the normal mode analysis of the complex.

for the CO<sub>2</sub>–DMCB complex assuming two different orientations within the C<sub>2</sub> symmetry framework: one with the axial orientation and the other with the parallel orientation to the DMCB plane (see also Fig. 6). The first scan (axial) reveals a repulsive potential of interaction between CO<sub>2</sub> and DMCB. On the other hand, the second one (parallel) shows a binding profile with an equilibrium minimum energy distance between the two molecules of ca. 3.44 Å and a stabilization energy of the complex equal to  $4.7 \text{ kJ mol}^{-1}$ . The main geometric parameters for both DOOBH and  $CO_2$ -DMCB complex are listed in Tables 6 and 7. Calculations performed at the DFT(B3LYP)/6-311++G(d,p) level of theory do not predict any other minima for the  $CO_2$ -DMCB system.

Fig. 8 clearly demonstrates the formation of the CO<sub>2</sub>– DMCB complex from DOOBH. DOOBH was produced first by irradiation of the matrix with UV light ( $\lambda > 315$  nm) and dominated in the sample at this stage of the experiment

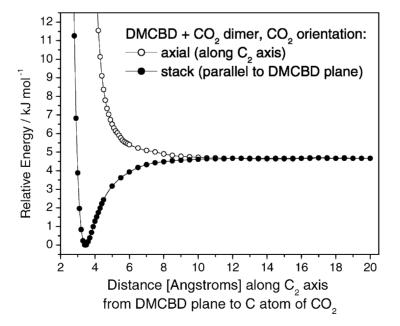


Fig. 7. Energy of the dimethylcyclobutadiene complex with  $CO_2$  as a function of the distance between the two molecules. Energies were calculated at the DFT(B3LYP)/6-311++G(d,p) level for the structures obtained by fixing the distance between DMCBD and  $CO_2$  and optimizing all other geometry parameters. Two relative orientations of the two molecules (parallel and axial—see Fig. 6) have been considered.  $C_2$  symmetry axis was conserved in all the calculations. Zero-level of energy corresponds to -422.039897 hartrees.

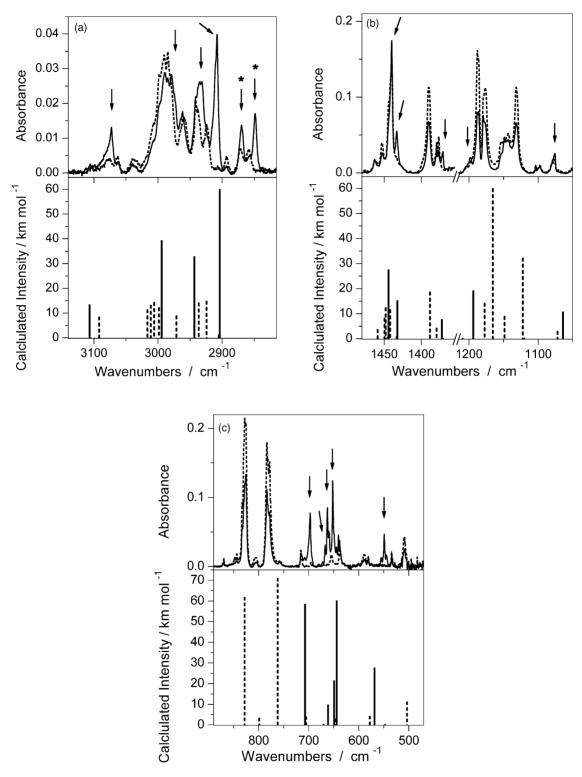


Fig. 8. Experimental evidence of photogeneration of the DMCB–CO<sub>2</sub> complex upon UV ( $\lambda > 235$  nm) irradiation of DOOBH isolated in Ar matrix (10 K). The experimental spectrum presented (in upper panel) with dashed line is dominated by the bands due to Dewar valence isomer (DOOBH), produced by preceding UV ( $\lambda > 315$  nm) irradiation. The spectrum presented with solid line was recorded after 360 min of subsequent UV ( $\lambda > 235$  nm) irradiation. The growing bands correspond to the DMCB–CO<sub>2</sub> dimer and are indicated with arrows. Asterisks indicate bands assigned to overtones or combination tones. The spectra theoretically predicted at the DFT(B3LYP)/6-311++G(d,p) level are presented (in lower panel) in the stick spectrum form. In all the regions dashed sticks correspond to the spectrum of DOOBH and solid sticks to the spectrum of the DMCB–CO<sub>2</sub> stacking complex. Theoretical frequencies were scaled using a factor equal to 0.978 (except for the CH stretching region, where the factor 0.968 was used).

Bond distances, bond angles in four-membered rings, and total electronic energy (E(RB + HF – LYP) = -422.0621471, in hartrees) at the geometry optimized at the B3LYP/6-311++G(d,p) level for DOOBH (atom numbering as in Fig. 3)

O1-C2	1.3695	C8–H9	1.0928
C2-C3	1.5313	C8-H10	1.0937
C3–C4	1.5352	C8-H11	1.0920
C4=C5	1.3435	C5-H12	1.0833
C5-C6	1.5151	C4-C13	1.4858
C6-O1	1.5078	C13-H14	1.0919
C6–C3	1.5523	C13-H15	1.0951
C2=O7	1.1926	C13-H16	1.0951
C6–C8	1.5001	C3-H17	1.0892
Bond angles <sup>a</sup> (°)			
C6-O1-C2	91.77	C6-C5-C4	93.90
C6-C3-C2	84.23	C6-C3-C4	85.32
O1-C6-C3	88.89	C5-C4-C3	93.94
O1-C2-C3	95.09	C5-C6-C3	86.83

<sup>a</sup> A–B–C is the angle between bonds A–C and B–C.

(dashed line in Fig. 8). Further irradiation with shorter wavelengths led to decrease of the bands originated in DOOBH and appearance of new bands due to the CO<sub>2</sub>-DMCB complex (the solid line in Fig. 8 corresponds to the spectrum obtained after 360 min of irradiation with  $\lambda > 235$  nm). A close inspection of the observed profile of the characteristic  $CO_2$  feature at 2341.7–2337.7 cm<sup>-1</sup>, due to the  $CO_2$  asymmetric stretching vibration is presented in Fig. 9, and clearly reveals that CO<sub>2</sub> must be associated. Indeed, this feature is observed as a broad complex band, which is red-shifted relatively to the doublet observed for the free CO<sub>2</sub> monomer isolated in argon. The later has components at 2345.0 and  $2339.1 \text{ cm}^{-1}$ , with full width at half maximum equal to only  $0.2 \text{ cm}^{-1}$  [17]. As shown in Fig. 8, the IR spectrum calculated for the stacked CO<sub>2</sub>-DMCB complex fits nicely the experimental spectrum of the photoproduct(s) generated

Table 7

Bond distances, bond angles in four-membered ring, and total electronic energy (E(RB + HF - LYP) = -422.03989725 Energy of the DMCB-CO<sub>2</sub> complex in hartrees) at the geometry optimized at the B3LYP/6-311++G(d,p) level for 1,3-dimethyl-cyclobutadiene in the DMCB-CO<sub>2</sub> complex (atom numbering as in Fig. 6)

Bond distances (Å	.)		
C1=C2	1.3382	C1–C8	1.4794
C2–C3	1.5776	C6–H9	1.0926
C3=C4	1.3382	C6-H10	1.0967
C4C1	1.5776	C6-H11	1.0972
C4-H5	1.0818	C8-H12	1.0926
C3–C6	1.4794	C8-H13	1.0967
C2-H7	1.0818	C8-H14	1.0972
Bond angles <sup>a</sup> (°)			
C1C2C3	89.67	C3-C4-C1	89.67
C2-C3-C4	90.33	C4C1C2	90.33

 $^{a}$  A–B–C is the angle between bonds A–C and B–C.

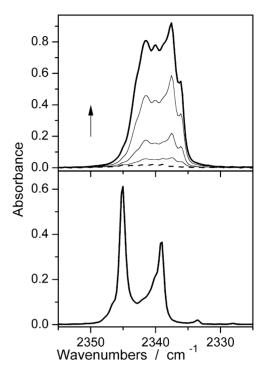


Fig. 9. Upper panel: Increase of the IR band due to CO<sub>2</sub> antisymmetric vibration during the progress (t = 0, 15, 60, 180 and 360 min) of UV ( $\lambda > 235$  nm) irradiation of DOOBH isolated in Ar matrix (10 K); lower panel: the same region of IR spectrum of a matrix obtained (in a separate experiment) by co-deposition of CO<sub>2</sub> and argon.

upon irradiation of the matrix with UV ( $\lambda > 235$  nm) light. Table 8 summarizes the proposed band assignments for this species. The most intense IR bands due to the bending vibrations of the methyl groups were found in the experimental spectrum at 1439.9, 1433.2 and  $1371.2 \text{ cm}^{-1}$  and the C-H asymmetric in-plane bending mode at 1198.3 cm<sup>-1</sup>, in good agreement with the theoretical predictions (1444.1, 1432.1, 1372.4 and 1193.6  $cm^{-1}$ , respectively), while the bands predicted at 1064.0, 1016.4 and  $661.3 \text{ cm}^{-1}$  have their experimental counterparts at 1075.6, 1012.3 and  $667.0 \,\mathrm{cm}^{-1}$ . The C–H out-of-plane bending vibration of the cyclobutadiene ring hydrogen atoms was predicted to give rise to a strong band at  $707.1 \text{ cm}^{-1}$  and was observed in the experimental spectrum at  $696.9 \text{ cm}^{-1}$ . On the other hand, the two  $CO_2$  bending vibrations, which in the  $CO_2$ monomer are degenerated by symmetry, were predicted by the calculations to appear in the CO<sub>2</sub>-DMCB at 649.2 and  $644.1 \text{ cm}^{-1}$  and, accordingly, were observed at 662.6/659.8and  $652.1 \,\mathrm{cm}^{-1}$ . The band due to the four-membered ring puckering ( $\tau$  ring; calculated frequency: 568.3 cm<sup>-1</sup>) was observed at 549.3/544.3 cm<sup>-1</sup>. This band is a direct analogue of the ring puckering band observed for unsubstituted cyclobutadiene at  $573 \text{ cm}^{-1}$  [18]. In the high frequency spectral region, the three intense methyl stretching bands were observed at 2973.7, 2931.2 and 2907.6  $cm^{-1}$ , while the C<sub>ring</sub>-H asymmetric stretching was found at 3072.4 cm<sup>-1</sup>, also in good agreement with the theoretical predictions (see

Observed and calculated vibrational frequencies, intensities and potential energy distributions (PED) for DMCB stacked to CO<sub>2</sub> monomer (frequencies  $(\nu)$  in cm<sup>-1</sup>, theoretical intensities (*I*) in km mol<sup>-1</sup>)

Observed (infrared, Ar matrix, $T = 10 \text{ K}$ )		Calculated (B3LYP/6-311++G(d,p))		Symmetry	PED <sup>a</sup> (%)
ν	ľ	v <sup>c</sup>	I		
3072.4	13.2	3106.6	13.1	В	$\nu$ (C–H) <sub>as</sub> (99.1)
		3106.4	0.03	А	$\nu$ (C–H) <sub>s</sub> (99.1)
		2994.7	0.007	А	v(CH <sub>3</sub> ) <sub>as</sub> (94.6)
2988.1				_	
2973.6	39.2	2994.3	39.1	В	$\nu(CH_3)'_{as}$ (94.5)
2962.0 2931.2	31.0	2943.2	32.6	А	$v(CH)^{\prime\prime}(00.6)$
2931.2	51.0	2943.0	0.1	B	$\nu(CH_3)'_{as}$ (99.6) $\nu(CH_3)''_{as}$ (99.6)
		2905.2	1.1	A	$\nu(CH_3)_{as} (99.0)$ $\nu(CH_3)_s (94.8)$
2907.6	32.3	2903.2	59.7	A B	$\nu(CH_3)_s (94.8)$ $\nu(CH_3)'_s (94.7)$
2907.0 2869.4	9.1	2903.8	39.1	Б	
2809.4 2848.6	9.1 11.7				Overtone $(\delta(CH_3)_{as}'')$ Overtone $(\delta(CH_3)_{as})$
2337.5	571.0	2340.7	511.1	В	$\nu(CO_2)_{as}$ (100.0)
2551.5	571.0	1663.6	32.5	B	$\nu(C=C)_{as}$ (74.4) + $\nu(C=CH_3)_{as}$ (12.3)
		1650.5	0.2	A	$\nu(C=C)_{as}(74.4) + \nu(C=CH_3)_{as}(12.5)$ $\nu(C=C)_s(77.7) + \nu(C=CH_3)_s(15.4)$
		1445.5			
1439.9	36.9		0.4 27.3	A	$\delta(CH_3)_{as}''(58.7) + \delta(CH_3)_{as}(31.2)$
1459.9	50.9	1444.1		B	$\delta(CH_3)_{as}^{\prime\prime\prime}$ (57.2) + $\delta(CH_3)_{as}^{\prime}$ (32.9)
1422.0	10.0	1433.3	0.7	B	$\delta(CH_3)'_{as}$ (58.2) + $\delta(CH_3)''_{as}$ (33.0)
1433.2	19.9	1432.1	15.0	A	$\delta(CH_3)_{as}$ (60.2) + $\delta(CH_3)''_{as}$ (31.4)
1371.2	8.3	1372.4	7.5	B	$\delta(CH_3)'_{s}$ (95.4)
		1371.0	0.03	A	$\delta(CH_3)_s$ (95.1)
		1342.6	0.1	A	$\nu(CO_2)_s$ (100.0) $\nu(C-C)_s$ (22.0) + $\nu(C-CH_2)_s$ (20.1) + $\delta(C-H)_s$ (10.4)
1100.2	07	1209.1	0.003	A	$\nu$ (C–C) <sub>s</sub> (33.0) + $\nu$ (C–CH <sub>3</sub> ) <sub>s</sub> (29.1) + $\delta$ (C–H) <sub>s</sub> (10.4)
1198.3	8.7	1193.6	19.0	В	$\delta(\text{C-H})_{\text{as}} (41.0) + \gamma(\text{CH}_3)'_{\text{as}} (21.4) + \nu(\text{C-C})_{\text{as}} (16.6) + \delta(\text{C-CH}_3)_{\text{as}} (10.6)$
		1120.4	0.004	A	$\delta$ (C–H) <sub>s</sub> (65.9) + $\gamma$ (CH <sub>3</sub> ) <sub>as</sub> (19.6)
1075.6	8.5	1064.0	10.5	В	$\delta(\text{C-H})_{\text{as}} (43.9) + \nu(\text{C-CH}_3)_{\text{as}} (30.1) + \gamma(\text{CH}_3)'_{\text{as}} (19.0)$
		1025.1	0.09	В	$\gamma(CH_3)_{as}^{''}$ (80.9) + $\gamma(C-CH_3)_{as}$ (10.3)
1012.3	2.6	1016.4	3.7	Α	$\gamma(CH_3)''_{as}$ (78.9) + $\gamma(C-CH_3)_s$ (10.2)
		997.3	0.007	Α	$\gamma$ (CH <sub>3</sub> ) <sub>as</sub> (34.8) + $\delta$ ring (28.6) + $\nu$ (C–CH <sub>3</sub> ) <sub>s</sub> (11.6)
		921.9	4.3	В	$\gamma(CH_3)'_{as}$ (40.3) + $\nu(C-CH_3)_{as}$ (35.1) + $\nu(C=C)_{as}$ (14.4)
		892.7	0.01	Α	$\nu$ (C–C) <sub>s</sub> (45.7) + $\delta$ ring (17.9) + $\gamma$ (CH <sub>3</sub> ) <sub>as</sub> (20.8) + $\delta$ (C–H) <sub>s</sub> (13.7)
696.9	37.6	707.1	58.2	Α	$\gamma$ (C–H) <sub>s</sub> (86.0) + $\gamma$ (C–CH <sub>3</sub> ) <sub>s</sub> (11.1)
		670.3	0.1	В	$\gamma$ (C–H) <sub>as</sub> (76.7)
667.0	4.6	661.3	9.5	В	$\nu$ (C–C) <sub>as</sub> (55.1) + $\nu$ (C–CH <sub>3</sub> ) <sub>as</sub> (13.4)
662.6	40.8	649.2	21.2	В	$\delta(\text{CO}_2)_{\text{as}}$ (88.5)
659.8	41.9	644.1	50.0	•	
652.1	41.8	644.1	59.9	А	$\delta(CO_2)_{s}$ (96.8)
644.5	3.5				
640.7	8.5				
555.1	3.8	569.2	27.2		
549.3 544.3	18.2	568.3	27.3	А	$\tau$ ring (76.7) + $\gamma$ (C–H) <sub>s</sub> (11.2)
511.5		547.3	0.03	А	$\delta$ ring (40.6) + $\nu$ (C–CH <sub>3</sub> ) <sub>s</sub> (37.6) + $\nu$ (C–C) <sub>s</sub> (11.4)
		363.9	0.01	А	$\delta(C-CH_3)_s$ (81.6)
		323.3	0.1	В	$\gamma$ (C-CH <sub>3</sub> ) <sub>as</sub> (92.5)
		224.3	2.2	В	$\delta$ (C-CH <sub>3</sub> ) <sub>as</sub> (79.4) + $\nu$ (C-C) <sub>as</sub> (13.1)
		209.7	0.01	В	$\tau$ (C-CH <sub>3</sub> )' (105.4)
		193.6	0.003	А	$\tau$ (C–CH <sub>3</sub> ) (95.3)
		137.9	9.5	А	$\gamma$ (C-CH <sub>3</sub> ) <sub>s</sub> (66.6) + $\tau$ ring (27.2)
		65.3	0.4	В	$CO_2$ -ring wag (104.8)
		42.0	0.005	Ā	CO <sub>2</sub> -ring symmetric stretching (98.6)
		38.7	0.003	В	$CO_2$ -ring asymmetric stretching (84.0) + $CO_2$ -ring wag (18.9)
		29.4	0.005	Ā	CO <sub>2</sub> -ring twist (99.3)
		20.2	0.0008	В	$CO_2$ -ring parallel shift (109.2) + $CO_2$ -ring asymmetric stretching (15.4)

<sup>a</sup> PED's lower than 10% are not included. Definition of symmetry coordinates is given in Table 3. See Fig. 6 for atom numbering.

<sup>b</sup> Relative integrated intensities.

<sup>c</sup> Theoretical positions of absorption bands above  $2900 \text{ cm}^{-1}$  were scaled down by a factor of 0.968 and below  $2000 \text{ cm}^{-1}$  by a factor 0.978.

Table 8 and Fig. 8). Finally, the band due to the antisymmetric stretching vibrations of the two C=C bonds is predicted to occur around  $1660 \text{ cm}^{-1}$ , but its assignment is uncertain since this spectral region does also exhibit relatively intense bands due to traces of matrix isolated water that could not be removed from the set up.

# 5. Conclusions

Several new results on photochemistry of DMAP were obtained by combined usage of matrix isolation technique, FTIR spectroscopy and theoretical simulations of IR spectra. Photochemical valence isomerization of DMAP to its Dewar isomer DOOBH strongly dominated for this molecule when compared with the ring opening, Norrish type I photoprocess leading to the conjugated ketene. In the present study, the photoproduced DOOBH was positively identified by comparison of its experimental and theoretically calculated IR spectrum. In turn, when irradiated using higher energy radiation ( $\lambda > 235$  nm) DOOBH is photolysed to DMCB plus CO<sub>2</sub>. In the matrix, these two photoproducts form a stacked (parallel) complex, that can be unequivocally identified by comparison of its experimentally observed IR spectrum with the spectrum simulated using the DFT method.

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