



Search for displaced vertices of oppositely charged leptons from decays of long-lived particles in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

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ABSTRACT

A search for long-lived particles decaying into an oppositely charged lepton pair, $\mu\mu$, ee , or $e\mu$, is presented using 32.8 fb^{-1} of pp collision data collected at $\sqrt{s} = 13$ TeV by the ATLAS detector at the LHC. Candidate leptons are required to form a vertex, within the inner tracking volume of ATLAS, displaced from the primary pp interaction region. No lepton pairs with an invariant mass greater than 12 GeV are observed, consistent with the background expectations derived from data. The detection efficiencies for generic resonances with lifetimes ($c\tau$) of 100–1000 mm decaying into a dilepton pair with masses between 0.1–1.0 TeV are presented as a function of p_T and decay radius of the resonances to allow the extraction of upper limits on the cross sections for theoretical models. The result is also interpreted in a supersymmetric model in which the lightest neutralino, produced via squark–antisquark production, decays into $\ell^+ \ell'^- \nu$ ($\ell, \ell' = e, \mu$) with a finite lifetime due to the presence of R-parity violating couplings. Cross-section limits are presented for specific squark and neutralino masses. For a 700 GeV squark, neutralinos with masses of 50–500 GeV and mean proper lifetimes corresponding to $c\tau$ values between 1 mm to 6 m are excluded. For a 1.6 TeV squark, $c\tau$ values between 3 mm to 1 m are excluded for 1.3 TeV neutralinos.

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1. Introduction

Many extensions to the Standard Model (SM) predict the production of weakly-coupled, long-lived particles (LLPs). In particular, several models, including supersymmetry (SUSY) [1–6] with R-parity violation (RPV) [7,8] or with gauge-mediated supersymmetry breaking (GMSB) [9–11], Hidden Valley models [12], dark-photon models [13] or models with long-lived right-handed neutrinos [14], predict the existence of LLPs that can decay into a pair of leptons. If the LLP has a lifetime of picoseconds to nanoseconds then its decay may be observed as a displaced vertex in the inner tracking volume of the ATLAS detector at the LHC.

This letter presents a search for displaced dilepton vertices originating from decays of LLPs into an oppositely charged $\mu\mu$, ee , or $e\mu$ pair, with an invariant mass of more than 12 GeV. The analysis uses data from proton–proton (pp) collisions recorded by the ATLAS experiment in 2016 at a center-of-mass energy of 13 TeV. Two signal models are used to study the sensitivity of the ATLAS detector to such LLPs. The first is a simplified RPV SUSY model in which a squark–antisquark pair is produced, each decaying into a long-

lived neutralino which results in a pair of charged leptons and a neutrino. The second is a toy model where a LLP, denoted by Z' , is produced in $q\bar{q}$ annihilations and decays into a pair of charged leptons. These models were selected to study how the kinematic properties of a three- or two-body decay affect the signal efficiencies. The search is intended to be as model-independent as possible and is not optimized for these particular signal models.

Previously, the ATLAS Collaboration searched for displaced dilepton vertices in the inner tracking volume of the ATLAS detector at $\sqrt{s} = 8$ TeV [15] and for oppositely charged dimuons using only muon-spectrometer tracks at $\sqrt{s} = 13$ TeV [16]. ATLAS also searched for LLPs with mass of less than 2 GeV at $\sqrt{s} = 8$ TeV by considering pairs of highly collimated leptons [17]. The CMS Collaboration searched for displaced dilepton vertices in the inner tracking volume of the CMS detector at $\sqrt{s} = 8$ TeV [18] and for electrons and muons with large impact parameters at $\sqrt{s} = 8$ TeV [19].

2. ATLAS detector

The ATLAS experiment [20–22] at the LHC is a multipurpose particle detector with a cylindrical geometry and a near 4π cov-

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erage in solid angle.¹ It consists of the inner detector (ID) tracking system surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field, electromagnetic and hadronic calorimeters, and a muon spectrometer (MS).

The ID extends from a radius of about 33 to 1100 mm and to $|z|$ of about 3100 mm. It provides tracking for charged particles within the pseudorapidity region $|\eta| < 2.5$. At small radii, silicon pixel layers and stereo pairs of silicon microstrip detectors provide high-resolution position measurements. The pixel system consists of four barrel layers, and three forward discs on either side of the ATLAS detector. The barrel pixel layers, which are positioned at radii of 33.3, 50.5, 88.5, and 122.5 mm are of particular relevance to this search. The silicon microstrip tracker (SCT) comprises four double layers in the barrel and nine forward discs on either side. The radial position of the innermost (outermost) SCT barrel layer is 299 mm (514 mm). The final component of the ID, the transition-radiation tracker (TRT), is positioned at a larger radius, with coverage up to $|\eta| = 2.0$.

Lead/liquid-argon (LAr) sampling calorimeters provide electromagnetic energy measurements with high granularity. A steel/scintillator-tile hadron calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The endcap and forward regions are instrumented with LAr calorimeters for electromagnetic and hadronic energy measurements up to $|\eta| = 4.9$.

The MS surrounds the calorimeters and is immersed in a toroidal magnetic field. It provides tracking for charged particles within the pseudorapidity region $|\eta| < 2.7$ and trigger information up to $|\eta| = 2.4$. Three layers of muon detectors are arranged in concentric shells at radii between 5 m and 10 m in the barrel region, and in wheels perpendicular to the beam axis at radii between 7.4 m and 21.5 m in the endcap regions where $|\eta| > 1.05$.

A two-level trigger system [23] is used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to at most 100 kHz. This is followed by a software-based, high-level trigger that reduces the accepted event rate to about 1 kHz.

3. Dataset and simulated events

This analysis uses the pp collision data recorded in 2016, corresponding to an integrated luminosity of 32.8 fb^{-1} obtained while all parts of the detector were operational. The uncertainty in the integrated luminosity is 2.2% [24], obtained using the LUCID-2 detector [25] for the primary luminosity measurements.

The sensitivity of the ATLAS detector to LLPs decaying into a pair of leptons was studied using Monte Carlo (MC) simulations of two different signals. In the RPV SUSY simplified model, a pair of left- and right-handed squarks of the first two generations was produced via the strong interaction. These eight squarks were assumed to be mass-degenerate and decay into their SM partner and the lightest supersymmetric particle (LSP), which is a bino-like neutralino ($\tilde{\chi}_1^0$). All other SUSY particles were assumed to be decoupled. The LSP decay was mediated by the following lepton-number-violating superpotential term [8]:

$$W_{\text{LLE}} = \frac{1}{2} \lambda_{ijk} L_i L_j \tilde{E}_k, \quad (1)$$

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the center of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r_{xy}, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

where λ_{ijk} are the coupling strengths of the interactions; i, j, k denote the fermion generation; and L and \tilde{E} are SU(2)-doublet and singlet superfields, respectively, which contain the SM leptons and their superpartners. In the presence of these interactions, the LSP decays via phase space into a lepton and a virtual slepton, whose decay is described by Eq. (1), resulting in the following three-body decay:

$$\tilde{\chi}_1^0 \rightarrow \ell_k^\pm \ell_{i/j}^\mp \nu_{j/i}.$$

This model was interpreted in two different scenarios, where a single dominant λ_{121} or λ_{122} coupling was assumed. Decays via a pure λ_{121} coupling have branching fractions $\mathcal{B}(\tilde{\chi}_1^0 \rightarrow ee\nu) = \mathcal{B}(\tilde{\chi}_1^0 \rightarrow e\mu\nu) = 0.5$, while decays via a pure λ_{122} coupling have $\mathcal{B}(\tilde{\chi}_1^0 \rightarrow \mu\mu\nu) = \mathcal{B}(\tilde{\chi}_1^0 \rightarrow e\mu\nu) = 0.5$.

The MC samples of the RPV SUSY model were generated with MADGRAPH5_aMC@NLO 2.2.3 [26] interfaced to PYTHIA 8.210 [27] using the A14 set of underlying event and hadronization parameters (A14 tune) [28] and the NNPDF23LO PDF set [29]. The matrix-element calculation was performed at the tree level and included the emission of up to two additional partons. The merging of matrix elements and parton showers was done with the CKKW-L algorithm [30], with a matching scale set to one quarter of the squark mass. The cross-sections in the SUSY model were calculated to approximate, next-to-next-to-leading order in the strong coupling constant, including the resummation of soft gluon emission at next-to-next-to-leading-logarithmic accuracy (approximate NNLO+NNLL) [31–38]. The nominal cross-section and its uncertainty were derived using the PDF4LHC15_mc PDF set, following the recommendations of Ref. [39]. The samples span four hypotheses for the squark and the LSP masses and mean proper lifetimes corresponding to $c\tau = 10\text{--}1000$ mm. Instead of producing MC samples for each of the two scenarios separately, combined samples were produced in which the LSP decays into the three final states ($ee\nu$, $e\mu\nu$ and $\mu\mu\nu$) at the same rate. The events were reweighted to match the chosen λ coupling.

The events of each MC sample of the RPV SUSY model were generated for a specific mean proper lifetime τ_{MC} of the LSP. To obtain the signal efficiency for a different mean proper lifetime τ , a weight was given to each LSP which decayed at the time t :

$$w_{\text{LSP}} = \frac{\tau_{\text{MC}}}{\tau} \exp\left(\frac{t}{\tau_{\text{MC}}} - \frac{t}{\tau}\right).$$

The final prediction of the signal efficiency at a given lifetime and mass was determined by applying this reweighting procedure individually to all MC samples of different lifetimes but the same mass, and calculating the weighted mean for the efficiency.

The second signal is a toy model of a Z' boson with mass of 100 to 1000 GeV and mean proper lifetime corresponding to a $c\tau$ value of 100, 250, or 500 mm. The natural width was based on a relativistic Breit–Wigner distribution, and was varied from 10 to 140 GeV. The samples were generated with PYTHIA 8.212 using the A14 tune and the NNPDF23LO PDF set. Such a directly produced Z' is expected to be excluded by searches for displaced hadronic jets, since it would decay into $q\bar{q}$ with a high branching fraction. The purpose of this model is to derive efficiencies for a two-body decay of an LLP that can be applied to other models of similar kinematics.

In this analysis, all backgrounds were estimated from data. A selection of SM MC samples were used to test and validate the background estimation techniques for random crossings of tracks in two-track vertices and to estimate the systematic uncertainties in vertexing and tracking. These included MC samples of $t\bar{t}$ events generated using SHERPA 2.2 [40] with the NNPDF30NNLO

Table 1

Requirements on the muon, photon and electron candidates that pass the triggers used in the preselection. e^* refers to electrons that are required to pass the ‘loose’ electron identification criteria [49].

Trigger	Candidate 1	p_T [GeV]	$ \eta $	$ d_0 $ [mm]	Candidate 2	p_T [GeV]	$ \eta $	$ d_0 $ [mm]
μ	μ	> 62	< 1.07	see text				
γ	γ	> 150	< 2.5	–	γ	> 10	< 2.5	–
					e	> 10	< 2.5	> 2.0
					μ	> 10	< 2.5	see text
γ	e	> 150	< 2.5	> 2.0				
$\gamma\gamma$	γ	> 55	< 2.5	–	γ	> 55	< 2.5	–
$\gamma\gamma$	e	> 55	< 2.5	> 2.0	e	> 55	< 2.5	> 2.0
$\gamma\gamma$	e^*	> 55	< 2.5	–	γ	> 55	< 2.5	–
$\gamma\gamma$	e	> 55	< 2.5	> 2.0	γ	> 55	< 2.5	–

PDF set [41], leptonic decays of dibosons (ZZ , WW , WZ) from SHERPA 2.1 with the CT10 PDF set [42], and dijet events generated using PYTHIA 8.186 [43,44] with the A14 tune and the NNPDF23LO PDF set [29]. These MC samples contain isolated leptons (from gauge boson decays) that have a high transverse momentum p_T and leptons and displaced tracks in b -jets, in addition to tracks from pile-up vertices for the study of background and systematic uncertainties.

The ATLAS detector simulation [45] is based on GEANT 4 [46]. The generation of the simulated event samples included the effect of multiple pp interactions per bunch crossing, as well as the effect on the detector response due to interactions from bunch crossings before or after the one containing the hard interaction. Multiple pp collisions per bunch crossing were simulated with the soft QCD processes of PYTHIA 8.186 using the A2 tune [47] and the MSTW2008LO PDF set [48].

4. Reconstruction and event selection

Measurements in the ID are used to reconstruct tracks and vertices, while information from the calorimeters and the muon spectrometer is used to identify electron and muon candidates. The analysis imposes selection criteria on the two leptons but not on other collision activity, except for requiring the presence of a primary vertex (PV), a vertex with at least two tracks and the highest scalar sum of the p_T^2 of associated tracks. This ensures that the analysis is sensitive to a wide range of new physics models. Events must also satisfy the standard data-quality criteria of ATLAS and must have a PV with z location that satisfies $|z_{PV}| < 200$ mm. The latter requirement has a negligible impact on signal efficiency.

4.1. Triggers and preselection

In order to improve the sensitivity to decays of LLPs, dedicated algorithms are used for the reconstruction of tracks and vertices. As these algorithms are CPU intensive and significantly increase the storage size of a recorded collision event, they are applied only to a subset of the data events identified using a set of preselection criteria.

The preselection criteria are based on triggers that do not use ID track information, because the track reconstruction applied at the trigger stage does not reconstruct tracks from particles produced with significant displacement from the pp interaction vertex. Displaced muon candidates are selected with a muon trigger requiring a MS track that has $p_T > 60$ GeV and $|\eta| < 1.05$. Photon triggers are used to select displaced electron candidates, triggering on the electromagnetic showers produced by the electrons. These include a single-photon trigger with photon $p_T > 140$ GeV and a diphoton trigger requiring both photons to have $p_T > 50$ GeV. The trigger efficiency is 1–90%, depending on the mass and lifetime of the LLP.

Table 2

Comparison of track requirements between the standard and large radius tracking.

	Standard	Large radius
Maximum $ d_0 $ [mm]	10	300
Maximum $ z_0 $ [mm]	250	1500
Maximum $ \eta $	2.7	5
Maximum shared silicon modules	1	2
Minimum unshared silicon hits	6	5
Minimum silicon hits	7	7

In order to be preselected, events that pass the triggers are required to satisfy an additional set of criteria, which is evaluated using the data reconstructed with the standard ATLAS algorithms, including the electron [49] and muon reconstruction algorithms [50]. The preselected events are stored in a raw data format so that the reconstruction optimized for LLP decays can be performed.

The preselection criteria for muon and electron candidates are summarized in Table 1. If the muon is reconstructed using an ID track and there is a good match with an MS track based on the χ^2 per degree of freedom, $\chi^2/\text{DoF} < 5$, then the transverse impact parameter, calculated relative to the beamline, must satisfy $|d_0| > 1.5$ mm. The preselection criteria for electron candidates depend on whether a single or diphoton trigger is passed. If a photon instead of an electron satisfies the single photon trigger, an additional candidate object is required to further limit the size of the data set. The reconstructed photons are required to pass the ‘loose’ photon identification criteria [51].

4.2. Reconstruction of long-lived particle decays

The standard tracking (ST) and vertex reconstruction in ATLAS are optimized for the reconstruction of particles that originate near the pp interaction vertex. The ST requirements on the transverse (d_0) and longitudinal (z_0) impact parameters, calculated relative to the beamline and beam spot respectively, limit the reconstruction efficiency for highly displaced decays of LLPs. In order to improve the efficiency, a second track reconstruction algorithm, large radius tracking (LRT) [52], as well as a dedicated displaced-vertex (DV) reconstruction is used.

The track requirements in the ST and LRT are compared in Table 2. Of particular note is the relaxation of the requirements on d_0 , z_0 , and the maximum number of shared silicon modules and the minimum number of unshared silicon hits between two tracks. The LRT reconstructs tracks based on the remaining hits that were not used by the ST, and the resulting tracks from both track reconstruction algorithms are used in the analysis.

The tracks used in the DV reconstruction must satisfy the following requirements: $p_T > 1.0$ GeV, $2 < |d_0| < 300$ mm, and $|z_0| < 1500$ mm. In addition, each track must have at least two SCT hits and either at least two pixel hits or at least one TRT hit, with no

more than two silicon hits shared with another track. The selected ID tracks are used to find two-track seed vertices. The vertex tracks are not allowed to have hits in pixel layers at radii smaller than the seed vertex, and are required to have a hit in the nearest pixel or SCT layer at larger radius. Multi-track DVs are formed iteratively using the collection of seed vertices. In this process, DVs that are close to each other are merged and poorly associated tracks are removed from DVs to ensure that each selected track is assigned to only one DV [53].

The performance and uncertainties associated to the LRT and vertexing are studied with reconstructed K_S^0 , as described in Section 6.

4.3. Displaced-vertex selection

DV candidates are required to satisfy several selection criteria in order to select DVs with good efficiency and low fake rate. The fit of each DV is required to have $\chi^2/\text{DoF} < 5$ to reject poorly reconstructed vertices. A minimum transverse displacement, $d_{xy} = \sqrt{(x_{DV} - x_{pp})^2 + (y_{DV} - y_{pp})^2} > 2$ mm, to all vertices from proton–proton collisions (x_{pp} , y_{pp}) is required in order to suppress background from prompt decays. DVs are selected only if their transverse radius $r_{xy} = \sqrt{x_{DV}^2 + y_{DV}^2}$ is less than 300 mm and $|z_{DV}| < 300$ mm (‘fiducial volume’). The invariant mass of the tracks forming the DV must satisfy $m_{DV} > 12$ GeV to suppress backgrounds from low-mass SM particles. Each DV is required to have at least one positively charged track and at least one negatively charged track. DVs are rejected if they are within the volume directly in front of disabled pixel modules [53], which reduces the fiducial volume by 2.3%. DVs with electrons that originate within tracking layers or support structures [53] are rejected (‘material veto’) to eliminate any remaining background not vetoed by the invariant mass requirement. About 42% of the fiducial volume is discarded by the material veto.

Each DV has to be associated with at least two leptons (see below). Non-leptonic DV tracks are not rejected and have to satisfy only the track requirements applied in the DV reconstruction. A DV associated with both electrons and muons is classified as an $e\mu$ vertex, while DVs associated with electrons or muons only are classified as ee or $\mu\mu$ vertices, respectively. One of the leptons of a DV is required to satisfy the trigger and preselection criteria discussed in Section 4.1, except in the case of the diphoton trigger, where both leptons must satisfy the preselection criteria.

4.4. Lepton identification and selection

In this search, the standard ATLAS electron and muon reconstruction and identification algorithms [49,50] are slightly modified to improve the efficiency at large impact parameters. In particular, any requirement on the number of pixel hits is removed. Leptons are reconstructed using ID tracks from both the ST and the LRT. Reconstruction ambiguities between electrons and muons are removed to ensure that an ID track is associated with only one muon or electron. In this procedure, any electron sharing an ID track with a muon is removed. In addition, the object with lower p_T is removed if two electrons or two muons share an ID track.

The electron candidates are identified using a modified version of the ‘loose’ operating point of the electron identification [49], where the d_0 information is not used in the calculation of the likelihoods. Electrons are rejected if they are associated with an energy cluster in the calorimeter that does not satisfy standard quality criteria. Muons are reconstructed [50] by matching MS tracks to ID tracks and by performing a global re-fit of all measurements. For muons with $0.1 < |\eta| < 1.9$, at least 10% of the TRT hits originally

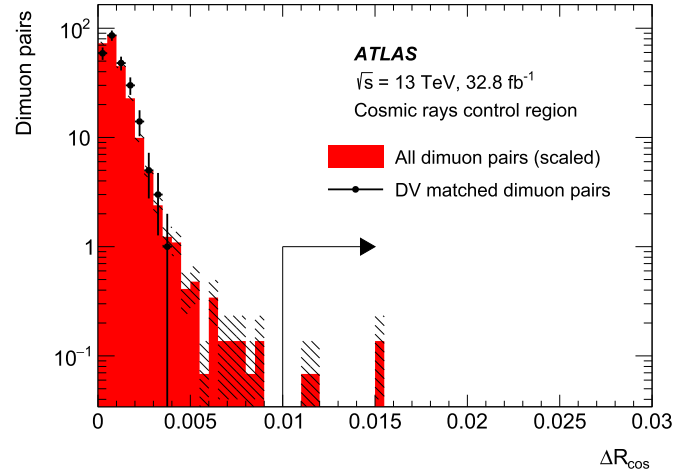


Fig. 1. ΔR_{cos} for dimuon events observed in the cosmic rays control region in collision data. The filled histogram includes all dimuon pairs, with the statistical uncertainties represented by the shaded bands. The dots indicate the subset forming a DV candidate. The distribution of all dimuon pairs is scaled to the DV distribution with $\Delta R_{\text{cos}} < 0.004$. The arrow illustrates the signal region requirement of $\Delta R_{\text{cos}} > 0.01$.

assigned to the ID track have to be included in the final fit [50]. Both the electrons and muons are required to have $p_T > 10$ GeV. In addition, electrons (muons) have to satisfy $|\eta| < 2.47$ (2.5).

A cosmic-ray muon passing through the ID during a collision event can be reconstructed as a pair of muons if one muon track segment is reconstructed in the wrong direction (i.e. opposite to its true direction of motion). The two reconstructed muons will appear to have opposite electric charges and form a high-mass DV. These back-to-back muons have $|\Delta\phi| \approx \pi$ and opposite η 's, $\eta_1 + \eta_2 \approx 0$. Cosmic-ray muons are therefore effectively rejected with a veto, $\Delta R_{\text{cos}} = \sqrt{(|\Delta\phi| - \pi)^2 + (\eta_1 + \eta_2)^2} < 0.01$. This veto is applied to any type of lepton pair in case a cosmic-ray muon is reconstructed as an electron. Only events with a single cosmic-ray muon are considered in the background estimation, as the possibility of having multiple cosmic-rays in an event was found to be negligible.

5. Background estimation

Owing to the DV mass requirement ($m_{DV} > 12$ GeV) low-mass SM processes are an insignificant source of background. The displacement requirements ($d_{xy} > 2$ mm and $|d_0| > 2$ mm) further reduce the contributions from other SM processes to negligible levels. Consequently, the largest background source is cosmic-ray muons and the random crossing of two uncorrelated leptons as discussed below.

The cosmic-ray background is estimated in a cosmic rays control region consisting of events with back-to-back muons. There are 3646 dimuons in this control region of which 246 dimuons form a displaced vertex. Fig. 1 shows the ΔR_{cos} distribution of the two samples with the former normalized to the latter using the integrals over $\Delta R_{\text{cos}} < 0.004$. The background estimate is given by the number of dimuons after normalization that satisfy the $\Delta R_{\text{cos}} > 0.01$ requirement of the signal region.

The procedure yields a background estimate of $N_{\text{bkg}}^{\text{cos}} = 0.27 \pm 0.14$ (stat.) ± 0.10 (syst.) events in the combined (ee , $e\mu$, $\mu\mu$) signal region. The systematic uncertainty is estimated by alternatively using the first three bins ($\Delta R_{\text{cos}} < 0.0015$) for the normalization and evaluating the difference in the region $0.0015 < \Delta R_{\text{cos}} < 0.004$. The signal contamination in the cosmic rays control region has been studied using the two signal models and found to be negligible.

To estimate the background contribution from randomly crossing tracks and/or leptons, two data-driven methods are used; *event mixing* and *track flipping*. The *event mixing* method randomly samples leptons from different events to measure how often a pair would form a vertex. The *track flipping* method performs secondary-vertex reconstruction on events after one track is randomly selected and flipped relative to the beam spot.²

The background estimate from the event mixing method is obtained by multiplying the number of track pairs observed in data by the probability that the two tracks will form a displaced vertex by chance. The sum of the background in the three dilepton samples is estimated to be 2.4×10^{-3} events. The track flipping method estimates a comparable background, 3.9×10^{-3} events. Both methods are tested with a MC sample of $t\bar{t}$ events using tracks without requiring lepton identification and yield about 20% more random crossing background. The final estimate uses the event mixing method because it relies on fewer assumptions in the background estimation, and the difference between these two methods, of 64%, is used as a systematic uncertainty. This yields a background estimate of $N_{\text{bkg}}^{\text{cross}} = 0.0024 \pm 0.0005$ (stat.) ± 0.0015 (syst.) events. This subdominant background contribution is neglected in final interpretations.

6. Signal systematic uncertainties

Five systematic uncertainties in the signal yields are considered. The uncertainties in the production cross-section of the SUSY model are 8.7% for a 700 GeV squark and 17.8% for a 1.6 TeV squark. The systematic uncertainty in the integrated luminosity of 32.8 fb^{-1} is 2.2%, as explained in Section 3.

The simulated events are reweighted to reproduce the distribution of the mean number of interactions per bunch crossing (μ) in the data. In this procedure, the μ measured in the data is divided by a factor of 1.09 ± 0.09 , which improves the agreement between the data and simulation for the observed number of primary vertices and reproduces the fraction of the visible cross-section of inelastic pp collisions as measured in the data [54]. The resulting uncertainty in the signal efficiency is usually less than 10%, but can be larger in cases with a small number of events in the MC simulations.

The trigger efficiencies in the simulations are scaled to those measured in the data using scale factors, which are obtained from a tag-and-probe method applied to Z +jets events. In this method, a pair of leptons is selected, where the tag lepton satisfies tight identification and isolation criteria and the probe lepton satisfies the same identification criteria as applied in the signal region. No requirements on the impact parameters of the probe and tag leptons are made.³ The requirements on the tag lepton and the invariant mass of the pair ($|m(\ell^+\ell^-) - m(Z)| < 10 \text{ GeV}$) ensure that lepton pairs from Z boson decays are selected with a low background, while the probe lepton is used to determine the trigger efficiency. The scale factors are derived with uncertainties of a few percent by comparing the trigger efficiencies measured in the data with those of a Z +jets MC sample. The average scale factor is 0.98–0.99 for the photon triggers and 0.88 for the muon trigger.

The decay $K_S^0 \rightarrow \pi^+\pi^-$ is used to study the systematic uncertainty in the LRT and vertexing, since this sample provides adequate statistics. The K_S^0 sample is selected from the same data set as used in the LLP search with the same triggers. Tracks originating

from a K_S^0 decay can be reconstructed by either the ST or the LRT algorithm. The systematic uncertainties associated with the standard track reconstruction [55] and the secondary vertexing [56] without the LRT are well understood. Systematic uncertainties due to the LRT are estimated by examining the differences in the K_S^0 yield between data and MC simulation as a function of vertex radius.

Events have to satisfy the event selection criteria described above except that the preselection is not explicitly required, as the high- p_T photon or muon triggers would reduce the statistical precision of this study significantly. Each K_S^0 candidate vertex must have exactly two tracks that are oppositely charged and an invariant mass between 350 and 650 MeV. They have to satisfy the standard vertex selection criteria described in Section 4.3. The difference of the longitudinal impact parameters between the two tracks, each relative to the K_S^0 vertex, is required to be less than 2 mm.

The numbers of K_S^0 vertices found in the data and background MC samples are binned in transverse vertex radius, r_{xy} . The K_S^0 yields in each bin are estimated after subtracting the background contributions using sidebands in the invariant mass distribution. The vertex yields of K_S^0 obtained with LRT are compared between the data and the background MC sample, where the data are normalized by adjusting the ST K_S^0 yield in data to that of the MC simulation. The largest difference in the r_{xy} distribution between data and MC simulation is chosen as the systematic uncertainty, and the statistical uncertainty in the difference is then included to yield a total systematic uncertainty of 6%.

In addition to the systematic uncertainty of 6% from LRT and vertexing, there is a systematic uncertainty of 2% from ST [55] and 1% from varying the parameters of the secondary vertexing [56]. The total uncertainty from tracking and vertexing is 10%, obtained by adding all the contributions linearly.

The uncertainty on the electron or muon identification efficiency can be neglected in comparison to the uncertainty in the LRT and vertexing.

7. Results and interpretations

No events are observed satisfying the selection criteria of the signal region, consistent with the estimated background of 0.27 ± 0.17 events.

Overall signal efficiencies for the Z' toy model are shown in Fig. 2. The modified reconstruction techniques used here significantly improve the signal sensitivity out to large decay radii, with decreasing sensitivity at larger decay radius. The signal efficiency increases with the heavier LLP because the leptons are produced with higher p_T , resulting in higher trigger efficiencies. This is particularly evident for the Z' with mass of 100 GeV in which the efficiency is significantly lower.

Overall, the LRT and DV techniques used in this study preserve detection efficiencies at large decay radii, showing sensitivity for LLPs with decay lengths of the order of the ID dimensions. The HepData entry [57] associated with this publication has additional efficiency maps and material to reinterpret the results of this search in other models.

For the interpretations of the RPV SUSY model, two independent scenarios are considered, each corresponding to an LSP decay that is mediated by a single dominant RPV coupling, λ_{121} or λ_{122} . Fig. 3 shows the overall signal efficiencies for each scenario, which can be as high as 40% for a 1.3 TeV neutralino and 12% for a 50 GeV neutralino. The selection criteria that have the largest impact on the signal efficiencies are the reconstruction efficiency of displaced vertices, the trigger and preselection requirements, and, in the case of ee and $e\mu$ vertices, the material veto.

² In the track flipping, the track parameters are transformed as: $d_0 \rightarrow -d_0$, $z_0 \rightarrow -z_0$, $\theta \rightarrow \pi - \theta$, $\phi \rightarrow \phi - \pi$ for $\phi > 0$ and $\phi \rightarrow \phi + \pi$ for $\phi \leq 0$.

³ Studies in signal MC samples have shown that a dependence of the trigger efficiencies on the impact parameters of the leptons can be neglected.

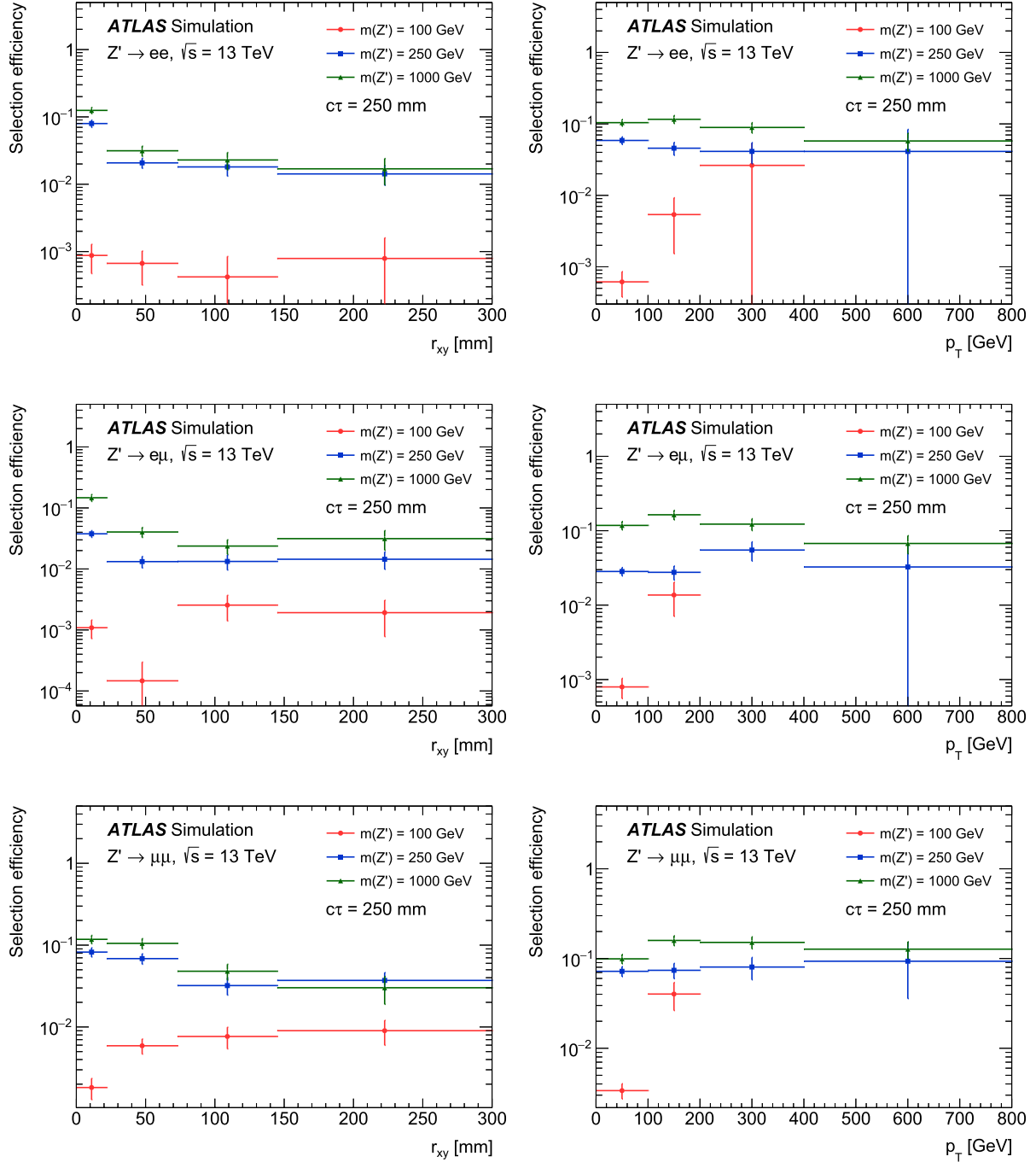


Fig. 2. Dependence of the overall signal efficiency on (left) the r_{xy} and (right) the p_T of the long-lived Z' for three masses, $c\tau = 250$ mm and (top) $Z' \rightarrow ee$, (middle) $Z' \rightarrow e\mu$ and (bottom) $Z' \rightarrow \mu\mu$. The error bars indicate the total uncertainties.

Upper limits at the 95% confidence level (CL) are calculated with the *HistFitter* [58] package by testing various signal hypotheses using the CL_s prescription [59] with a profile likelihood ratio as the test statistic, whose probability density function is determined with pseudo experiments.

The model-independent upper limit on the visible cross-section, which is defined as the production cross-section times the overall signal efficiency, $\langle \epsilon \sigma \rangle_{\text{obs}}^{95}$ is 0.09 fb and on the visible number of signal events S_{obs}^{95} is 3.0 events, consistent with the expectation of $3.0^{+0.6}_{-0.0}$ events.

The upper limits on the production cross-section of the RPV SUSY model are shown in Figs. 4 and 5 for a 700 GeV squark and

a 1.6 TeV squark, respectively. The observed limits are consistent with the expected limits within the uncertainties. Mean proper neutralino lifetimes which correspond to $c\tau$ between 1 mm and 6 m are excluded for a 700 GeV squark. If the squark mass is much higher, 1.6 TeV, the excluded $c\tau$ region is 3 mm to 1 m for a 1300 GeV neutralino, while the $c\tau$ of a 50 GeV neutralino can be constrained only in the λ_{122} scenario, where $c\tau$ values between 4 and 30 mm are excluded.

8. Conclusion

A search for a long-lived particle with a mass of more than 12 GeV that decays into an oppositely charged ee , $\mu\mu$ or $e\mu$ pair

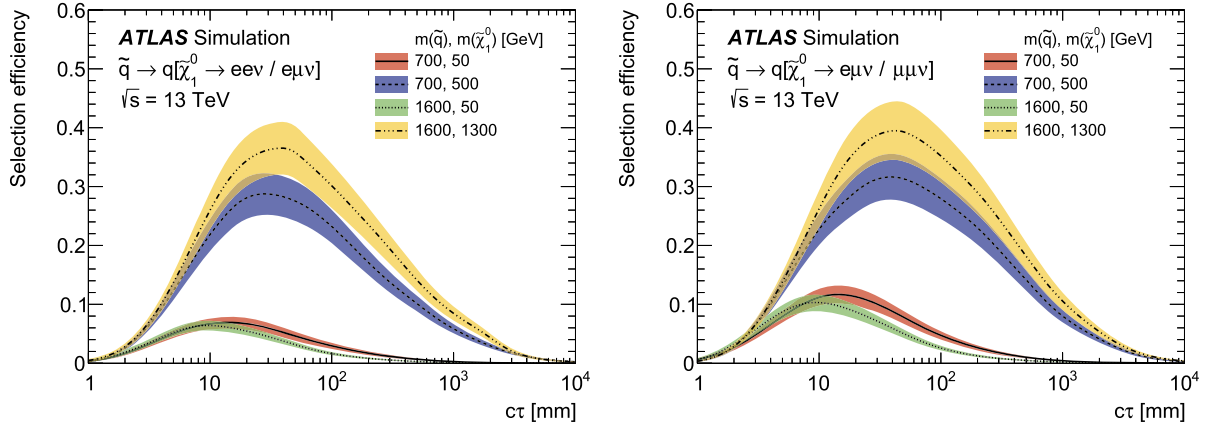


Fig. 3. Overall signal efficiency as a function of the mean proper lifetime of the $\tilde{\chi}_1^0$, in units of $c\tau$, for (left) the λ_{121} and (right) the λ_{122} scenarios with various combinations of squark and $\tilde{\chi}_1^0$ masses. The shaded bands indicate the total uncertainties.

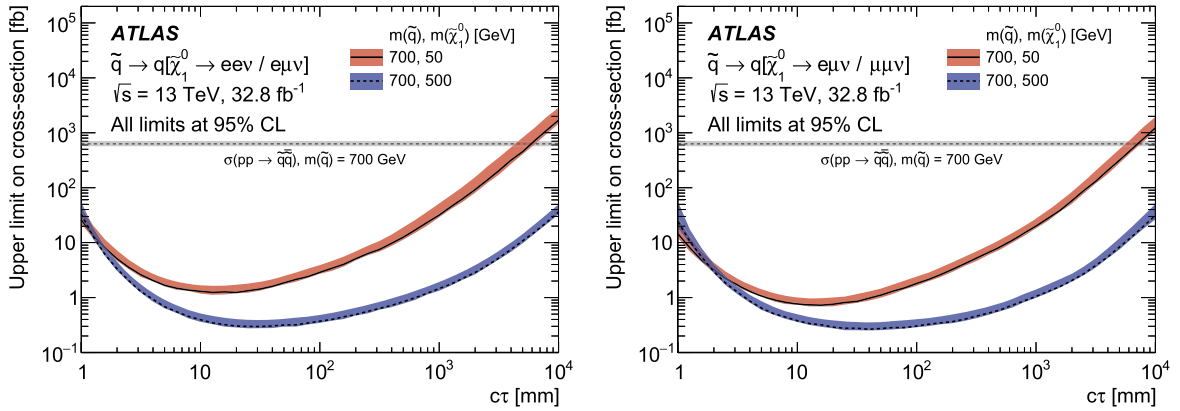


Fig. 4. Observed upper limits at 95% CL on the squark–antisquark production cross-section as a function of the mean proper lifetime of the $\tilde{\chi}_1^0$, in units of $c\tau$, for (left) the λ_{121} and (right) the λ_{122} scenarios and a 700 GeV squark. The horizontal lines indicate the theoretical (approx. NNLO+NNLL) cross-sections with the uncertainties shown as shaded bands. The shaded bands around the observed limits indicate the $\pm 1\sigma$ variations in the expected limit.

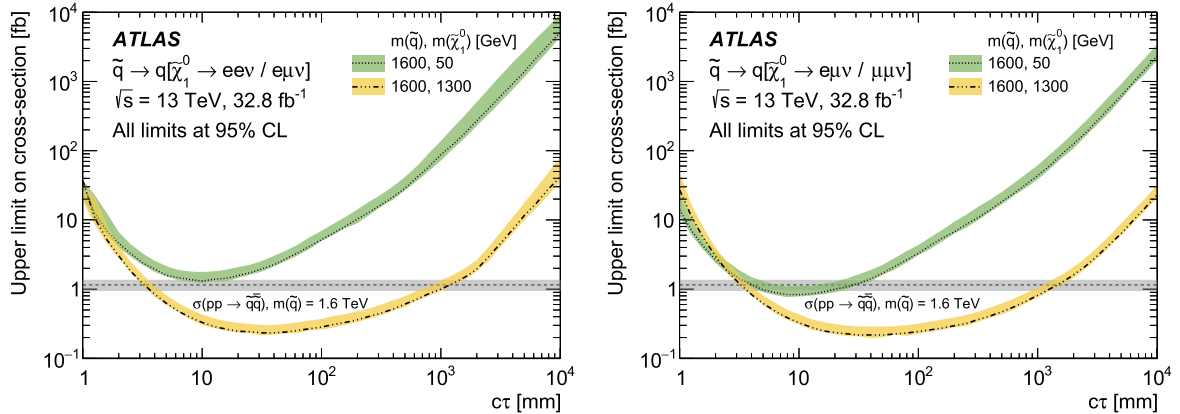


Fig. 5. Observed upper limits at 95% CL on the squark–antisquark production cross-section as a function of the mean proper lifetime of the $\tilde{\chi}_1^0$, in units of $c\tau$, for (left) the λ_{121} and (right) the λ_{122} scenarios and a 1.6 TeV squark. The horizontal lines indicate the theoretical (approx. NNLO+NNLL) cross-sections with the uncertainties shown as shaded bands. The shaded bands around the observed limits indicate the $\pm 1\sigma$ variations in the expected limit.

within the ATLAS inner detector has been performed on 32.8 fb^{-1} of pp collision data at $\sqrt{s} = 13 \text{ TeV}$ collected in 2016. No events are observed in the signal region, consistent with the expected background of 0.27 ± 0.17 events. The detection efficiencies for generic resonances with lifetimes ($c\tau$) of 100–1000 mm decaying into a dilepton pair with masses between 0.1–1.0 TeV are presented as a function of p_T and decay radius of the resonances to

allow the extraction of upper limits on the cross sections for theoretical models. Exclusion limits are derived for a supersymmetric signal model in which long-lived neutralinos ($\tilde{\chi}_1^0$) are produced through squark–antisquark production and decay into two charged leptons and one neutrino via R-parity violating couplings. If the eight left- and right-handed squarks of the first two generations have a common mass of 700 GeV, mean proper lifetimes for long-

lived neutralinos which correspond to $c\tau$ values between 1 mm and 6 m are excluded. If the squark mass is 1.6 TeV, $c\tau$ values between 3 mm and 1 m are excluded for a 1.3 TeV neutralino, while for a 50 GeV neutralino the limits are significantly weaker.

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