
#### Abstract

$A b$ initio $6-31 G^{*}$ SCF-MO calculations have been carried out on $S$-ethyl thioacrylate $\left[\mathrm{CH}_{2}=\mathrm{CHC}(\mathrm{O}) \mathrm{SCH}_{2} \mathrm{CH}_{3}\right] . \mathrm{Fully}^{2}$ optimized geometries, relative stabilities, dipole moments and harmonic force fields for several conformers of this molecule have been determined and the results compared with those for similar molecules. Together with FTIR and Raman spectroscopic data, the theoretical results demonstrate that $S$-ethyl thioacrylate exists in two different conformations about the $\mathrm{C}_{\alpha}-\mathrm{C}$ bond (the $s$-cis and $s$-trans forms, with $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ dihedral angles equal to $0^{\circ}$ and $180^{\circ}$, respectively); the $s$-cis conformation being more stable than the s-trans form by $c a .6 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for the isolated molecule. Comparison of the experimental and theoretical vibrational spectra confirms that, as concluded from our previous study on the analogous trans-S-ethyl thiocrotonate molecule (R. Fausto, P. J. Tonge and P. R. Carey, J. Chem. Soc., Faraday Trans., 1994, 90, 3491), the presence of the $s$-trans isomer of an $\alpha, \beta$-unsaturated thioester can be successfully monitored by the IR band at $c a .1170 \mathrm{~cm}^{-1}$, ascribed to the $\mathrm{C}_{\alpha}-\mathrm{C}$ stretching mode of this form. In addition, we were also able to identify some IR bands sensitive to the conformation of the ethyl group that may be used as spectroscopic probes to study conformational equilibria associated with this internal degree of freedom in more complex $S$-ethyl thioesters.


The $\alpha, \beta$-unsaturated acyl group, $-\mathrm{CH}=\mathrm{CH}-\mathrm{C}(=\mathrm{O})-$, provides an ideal chromophoric label by which to generate the resonance Raman (RR), ${ }^{1,2}$ or Raman difference, ${ }^{3}$ spectra of transient enzyme-substrate complexes formed in the active site during the catalytic action of serine or cysteine proteases. In the case of the serine protease group of enzymes, where the $\alpha, \beta$-unsaturated acyl group is covalently linked to an active site serine residue, RR spectroscopic studies have already provided important structural and mechanistic detail on key catalytic groups in the active site. ${ }^{2}$ These studies are presently being extended to the cysteine protease family of enzymes, via intermediates of the type $-\mathrm{CH}=\mathrm{CH}-\mathrm{C}(=\mathrm{O})-\mathrm{S}-$ enzyme and using Raman difference spectroscopy, ${ }^{4,5}$ to investigate common as well as novel mechanistic features.
A factor which limits our interpretation of the Raman spectra of $\alpha, \beta$-unsaturated thioesters, free in solution or bound to their target enzyme, is the paucity of information in the literature concerning the conformational and spectroscopic properties of these compounds. The present paper follows a study of cis- and trans-S-ethyl thiocrotonate molecules, ${ }^{6}$ as part of a series whose purpose is to fill that void. In spite of their importance in a variety of biochemical systems (e.g. see ref. 7), simple $\alpha, \beta$-unsaturated carbonyl thioesters have not been the object of detailed conformational and vibrational spectroscopic analysis. This may be due, in part, to compound instability and difficulties in purifying some members of this family. In addition, to the best of our knowledge, no systematic theoretical studies of this family of compounds have been undertaken until now. Thus, in the present study, the conformational preferences and vibrational properties of $S$-ethyl thioacrylate (abbreviated ETA) were studied by a combined $a b$ initio SCF-MO calculations/vibrational spectroscopic approach, as a step towards the general understanding of the

[^0]conformational and vibrational properties of $\alpha, \beta$-unsaturated thioesters.

## Experimental and computational methods

## Synthesis and equipment

$S$-Ethyl thioacrylate was synthesized by dropwise addition of acrylyl chloride to a dry solution of ethanethiol in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$, following the general procedure described in ref. 6. Purification was undertaken by high-performance liquid chromatography (HPLC) on a silica column (Supercosil LC-SI semiprep, $25.0 \mathrm{~cm} \times 10 \mathrm{~mm}, 5 \mu \mathrm{~m}$ pore) using hexane as eluent ( $3 \mathrm{ml} \mathrm{min}^{-1}$ ), and the purity of the sample was checked by UV and NMR spectroscopies. The Raman and FTIR spectroscopic studies were performed directly on the HPLC fractions in hexane. FTIR spectra were obtained using a Digilab FTS 40A spectrometer equipped with a DTGS (deuterated triglycine sulfate) pyroelectric detector. Data collection was performed using a demountable liquid cell equipped with KBr windows. For each spectrum, 64 scans were obtained and coadded. Raman spectra were obtained using a single monochromator, equipped with a charge-coupled device (CCD) detector and a super notch filter. ${ }^{8}$ Data collection was performed using $90^{\circ}$ sampling geometry with 300 mW 647.1 nm laser excitation. For each spectrum, 20 scans, each 10 s, were acquired and co-added.

## Computational details

The ab initio SCF-MO calculations were carried out with a $6-31 \mathrm{G}^{*}$ basis set ${ }^{9}$ using the GAUSSIAN 92 program system ${ }^{10}$ running on a DEC ALPHA 7000 computer. Molecular geometries were fully optimized by the force gradient method using Berny's algorithm. ${ }^{11}$ The largest residual internal coordinate forces were always less than $3 \times 10^{-4} E_{\mathrm{h}} a_{0}^{-1}$ or $E_{\mathrm{h}}$ $\mathrm{rad}^{-1}$, for bond stretches and angle bends, respectively. The stopping criterion for the SCF iterative process required a
density matrix convergence of less than $10^{-8}$. The force constants (symmetry internal coordinates) for use in the normal coordinate analysis were obtained from the ab initio Cartesian harmonic force constants using the program TRANSFORMER. ${ }^{12}$ This program was also used to prepare the input data for the normal coordinate analysis programs used in this study (BUILD-G and VIBRAT ${ }^{13}$ ). The calculated force fields were scaled down in order to adjust the ab initio calculated wavenumbers $\left(\mathrm{cm}^{-1}\right)$ to the experimental ones by using the linear equation $v_{\text {scaled }}=0.8996 v_{a b \text { initio }}-14.4$ previously derived for $S$-ethyl thiocrotonate. ${ }^{6}$

## Results and Discussion

## Geometries and energies

Table 1 shows the $6-31 \mathrm{G}^{*}$ energies and optimized geometries for the relevant conformers of ETA. The atom numbering is presented in Fig. 1.
ETA has three different internal degrees of freedom that can give rise to conformational isomerism: internal rotations about the $\mathrm{C}_{\alpha}-\mathrm{C}, \mathrm{C}-\mathrm{S}$ and $\mathrm{S}-\mathrm{C}$ bonds. As shown in previous studies on thioesters and thioacids, ${ }^{14-17}$ molecules having an $s$-trans conformation about the $\mathrm{O}=\mathrm{C}-\mathrm{S}-\mathrm{R}$ ( $\mathrm{R}=\mathrm{H}$ or alkyl) axis have a much higher energy than those

Table 1 6-31 $G^{*}$ calculated optimized geometries, energies and electric dipole moments for the relevant forms of ETA ${ }^{a}$

| parameter | Cg | Ca | Tg | Ta |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{C}_{\alpha}-\mathrm{C}$ | 148.97 | 148.99 | 148.92 | 148.89 |
| $\mathrm{C}=\mathrm{O}$ | 119.06 | 119.04 | 119.06 | 119.06 |
| $\mathrm{C}-\mathrm{S}$ | 177.94 | 177.82 | 178.57 | 178.45 |
| S-C | 181.99 | 182.03 | 182.04 | 182.11 |
| $\mathrm{C}-\mathrm{C}$ | 152.46 | 152.55 | 152.42 | 152.57 |
| $\mathrm{C}=\mathrm{C}$ | 132.00 | 131.98 | 131.97 | 131.96 |
| $\mathrm{C}_{\alpha}-\mathrm{H}$ | 107.51 | 107.51 | 107.54 | 107.54 |
| $\mathrm{C}_{\mathrm{B}}-\mathrm{H} 7$ | 107.46 | 107.46 | 107.47 | 107.47 |
| $\mathrm{C}_{8}-\mathrm{H} 8$ | 107.41 | 107.41 | 107.44 | 107.46 |
| $\mathrm{C}-\mathrm{H} 11$ | 108.02 | 108.04 | 107.98 | 108.01 |
| C-H12 | 108.29 | 108.04 | 108.30 | 108.01 |
| C-H13 | 108.62 | 108.48 | 108.62 | 108.49 |
| C-H14 | 108.42 | 108.44 | 108.42 | 108.44 |
| C-H15 | 108.22 | 108.44 | 108.22 | 108.44 |
| $\mathrm{C}-\mathrm{C}=\mathrm{O}$ | 123.81 | 123.90 | 120.68 | 120.77 |
| $\mathrm{O}=\mathrm{C}-\mathrm{S}$ | 122.78 | 122.45 | 121.95 | 121.57 |
| $\mathrm{C}-\mathrm{C}-\mathrm{S}$ | 113.41 | 113.65 | 117.37 | 117.66 |
| $\mathrm{C}-\mathrm{S}-\mathrm{C}$ | 101.07 | 100.53 | 100.56 | 99.98 |
| $\mathrm{S}-\mathrm{C}-\mathrm{C}$ | 113.90 | 109.92 | 114.05 | 109.71 |
| $\mathrm{C}=\mathrm{C}-\mathrm{C}$ | 120.40 | 120.42 | 126.05 | 126.05 |
| $\mathrm{H}-\mathrm{C}_{\alpha}-\mathrm{C}$ | 117.70 | 117.70 | 112.52 | 112.52 |
| $\mathrm{H} 7-{ }_{\mathrm{C}}=\mathrm{C}$ | 121.49 | 121.51 | 120.66 | 120.69 |
| $\mathrm{H} 8-\mathrm{C}=\mathrm{C}$ | 120.86 | 120.85 | 122.94 | 122.97 |
| H11-C-S | 107.45 | 108.18 | 107.61 | 108.29 |
| H12-C-S | 104.81 | 108.18 | 104.47 | 108.29 |
| H13-C-C | 109.65 | 109.47 | 109.61 | 109.46 |
| H14-C-C | 111.21 | 111.48 | 111.20 | 111.49 |
| H15-C-C | 110.84 | 111.48 | 110.90 | 111.49 |
| $\mathrm{C}-\mathrm{C}(=\mathrm{O})-\mathrm{S}$ | 179.67 | 180.00 | 179.28 | 180.00 |
| $\mathrm{O}=\mathrm{C}-\mathrm{S}-\mathrm{C}$ | 0.86 | 0.00 | 3.09 | 0.00 |
| $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ | -0.14 | 0.00 | -173.93 | 180.00 |
| $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ | 80.24 | 180.00 | 79.58 | 180.00 |
| H11-C-S-C | -43.29 | 58.31 | -44.26 | 58.39 |
| H12-C-S-C | -158.17 | -58.31 | -159.06 | -58.39 |
| H13-C-C-S | 178.47 | 180.00 | 178.20 | 180.00 |
| H14-C-C-S | 58.80 | 60.56 | 58.55 | 60.57 |
| H15-C-C-S | -62.20 | -60.56 | -62.49 | -60.57 |
| $\Delta E$ | - | 1.415 | 5.706 | 6.934 |
| $\|\mu\|$ | 1.010 | 1.021 | 1.674 | 1.733 |

${ }^{a}$ Bond lengths in pm , angles in degrees, energies in $\mathrm{kJ} \mathrm{mol}^{-1}$, dipole moments in Debye; see Fig. 1 for atom numbering. ${ }^{b}$ Energies relative to the most stable conformer; the total energy for the most stable form is $-666.3559832 E_{\mathrm{h}}$.
dien Online adopting an $s$-cis conformation about this axis $\left(\Delta E_{[s-t r a n s(C-S)]-[s-c i s(C-S)]}>13 \mathrm{~kJ} \mathrm{~mol}^{-1}\right)$ and thus were not considered in this study.

We have shown recently ${ }^{6}$ that internal rotation about the $\mathrm{C}_{\alpha}-\mathrm{C}$ bond in trans-S-ethyl thiocrotonate (trans-ETC) gives rise to two conformers: the $s$-cis and s-trans forms, with $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ dihedral angles equal to 0 and $180^{\circ}$, respectively. For the isolated molecule, the $s$-cis form was found to be $c a .7 \mathrm{~kJ} \mathrm{~mol}^{-1}$ more stable than the $s$-trans form. ${ }^{6}$ In addition, these two conformers have also been found to correspond to the stable conformations about the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis in the case of the analogous $\alpha, \beta$-unsaturated carboxylic esters and acids. ${ }^{18-21}$ This general trend can be easily correlated with the increased importance of the mesomerism associated with the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ fragment for the planar conformations, owing to a more efficient overlap of the $\mathrm{p}(\pi)$ orbitals. ${ }^{18-21}$


The factors that determine the relative energy of the two stable conformations have been pointed out elsewhere. ${ }^{18-21}$ These are, besides the mesomerism involving the two double bonds (that was found to favour the $s$-trans conformer ${ }^{18,19}$ ), mesomerism within the thiocarboxylate group,

and steric and electrostatic repulsions between the vinyl moiety and the O or S atoms (both favouring the $s$-cis form ${ }^{18,19}$ ). In the case of ETA, the s-trans conformation about the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis is found to be $c a .5 .6 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy than the $s$-cis form ( $\Delta E_{\mathrm{Tg}-\mathrm{Cg}}=5.706 \mathrm{~kJ} \mathrm{~mol}^{-1}$; $\Delta E_{\mathrm{Ta}-\mathrm{Ca}}=5.519 \mathrm{~kJ} \mathrm{~mol}^{-1}$; see Table 1). The slightly smaller $\Delta E_{\mathrm{T}-\mathrm{C}}$ values found for ETA when compared with the previously reported data for trans-ETC ${ }^{6}$ result from a delicate balance between the above mentioned factors: (i) the presence of the electron releasing $\beta$-methyl substituent in trans-ETC leads to an increased importance of the mesomerism within the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ moiety, thus giving rise to an additional stabilization of the s-trans forms relative to the most stable $s$-cis conformers (this gives rise to longer $\mathrm{C}=\mathrm{C}$ and shorter $\mathrm{C}_{\alpha}-\mathrm{C}$ bond lengths in trans-ETC ${ }^{6}$ by 0.4 and 0.5 pm , respectively, when compared with those bond lengths in ETA); (ii) since the mesomerism associated with the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ fragment competes with the mesomerism within the thiocarboxylic group, the latter effect is less important in trans-ETC than in ETA (consequently, the $\mathrm{C}-\mathrm{S}$ bond is longer in trans-ETC ${ }^{6}$


Ca
Ta

Cg

Tg

Fig. 1 Numbering of atoms for the relevant conformers of ETA
than in ETA by 0.4 pm ), and it also contributes to a slight additional stabilization in trans-ETC of the s-trans conformers; (iii) the dominating effect is, however, the electrostatic destabilizing repulsions between the more positively charged $\beta$-methyl substituent and the slightly positive thioester sulfur atom ${ }^{22,23}$ in $s$-trans trans-ETC when compared with the vinyl-S electrostatic repulsion in s-trans ETA, thus leading to the observed relative values of $\Delta E_{\mathrm{T}-\mathrm{C}}$.

The greater importance of the mesomerism involving the two double bonds in the s-trans conformers [ Tg and Ta (see Fig. 1)] when compared with the $s$-cis forms ( Cg and Ca ) gives rise to the shorter $\mathrm{C}_{\alpha}-\mathrm{C}$ bond lengths calculated for the $s$-trans forms (see Table 1). On the other hand, the $\mathrm{C}=\mathrm{C}$ bond lengths are predicted to be very similar in all conformers, indicating that the structural consequences of the mesomeric effect are, for this bond, compensated by the more important steric and electrostatic repulsions between the vinyl moiety and the S atom in the $s$-trans forms, when compared to the vinyl-O steric interaction which is operating in the s-cis conformers.
The fundamental importance of the vinyl-S and vinyl-O steric or electrostatic repulsions to the conformational preferences of the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis in ETA is reinforced by the calculated changes in the structural parameters associated with the $s$-cis $\rightarrow s$-trans isomerization (see Table 1). Thus the $\mathrm{C}-\mathrm{C}-\mathrm{S}, \mathrm{C}=\mathrm{C}-\mathrm{C}$ and $\mathrm{H} 8-\mathrm{C}=\mathrm{C}$ angles increase considerably (ca. 4.0, 5.5 and $2.0^{\circ}$, respectively), whereas the $\mathrm{C}-\mathrm{C}=\mathrm{O}$, $\mathrm{O}=\mathrm{C}-\mathrm{S}, \quad \mathrm{H}-\mathrm{C}_{\alpha}-\mathrm{C}, \quad \mathrm{H}-\mathrm{C}_{\alpha}=\mathrm{C}, \quad \mathrm{H} 7-\mathrm{C}=\mathrm{C} \quad$ and $\mathrm{H} 8-\mathrm{C}-\mathrm{H} 7$ angles decrease to compensate those changes. In turn, no significant changes are observed in the bond lengths, which usually are mainly determined by electronic effects. ${ }^{13-17}$
Besides the forms resulting from internal rotation about the $\mathrm{C}_{\alpha}-\mathrm{C}$ axis, ETA may also exist in two different-by-symmetry stable conformations differing in the relative orientation of the terminal methyl group of the $\mathrm{SCH}_{2} \mathrm{CH}_{3}$ fragment. As was previously found for trans-ETC, ${ }^{6}$ the most stable of these conformations in the studied molecule is a doubly degenerate conformational state (point group $C_{1}$ at equilibrium), corresponding to $\mathrm{C} 1-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axes of $c a . \pm 80^{\circ}$ (gauche forms). The second conformer is the $C_{\mathrm{s}}$ symmetric anti form ( $\mathrm{C} 1-\mathrm{S}-\mathrm{C}-\mathrm{C}$ dihedral angle equal to $180^{\circ}$ ), which is $c a .1 .4$ $\mathrm{kJ} \mathrm{mol}{ }^{-1}$ less stable than the gauche conformation. The fact that the preferred conformations assumed by the $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axis are not significantly affected by the conformation of the acyl fragment follows the same pattern observed for trans-ETC ${ }^{6}$ and is in consonance with our previous studies on ethyl dithioacetate and ethyl dithiopropionate, ${ }^{22,23}$ where the conformational preferences of this axis (in particular its tendency to adopt gauche conformations) have been analysed in detail. The main structural changes associated with the anti $\rightarrow$ gauche rotamerization occur in the $\mathrm{C} 10-\mathrm{C}-\mathrm{S}$, $\mathrm{H} 11-\mathrm{C}-\mathrm{S}$ and $\mathrm{H} 12-\mathrm{C}-\mathrm{S}$ angles (i.e., those angles involving the atoms directly affected by the internal rotation), which increase by $c a .4^{\circ}$ and decrease by $c a .1$ and $4^{\circ}$, respectively. Indeed, these changes follow the general trends reported in previous studies on ethyl dithioesters. ${ }^{22,23}$
The calculated electric dipole moments of the s-trans conformers of ETA are larger than those of the $s$-cis forms (Table 1). This result can be understood by considering the closer proximity of the overall positively charged $\mathrm{H}_{2} \mathrm{C}=$ to the negatively charged carbonyl oxygen in the s-cis conformers. On the other hand, the electric dipole moments associated with an anti or gauche $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axis are very similar. Thus, it can be predicted that the relative population about the s-cis/s-trans $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis will be sensitive to the polarity of the solvent, while the gauche/anti $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axis population ratio will not vary significantly upon changing the solvent polarity.
In summary, ETA exists in four different conformational states differing by the relative orientation of the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$
and $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axes: besides the most stable $s$-cis conformers about the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis [ Ca and Cg (Fig. 1)], two higher energy forms having the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis in the $s$-trans conformation also participate in the conformational equilibrium [the Ta and Tg conformers (Fig. 1)]. The structural parameters (bond lengths and, in particular, bond angles) of the different conformers studied are strongly influenced by the presence of either the strong repulsive steric and electrostatic interactions between the vinyl group and the sulfur atom (in the s-trans $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ conformers) or those involving the vinyl group and the carbonyl oxygen atom. Finally, considering the relative values of the dipole moments of the various conformers, the calculations indicate that solvent variation may be used successfully to change the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O} s-c i s / s$ trans relative populations, but does not provide a useful way to modify significantly the anti/gauche $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ conformational equilibrium.

## Vibrational studies

ETA has 39 fundamental vibrations. In the case of the symmetric ( $C_{\mathrm{s}}$ ) conformers ( $\mathrm{Ca}, \mathrm{Ta}$ ), the normal modes will span the irreducible representations $25 A^{\prime}+14 A^{\prime \prime}$, while those of non-symmetric gauche forms ( Cg and $\mathrm{Tg}, C_{1}$ point group) belong to the $A$ symmetry species. Hence, all vibrations are active in both Raman and IR.

As mentioned in the Experimental and computational methods section, calculated spectra were obtained from the $6-31 G^{*} a b$ initio wavefunctions at the equilibrium geometries


Fig. 2 FTIR spectrum ( $1100-1800 \mathrm{~cm}^{-1}$ region) of ETA (a) in hexane solution at room temperature and (b) $6-31 G^{*}$ calculated IR spectrum ( Cg and Tg scaled frequencies): $\square \mathrm{Cg} ; \square \mathrm{Tg}$. The calculated intensities of the bands due to the $s$-trans conformer are multiplied by the factor $1 / 2.25$ (see text). In this spectral region, the calculated spectra of the anti conformers ( Ca and Ta ) superimpose so closely to those of the corresponding most stable gauche forms $(\mathrm{Cg}$ and Tg ) that no bands exclusively due to these conformers are predicted to occur in the observed spectrum.
of the various conformers considered, and then scaled down in order to reproduce the experimental wavenumbers $\left(\mathrm{cm}^{-1}\right)$ by using the linear relationship $v_{\text {scaled }}=0.8996 v_{a b}$ initio -14.4 previously derived for $S$-ethyl thiocrotonate. ${ }^{6}$ Table 2 summarizes the vibrational results of this study (additional tables containing full details of the normal coordinate analysis calculations are available from the corresponding author). The experimental and calculated spectra are presented in Fig. 2-4.

The FTIR spectrum of ETA in hexane solution (Fig. 2 and 3 ) is dominated by the intense bands at 1684 and $1010 \mathrm{~cm}^{-1}$, the first band showing a shoulder in its low wavenumber wing ( $1676 \mathrm{~cm}^{-1}$ ). These two intense bands are easily assigned to the $v(\mathrm{C}=\mathrm{O})$ and $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ stretching modes of the more stable $s$-cis conformers, in agreement with the theoretical results (Table 2). It should be noted, however, that the calculations predict that the $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ band should be the most intense IR band, while in the observed spectrum this band is less intense than, for instance, that originating from the carbonyl stretching mode. Indeed, the results seem to indicate that the $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ fundamental in the $s$-cis conformers (or at least in the most stable Cg form) interacts by Fermi resonance with the combination mode involving the relatively intense $v(\mathrm{C}-\mathrm{S})$ and $\delta(\mathrm{S}-\mathrm{C}-\mathrm{C})$ vibrations, whose fundamentals occur at 727 $\mathrm{cm}^{-1}$ and ca. $300 \mathrm{~cm}^{-1}$. The second component of the expected Fermi resonance doublet is here considered to be the band appearing at $976 \mathrm{~cm}^{-1}$ [that has also a contribution to


Fig. 3 FTIR spectrum ( $400-1100 \mathrm{~cm}^{-1}$ region) of ETA (a) in hexane solution at room temperature and (b) $6-31 \mathrm{G}^{*}$ calculated IR spectrum ( Cg and Tg scaled frequencies); $\square \mathrm{Cg} ; \square \mathrm{Tg}$. The calculated intensities of the bands due to the s-trans conformer are multiplied by the factor $1 / 2.25$ (see text). In this spectral region, the anti forms give rise to observable bands at $664 \mathrm{~cm}^{-1}$ and $960 \mathrm{~cm}^{-1}$; the calculated (scaled) bands appear ( $\mathrm{Ca} ; \mathrm{Ta}$ ) at $\left(664 ; 666 \mathrm{~cm}^{-1}\right)$ and (943; 943 $\mathrm{cm}^{-1}$ ), respectively, and their positions are indicated by the two vertical arrows ( $\downarrow$ ). Note also the appreciable frequency underestimation by the calculations of the $v(\mathrm{C}-\mathrm{S})$ bands (noted by the symbol *) and the absence in the calculated spectrum of the band at $c a .976 \mathrm{~cm}^{-1}$ that corresponds to the second band of the Fermi doublet associated with the $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ mode of the Cg and Ca conformers (see text).


Fig. 4 Raman spectrum (1200-1800 $\mathrm{cm}^{-1}$ region) of ETA (a) in hexane solution at room temperature and (b) $6-31 G^{*}$ calculated Raman spectrum ( Cg and Tg scaled frequencies): $\square \mathrm{Cg} ; \square \mathrm{Tg}$. The calculated intensities of the bands due to the s-trans conformer are multiplied by the factor $1 / 2.25$ (see text). The anti forms give rise to observable band at $1243 \mathrm{~cm}^{-1}$; the corresponding calculated (scaled) band at $c a .1232 \mathrm{~cm}^{-1}$ (see Table 2) is indicated by the vertical arrow ( $\downarrow$ ).
its total intensity from the $\gamma\left(\mathrm{C}_{\alpha}-\mathrm{H}\right)$ fundamental vibration from the $s$-cis conformers; see Table 2].

Apart from the forementioned point, the general agreement between the calculated and experimental IR spectra of ETA (both frequencies and intensities) is remarkable (see Fig. 2 and 3, and Table 2). As the energy differences between the various conformers of this compound are predicted to be small, all conformers are present in significant amounts at room temperature and thus, they must contribute to the observed spectrum. Assuming a Boltzmann distribution, and taking into consideration the relative conformational energies obtained theoretically (see Table 1) as well as the fact that the gauche conformers are by symmetry doubly degenerate forms; the expected relative conformational populations (\%) at room temperature can be easily calculated, yielding $\mathrm{Cg}: \mathrm{Ca}: \mathrm{Tg}: \mathrm{Ta}=71: 20: 7: 2$. Thus, forms Cg and Ca must provide the main contributions to the bands observed. However, in general, the calculated spectra of forms differing in the relative orientation of the $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ axis do not differ very much, and thus only a few bands could be ascribed to individual conformers differing in the relative orientation of this axis. In fact, with a few exceptions, all modes which are predicted to occur at significantly different frequencies in the gauche and anti $\mathrm{C}-\mathrm{S}-\mathrm{C}-\mathrm{C}$ conformers have frequencies nearly coincident with those of different modes, and have been ascribed to bands which also have contributions from these latter vibrations [exceptions are the $v(\mathrm{C}-\mathrm{C})$ and $v(\mathrm{~S}-\mathrm{C})$

Table 2 Experimental frequencies and calculated frequencies and intensities for ETA ${ }^{a}$

| approximate description |  | calculated |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  | experimental ${ }^{\text {b }}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Cg |  |  |  | Ca |  |  |  | Tg |  |  |  | Ta |  |  |  |  |  |
|  |  | ab initio | scaled | $I_{(\mathbb{R})}$ | $I_{\text {(Raman) }}$ | ab initio | scaled | $I_{(\mathbb{R})}$ | $I_{\text {(Raman) }}$ | ab initio | scaled | $I_{(\mathbb{R})}$ | $I_{\text {(Raman) }}$ | ab initio | scaled | $I_{\text {(IR) }}$ | $I_{\text {(Raman) }}$ | IR | Raman |
| $v\left(=\mathrm{CH}_{2}\right)$ as | A | 3445 | 3085 | 3.2 | 68.5 | 3444 | 3084 | 3.2 | 69.4 | 3431 | 3072 | 10.2 | 59.3 | 3434 | 3075 | 11.4 | 66.5 |  |  |
| $v\left(\mathrm{C}_{\alpha}-\mathrm{H}\right)$ | $A^{\prime}$ | 3381 | 3027 | 3.9 | 75.1 | 3380 | 3026 | 4.3 | 76.4 | 3382 | 3028 | 4.2 | 125.9 | 3382 | 3028 | 4.1 | 133.6 |  |  |
| $v\left(=\mathrm{CH}_{2}\right) \mathrm{s}$ | $A^{\prime}$ | 3352 | 3001 | 10.4 | 72.2 | 3352 | 3001 | 10.7 | 73.3 | 3346 | 2996 | 4.6 | 82.2 | 3351 | 3000 | 4.3 | 74.4 |  |  |
| $v\left(\mathrm{CH}_{2}\right)$ as | $A^{\prime \prime}$ | 3326 | 2978 | 13.3 | 21.9 | 3329 | 2980 | 15.0 | 32.9 | 3329 | 2980 | 11.0 | 24.6 | 3332 | 2983 | 14.4 | 35.0 |  |  |
| $v\left(\mathrm{CH}_{3}\right)$ as | $A^{\prime \prime}$ | 3303 | 2957 | 15.8 | 61.8 | 3280 | 2936 | 21.7 | 80.7 | 3304 | 2958 | 17.3 | 57.5 | 3280 | 2936 | 22.3 | 75.6 |  |  |
| $v\left(\mathrm{CH}_{3}\right)$ as | $A^{\prime}$ | 3273 | 2930 | 36.1 | 95.5 | 3283 | 2939 | 43.8 | 91.7 | 3272 | 2929 | 35.5 | 92.3 | 3284 | 2940 | 43.4 | 86.2 |  |  |
| $v\left(\mathrm{CH}_{2}\right) \mathrm{s}$ | $A^{\prime}$ | 3256 | 2915 | 14.6 | 121.9 | 3269 | 2926 | 0.4 | 91.0 | 3256 | 2915 | 14.4 | 114.8 | 3270 | 2927 | 0.5 | 96.4 |  |  |
| $v\left(\mathrm{CH}_{3}\right) \mathrm{s}$ | $A^{\prime}$ | 3211 | 2874 | 38.3 | 124.2 | 3213 | 2876 | 33.0 | 120.3 | 3211 | 2874 | 37.7 | 123.6 | 3213 | 2876 | 31.1 | 117.9 |  |  |
| $v(\mathrm{C}=\mathrm{O})$ | $A^{\prime}$ | 1956 | 1745 | 200.4 | 3.9 | 1958 | 1747 | 219.8 | 3.9 | 1958 | 1747 | 296.9 | 45.7 | 1959 | 1748 | 315.0 | 46.2 | 1684, 1676* | 1685, 1677* |
| $v(\mathrm{C}=\mathrm{C})$ | $A^{\prime}$ | 1847 | 1647 | 67.5 | 104.0 | 1848 | 1648 | 65.4 | 103.8 | 1854 | 1653 | 10.6 | 40.0 | 1854 | 1653 | 11.2 | 38.5 | 1617 | 1619 |
| $\delta\left(\mathrm{CH}_{3}\right)$ as | $A^{\prime}$ | 1647 | 1467 | 3.6 | 11.8 | 1646 | 1466 | 2.0 | 10.6 | 1647 | 1467 | 3.6 | 12.1 | 1646 | 1466 | 1.9 | 9.8 | 1460 | 1460 |
| $\delta\left(\mathrm{CH}_{3}\right)$ as | $A^{\prime \prime}$ | 1636 | 1457 | 5.8 | 21.6 | 1640 | 1461 | 6.6 | 23.1 | 1636 | 1457 | 6.7 | 19.9 | 1640 | 1461 | 6.6 | 22.4 | 1460 | 1460 |
| $\delta\left(\mathrm{CH}_{2}\right)$ | $A^{\prime}$ | 1612 | 1436 | 4.7 | 18.2 | 1629 | 1451 | 7.5 | 13.8 | 1612 | 1436 | 6.1 | 16.2 | 1629 | 1451 | 7.3 | 14.0 | 1460 \#, 1436 | 1460\#, 1436 |
| $\delta\left(=\mathrm{CH}_{2}\right)$ | $A^{\prime}$ | 1572 | 1400 | 64.4 | 25.1 | 1572 | 1400 | 62.1 | 27.4 | 1582 | 1409 | 18.1 | 27.4 | 1583 | 1410 | 17.5 | 26.2 | 1410*, 1396 | 1396 |
| $\delta\left(\mathrm{CH}_{3}\right) \mathrm{s}$ | $A^{\prime}$ | 1562 | 1391 | 3.5 | 3.6 | 1567 | 1395 | 4.8 | 2.5 | 1562 | 1391 | 3.6 | 3.6 | 1567 | 1395 | 2.2 | 4.7 | 1396 | 1396 |
| $\omega\left(\mathrm{CH}_{2}\right)$ | $A^{\prime}$ | 1452 | 1292 | 27.6 | 1.3 | 1450 | 1290 | 12.4 | 10.7 | 1452 | 1292 | 22.6 | 2.3 | 1446 | 1286 | 23.1 | 1.6 | 1268 | 1289 |
| $\delta\left(\mathrm{C}_{\alpha}-\mathrm{H}\right)$ | $A^{\prime}$ | 1446 | 1286 | 1.7 | 28.3 | 1442 | 1283 | 19.9 | 21.0 | 1426 | 1268 | 3.6 | 14.5 | 1426 | 1268 | 2.3 | 14.0 | 1268 | 1289, 1273* |
| $\mathrm{tw}\left(\mathrm{CH}_{2}\right)$ | $A^{\prime \prime}$ | 1408 | 1252 | 0.2 | 8.3 | 1384 | 1231 | <0.1 | 8.6 | 1408 | 1252 | 0.2 | 6.4 | 1385 | 1232 | <0.1 | 8.1 |  | 1264, 1243 \# |
| $\omega\left(=\mathrm{CH}_{2}\right)$ | $A^{\prime}$ | 1190 | 1056 | 21.5 | 10.6 | 1186 | 1053 | 22.9 | 7.9 | 1054 | 934 | 109.3 | 1.0 | 1053 | 933 | 106.0 | 2.3 | 1063, 946* |  |
| $\gamma\left(\mathrm{CH}_{3}\right)$ | $A^{\prime}$ | 1173 | 1041 | 0.6 | 4.2 | 1169 | 1037 | 22.7 | 2.9 | 1172 | 1040 | 2.4 | 4.3 | 1174 | 1042 | 4.6 | 6.3 | 1044 |  |
| $\gamma\left(\mathrm{CH}_{3}\right)$ | $A^{\prime \prime}$ | 1169 | 1037 | 10.0 | 0.7 | 1164 | 1033 | 0.4 | 4.7 | 1178 | 1045 | 2.7 | 4.4 | 1164 | 1033 | 0.3 | 4.9 | 1044 |  |
| $\gamma\left(=\mathrm{CH}_{2}\right)$ | $A^{\prime \prime}$ | 1147 | 1017 | 31.4 | 3.4 | 1146 | 1017 | 31.4 | 3.6 | 1124 | 997 | 53.1 | 0.4 | 1123 | 996 | 55.8 | 0.3 | 1010 |  |
| $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ | $A^{\prime}$ | 1123 | 996 | 353.2 | 1.1 | 1126 | 999 | 347.5 | 2.3 | 1309 | 1163 | 173.0 | 8.1 | 1309 | 1163 | 189.4 | 8.8 | 1170*, 1010/976 ${ }^{\text {c }}$ |  |
| $\gamma\left(\mathrm{C}_{\alpha}-\mathrm{H}\right)$ | $A^{\prime \prime}$ | 1115 | 989 | 23.1 | 1.5 | 1115 | 989 | 23.1 | 1.5 | 1137 | 1008 | 7.1 | 3.4 | 1135 | 1007 | 4.3 | 3.4 | 1010*, 976 |  |
| $v(\mathrm{C}-\mathrm{C})$ | $A^{\prime}$ | 1058 | 937 | 7.3 | 7.2 | 1064 | 943 | 9.2 | 8.4 | 1059 | 938 | 2.2 | 7.0 | 1064 | 943 | 2.8 | 5.6 | 962 \#, 952 |  |
| $\gamma\left(\mathrm{CH}_{2}\right)$ | $A^{\prime \prime}$ | 834 | 736 | 7.0 | 2.8 | 851 | 751 | 0.1 | 0.1 | 831 | 733 | 3.7 | 2.1 | 854 | 754 | 0.2 | 0.9 | 740 |  |
| $\tau(\mathrm{C}=\mathrm{C})$ | $A^{\prime \prime}$ | 828 | 730 | 23.7 | 6.8 | 829 | 731 | 31.7 | 8.5 | 838 | 739 | 22.7 | 8.4 | 835 | 737 | 26.7 | 8.7 | 736 |  |
| $v(\mathrm{C}-\mathrm{S})$ | $A^{\prime}$ | 798 | 703 | 39.0 | 2.5 | 802 | 707 | 37.7 | 4.8 | 681 | 598 | 21.9 | 16.1 | 695 | 611 | 32.6 | 8.0 | 727, 631* |  |
| $\nu(\mathrm{S}-\mathrm{C})$ | $A^{\prime}$ | 727 | 640 | 3.5 | 23.6 | 754 | 664 | 2.7 | 21.4 | 631 | 643 | 15.4 | 14.8 | 756 | 666 | 3.7 | 21.5 | 664\#, 631 |  |
| $\delta(\mathrm{O}=\mathrm{C}-\mathrm{S})$ | $A^{\prime}$ | 545 | 476 | 0.8 | 3.2 | 564 | 493 | 0.2 | 2.0 | 552 | 482 | 21.8 | 1.1 | 542 | 473 | 21.8 | 2.7 | 500 |  |
| $\gamma(\mathrm{C}=\mathrm{O})$ | $A^{\prime \prime}$ | 538 | 470 | 0.5 | 4.5 | 537 | 469 | 0.5 | 4.3 | 485 | 422 | 1.0 | 5.2 | 486 | 423 | 0.4 | 5.4 | 500, 420* |  |
| $\delta\left(\mathrm{C}=\mathrm{C}_{\alpha}-\mathrm{C}\right)$ | $A^{\prime}$ | 423 | 366 | 5.9 | 3.8 | 415 | 359 | 12.0 | 3.3 | 514 | 448 | 1.8 | 3.4 | 523 | 456 | 3.0 | 2.9 | 480* |  |
| $\delta(\mathrm{S}-\mathrm{C}-\mathrm{C})$ | $A^{\prime}$ | 365 | 314 | 8.5 | 2.6 | 328 | 281 | 1.9 | 6.6 | 371 | 319 | 3.7 | 0.6 | 345 | 296 | 2.1 | 2.5 |  |  |
| $\delta(\mathrm{C}-\mathrm{C}=\mathrm{O})$ | $A^{\prime}$ | 307 | 262 | 3.8 | 0.7 | 293 | 249 | 4.5 | 0.6 | 245 | 206 | 2.0 | 3.2 | 259 | 219 | 4.4 | 4.5 |  |  |
| $\tau(\mathrm{C}-\mathrm{C})$ | $A^{\prime \prime}$ | 262 | 221 | 1.4 | 0.2 | 265 | 224 | 0.2 | <0.1 | 298 | 254 | 3.7 | 0.9 | 265 | 224 | 0.2 | 0.1 |  |  |
| $\delta(\mathrm{C}-\mathrm{S}-\mathrm{C})$ | $A^{\prime}$ | 182 | 149 | 3.1 | 1.3 | 137 | 109 | 3.3 | 1.1 | 193 | 159 | 3.4 | 0.7 | 146 | 117 | 3.6 | 1.1 |  |  |
| $\tau\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ | $A^{\prime \prime}$ | 118 | 92 | 0.3 | 5.2 | 87 | 64 | 3.6 | 2.3 | 122 | 95 | 3.1 | 2.6 | 32 | 14 | 1.4 | 2.1 |  |  |
| $\tau(\mathrm{S}-\mathrm{C})$ | $A^{\prime \prime}$ | 96 | 72 | 1.7 | 1.5 | 48 | 29 | 0.7 | 2.1 | 100 | 76 | 0.4 | 1.8 | 61 | 40 | <0.1 | 1.5 |  |  |
| $\tau(\mathrm{C}-\mathrm{S})$ | $A^{\prime \prime}$ | 69 | 48 | 3.8 | 3.3 | 125 | 98 | 2.0 | 4.5 | 32 | 14 | 1.8 | 3.7 | 131 | 103 | 4.2 | 2.8 |  |  |

[^1]stretching modes of the anti forms, ascribed to the bands at $962 \mathrm{~cm}^{-1}$ and $664 \mathrm{~cm}^{-1}$, respectively; the $\delta\left(\mathrm{CH}_{2}\right)$ bending mode of the gauche forms ( $1436 \mathrm{~cm}^{-1}$ ); and the low intensity Raman doublet of bands is $1264 \mathrm{~cm}^{-1}$ and $1243 \mathrm{~cm}^{-1}$, ascribed to the $\operatorname{tw}\left(\mathrm{CH}_{2}\right)$ vibration in gauche and anti forms, respectively]. On the other hand, the calculations predict that the s-trans forms about the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis are spectroscopically distinct from the corresponding $s$-cis conformers (Table 2). Indeed, comparison of calculated and observed IR spectra shown in Fig. 2 and 3 shows that bands at 1676, 1410, 1170 and $946 \mathrm{~cm}^{-1}$ must be respectively assigned to the $v(\mathrm{C}=\mathrm{O}), \delta\left(=\mathrm{CH}_{2}\right), v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ and $\omega\left(=\mathrm{CH}_{2}\right)$ vibrations in the $s$-trans conformers. The $s$-trans $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ stretching band at $c a$. $1170 \mathrm{~cm}^{-1}$ is of note, as it corresponds to a very intense band appearing in a 'clean' spectral region (see Fig. 3). Indeed, this result gives further support to the conclusion presented in our previous study on trans-ETC, ${ }^{6}$ where it is stressed that this band may be used as a probe of the presence of s-trans conformers of $\alpha, \beta$-unsaturated thioesters in more complex situations (for example, in enzyme active sites), since it is not expected that the $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ stretching mode will change much in frequency within this family of molecules.
Both the experimental Raman and IR spectra obtained in hexane solution are fitted well by the calculations when the calculated spectra of $s$-cis and $s$-trans conformers are added with the s-trans spectra multiplied by a factor of $1 / 2.25$ (Fig. $2-4$ ). Assuming a Boltzman population distribution, the ( $s$ -trans)-(s-cis) energy difference may be estimated to be ca. 2.0 $\mathrm{kJ} \mathrm{mol}^{-1}$. Although this estimation may have a substantial error considering that intrinsic band intensities may be significantly different in the isolated molecule situation and for the solution phases, the magnitude of the differences in energy found is large enough to enable us to conclude that the experimental data point to a considerably smaller energy difference between the two stable conformations about the $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis for ETA in solution than for the isolated molecule ( $c a .6 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ). Note that this result follows the same trend as that previously reported for trans-ETC, where the influence of the polarity of the solvent on the conformational equilibrium was analysed and it was found that, as expected, more polar solvents tend to stabilize the more polar $s$-trans conformers. ${ }^{6}$
Also, as was found for trans-ETC, the experimental frequencies of $v(\mathrm{C}=\mathrm{C})$ and $v(\mathrm{C}=\mathrm{O})$ in ETA are significantly smaller than the calculated (scaled) values, pointing to an increased polarity of these bonds (less double bond character) owing to intermolecular interactions with the solvent. As expected, the red shift in frequency is considerably more pronounced for $v(\mathrm{C}=\mathrm{O})$ than for $v(\mathrm{C}=\mathrm{C})$ and, for the same mode, for the more polar s-trans conformers (Table 2). A solvent dependent effect can also be observed for $v(\mathrm{C}-\mathrm{S})$ in both conformers, but this time leading to a blue shift. The observed increase of the frequency of this mode is also more pronounced for the more polar s-trans forms $\left[\Delta v(\mathrm{C}-\mathrm{S})_{\text {(obs-calc) }} 30 \mathrm{~cm}^{-1}\right.$ in $s$-trans forms $v s . c a .20 \mathrm{~cm}^{-1}$ in $s$-cis conformers].

Finally, it is also interesting to note that, as for other thioesters, ${ }^{22,23}$ in ETA $v(\mathrm{~S}-\mathrm{C})$ is an essentially pure mode whichever conformer is considered [however, in the gauche forms this coordinate mixes somewhat with one of the $\gamma\left(\mathrm{CH}_{3}\right)$ rocking coordinates], while both $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ and $v(\mathrm{C}-\mathrm{S})$ are considerably mixed modes. Particularly relevant is the fact that $v(\mathrm{C}-\mathrm{S})$ reduces its frequency by $c a .100 \mathrm{~cm}^{-1}$ in going from the $s$-cis to the $s$-trans $\mathrm{C}=\mathrm{C}-\mathrm{C}=\mathrm{O}$ axis conformation (i.e., s-cis $\rightarrow s$-trans isomerization about the $\mathrm{C}_{\alpha}-\mathrm{C}$ axis leads to a weakening of the $\mathrm{C}-\mathrm{S}$ bond), since this result may be important to the analysis of the $\alpha, \beta$-unsaturated thioacylenzyme vibrational data. ${ }^{3-5}$ The importance of this result is
underlined by the fact that the $\mathrm{C}-\mathrm{S}$ bond is the catalytically relevant bond that undergoes cleavage during the deacylation step of the catalytic reaction, ${ }^{3-5}$ as well as because the structural data do not reflect very clearly such a weakening of the $\mathrm{C}-\mathrm{S}$ bond length upon conformation change (as discussed before, the changes with conformation of the relevant intramolecular interactions, which determine the bond lengths in the two stable conformations, partially compensate each other, leading to a $\mathrm{C}-\mathrm{S}$ bond length in the $s$-trans form that is larger than in the $s$-cis conformer by only $c a .0 .4 \mathrm{pm}$; Table 1).
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## References

P. R. Carey and P. J. Tonge, Chem. Soc. Rev., 1990, 19, 293.
P. R. Carey and P. J. Tonge, Acc. Chem. Res., 1995, 28, 8.

3 M. J. O'Connor, R. B. Dunlap, J. D. Odom, D. Hilvert, M. Pusztai-Carey, B. C. Shenoy and P. R. Carey, J. Am. Chem. Soc., 1996, 118, 239.
4 J. D. Doran, P. J. Tonge, J. S. Mort and P. R. Carey, Biochemistry, 1996, 35, 12487.
5 J. D. Doran and P. R. Carey, Biochemistry, 1996, 35, 12495.
6 R. Fausto, P. J. Tonge and P. R. Carey, J. Chem. Soc., Faraday Trans., 1994, 90, 3491.
7 R. L. D'Ordine, P. J. Tonge, P. R. Carey and V. E. Anderson, Biochemistry, 1994, 33, 12635.
8 M. Kim, H. Owen and P. R. Carey, Appl. Spectrosc., 1993, 47, 1780.

9 W. J. Hehre, R. Ditchfield and J. A. Pople, J. Chem. Phys., 1972, 56, 2257.
10 M. J. Frisch, G. W. Trucks, M. Head-Gordon, P. M. W. Gill, M. W. Wong, J. B. Foresman, B. J. Johnson, H. B. Schlegel, M. A. Robb, E. S. Replogle, R. Gomperts, J. L. Andres, K. Raghavachari, J. S. Binkley, C. Gonzalez, R. L. Martin, D. J. Fox, D. J Defrees, J. Baker, J. J. P. Stewart and J. A. Pople, GAUSSIAN 92, revision C, Gaussian Inc., Pittsburgh, PA, 1992.
11 H. B. Schlegel, PhD Thesis, Queen's University, Kingston, Ontario, 1975.
12 M. D. G. Faria and R. Fausto, TRANSFORMER (version 1.0), Departamento de Química, Universidade de Coimbra, 1990.
13 M. D. G. Faria and R. Fausto, BUILD-G and VIBRAT, Departamento de Química, Universidade de Coimbra, 1990. (These programs incorporate several routines from programs GMAT and FPERT; H. Fuher, V. B. Kartha, K. G. Kidd, P. J. Krueger and H. H. Mantsch, Natl. Res. Counc. Can. Bull., 1976, 15, 1.)
14 R. Fausto, L. A. E. Batista de Carvalho, J. J. C. Teixeira-Dias and M. N. Ramos, J. Chem. Soc., Faraday Trans. 2, 1989, 85, 1945.

15 R. Fausto, L. A. E. Batista de Carvalho and J. J. C. TeixeiraDias, THEOCHEM, 1990, 207, 67.
16 R. Fausto, THEOCHEM, 1994, 315, 123.
17 R. Fausto, Rev. Port. Quím., 1996, 3, 59.
18 M. D. G. Faria, J. J. C. Teixeira-Dias and R. Fausto, Vibrat. Spectrosc., 1991, 1, 43.
19 M. D. G. Faria, J. J. C. Teixeira-Dias and R. Fausto, Vibrat. Spectrosc., 1991, 2, 107.
20 A. Kulbida, M. N. Ramos, M. Rasanen, J. Nieminen, O. Schrems and R. Fausto, J. Chem. Soc., Faraday Trans., 1995, 91, 1571.
21 R. Fausto, A. Kulbida and O. Schrems, J. Chem. Soc., Faraday Trans., 1995, 91, 3755.
22 R. Fausto, A. G. Martins, J. J. C. Teixeira-Dias, P. J. Tonge and P. R. Carey, J. Mol. Struct., 1994, 323, 59.

23 R. Fausto, A. G. Martins, J. J. C. Teixeira-Dias, P. J. Tonge and P. R. Carey, J. Phys. Chem., 1994, 98, 3592.

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[^1]:     equation $v_{\text {scaled }}=0.8996 v_{\text {ab initio }}-14.4$. ${ }^{6}$ When different bands are assigned to the same mode in differ
    respectively. ${ }^{c}$ The $v\left(\mathrm{C}_{\alpha}-\mathrm{C}\right)$ mode in the $s$-cis conformers gives rise to a Fermi resonance doublet (see text).

