Alternating Binaphthyl—Thiophene Copolymers: Synthesis, Spectroscopy, and Photophysics and Their Relevance to the Question of Energy Migration versus Conformational Relaxation


Department of Chemistry, University of Coimbra, Rua Larga, 3004-535 Coimbra, Portugal, Departamento de Engenharia Química e Biológica, Centro de Química Estrutural, Instituto Superior Técnico, Av. Rovisco Pais, s/n, Lisboa, Portugal, and Makromolekulare Chemie, Bergische Universität Wuppertal, Gaussstr. 20, 42097 Wuppertal, Germany

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ABSTRACT: The synthesis and a comprehensive spectroscopic and photophysical study are presented of four alternating binaphthyl—oligothiophene copolymers (DP: 10–15 repeat units) in solution at room and low temperature and in the solid state (thin films). Detailed results are presented on absorption, emission, and triplet—triplet absorption spectra together with all relevant quantum yields (fluorescence, intersystem crossing, internal conversion, and singlet oxygen formation), excited-state lifetimes, and singlet and triplet energies. From these, several conclusions can be drawn. First, the main deactivation channels for the molecules in solution are the radiationless processes (S→S internal conversion and S→T intersystem crossing). Second, in the solid state the fluorescence quantum yields are smaller than those in solution. From time-resolved fluorescence decays in the picosecond time domain, three decay components are seen: a fast decay (40–60 ps) at short wavelengths, which becomes a rising component at longer wavelengths, an intermediate decay component (330–477 ps) associated with an ensemble of isolated segment-like units, which is dominant at the initial part of the emissive spectra and progressively decreases for longer emissions, and a third exponential related to the emission of the fully relaxed polymer. Together with steady-state anisotropy studies, this is discussed in terms of the possibilities of energy migration/transfer along the polymer chain and of the conformational (torsional) relaxation of the systems studied.

Introduction

Conjugated organic polymers are playing a major role within the developing field of plastic electronics, due in particular to their applications in light-emitting diodes, thin film transistors, photovoltaic devices, sensors, etc. Considerable research effort has been devoted to the improvement of the light-emitting properties and color tuning of photoluminescence devices involving conjugated polymers. One approach for achieving this goal involves copolymerization with chromophores displaying high fluorescence emission yields and good electronic properties.

Pure binaphthyl enantiomers are commercially available. Proper substitution of the binaphthyl core has provided suitable monomers with different side chains in 2,2’-position for our copolymerization experiments. In this study, the synthesis of fluorescent alternating copolymers containing both binaphthyl and oligothiophene moieties will be described. Oligothiophenes have been chosen as building blocks because of their unique properties, such as high (photo)stability and good hole transport properties. However, oligothiophenes, and more particularly polythiophenes, display low fluorescence quantum yields due to the high density of vibrational modes of the excited-state species that can lead to other modes of decay than photon emission (such as thermal decay and intersystem crossing), in addition to their efficient intersystem crossing induced by the heavy atom effect of the sulfur atom. As indicated above, it may be possible to overcome the low fluorescence quantum yields and improve the emission efficiency of oligothiophenes by copolymerization with highly fluorescent moieties. This approach can decrease the vibrational freedom of the singlet excited state and favor the excited-state radiative channel. In addition, the incorporation of nonplanar units, such as binaphthyl, in the polymer backbone should suppress interchain interactions, which could prevent polymer chains from packing, thereby reducing self-quenching processes and thus improving the fluorescence quantum efficiency. In the case of binaphthyl, the pseudo-orthogonal nature of this chromophore (with the dihedral angle between the two naphthalene units ranging from 60° to 120°) should suppress the crystallization and facilitate the formation of an amorphous glassy state.

One important current topic of discussion in the field of conjugated organic polymers is the exact nature of fast relaxation processes occurring following excitation of the polymer. The possibilities which have been suggested include intrachain and interchain energy transfer, together with more localized phenomena such as intramolecular vibrational relaxation and/or conformational (torsional) relaxation within the polymer skeleton. In general, conjugated polymers typically contain between 10 and 1000 repeat units, and consequently intrachain processes may involve different conjugation segments of the polymer, which may result from conformational defects, such that they are likely to be important with isolated chains, whereas interchain processes, which are particularly relevant in the solid state, involve different polymer units/chains. These processes may be interdependent. Intrachain processes depend on the size of the polymer backbone chain and consequently on the number of conjugated segments (chromophores) along the chain. The existence of conformational defects is often responsible for the (de)localization of π-electrons onto individual segments. Depending on the different chromophore structures, each segment can emit within the same polymer.
Two distinct explanations have been given to the fast component in the fluorescence decays of organic conjugated polymers; either it results from resonance energy transfer between different conjugated segments, or it is due to conformational relaxation within the polymer backbone. A differentiation between the relative contributions of these two competitive processes to the fast component decay can be made by investigating oligomeric compounds where energy transfer is absent. In general, it is accepted that energy transfer can occur between different chromophores within the polymer chain, and these processes can be faster or slower, ranging from values <1 ps to some hundreds of ps, depending on the relative energy of these segmental units.

In the present study, we have investigated four polymers possessing a relatively small number of binaphthyl—oligothiophene repeat units (10–15). For this reason, and due to the function of the binaphthyl moieties as effective conjugation barriers, the extent of intrachain energy transfer between the conjugated segments should be drastically reduced. In this particular case, the shortest living component of the fluorescence decay should preferentially be associated with conformational relaxation processes. Thus, the present investigation of the photophysical and time-resolved fluorescence decay profiles in solution of four alternating binaphthyl—oligothiophene copolymers may provide new insights into the fluorescence relaxation phenomena in conjugated polymers. In addition, thiophene-containing copolymers are now important materials for solar cell applications and charge transport applications, and understanding of the photophysical and energy migration properties is also relevant for these.

Experimental Section

Materials. Unless otherwise indicated, all starting materials were purchased from commercial sources (Aldrich, Fischer, EM Science, Lancaster, ABCR, Strem) and used without further purification. All reactions were assembled under an inert atmosphere either in a screw-capped vial or in a resealable Schlenk tube. Analytical thin-layer chromatography (TLC) was performed on Merck silica gel F254 plates with silica gel (230 plates with F-254 indicator. Visualization was accomplished by an ultraviolet lamp (254 nm). Silica gel column chromatography was carried out with silica gel (230–400 mesh) from EM Science. 1H and 13C NMR spectra were recorded on a Bruker ARX 400 spectrometer. Chemical shifts (δ) are reported in parts per million (ppm) with respect to residual solvents protons as internal standards. The coupling constants are reported in hertz (Hz). Splitting patterns are designated as s (singlet), d (doublet), t (triplet), q (quartet), bs (broad singlet), m (multiplet), and bm (broad multiplet). Low-resolution mass spectrometry was obtained on a Varian MAT 311A operating at 70 eV (electron impact, EI). Mass spectra are reported as relative intensity. Molecular weight determinations were performed using a Spectra 100 GPC (vs polystyrene standards in THF): Mn = 4930, Mw/Mn = 1.7. 1H NMR (400 MHz, CDCl3, 80 °C): δ = 8.11–7.71 (m, 4H), 7.66–6.94 (m, 10H), 4.02–3.57 (m, 4H), 1.68–0.55 (m, 30H). 13C NMR (100 MHz, CDCl3, 80 °C): 155.3, 143.9, 136.8, 134.0, 129.7, 129.6, 129.4, 129.2, 126.6, 126.0, 124.8, 124.5, 124.4, 124.0, 121.1, 120.6, 117.9, 99.8, 70.0, 32.0, 29.7, 29.5, 26.0, 23.0, 14.5 ppm. 13C NMR (100 MHz, CDCl3, 80 °C): 155.3, 143.9, 136.8, 134.0, 129.7, 129.6, 129.4, 129.2, 126.6, 126.0, 124.8, 124.5, 124.4, 124.0, 121.1, 120.6, 117.9, 99.8, 70.0, 32.0, 29.7, 29.5, 26.0, 23.0, 14.5 ppm.

Thiophene Copolymers

Poly(6,6′-dibromo-2,2′-dioxycyclobenzoyl)-alt-2.5-bithiophene (BNpα2). The general procedure was followed using 6,6′-dibromo-2,2′-bis(octyloxy)-1,1′-binaphthyl (80 mg, 0.120 mmol), 5,5′-bis(trimethylstannyl)-2,2′-bithiophene (59 mg, 0.120 mmol), PdCl2dppf (9 mg, 0.012 mmol), and KF (70 mg, 1.197 mmol). Yield: 17%. 1H NMR (400 MHz, CDCl3, 80 °C): δ = 7.89–7.34 (m, 4H), 6.75 (d, 8.0 Hz, 4H), 4.91 (s, 4H), 2.42 (m, 4H), 1.45 (m, 4H), 1.19 (bm, 23H), 0.80 (m, 3H) ppm. 13C NMR (100 MHz, CDCl3, 80 °C): 151.4, 142.7, 134.4, 132.8, 130.7, 130.1, 129.9, 128.9, 128.6, 127.5, 127.0, 120.6, 117.8, 117.5, 71.6, 35.9, 32.2, 31.7, 29.7, 29.5, 23.0, 19.5, 14.5 ppm.

Poly(6,6′-dibromo-2,2′-dioxycyclobenzoyl)-alt-2.5-bithiophene (BNpα2). Following the general procedure (R)-6,6′-dibromo-2,2′-bis(octyloxy)-1,1′-binaphthyl (60 mg, 0.090 mmol), 5,5′-bis(trimethylstannyl)-2,2′-5,5′-2′-terthiophene (52 mg, 0.090 mmol), PdCl2dppf (7 mg, 0.009 mmol), and KF (52 mg, 0.898 mmol) were reacted. Yield: 62%. 1H NMR (400 MHz, CDCl3, 80 °C): δ = 7.85 (d, 9.0 Hz, 2H), 7.78 (d, 8.1 Hz, 2H), 7.34 (d, 9.0 Hz, 2H), 7.24 (m, 2H), 7.12 (t, J = 7.4, 7.4 Hz, 2H), 7.07 (d, J = 8.1 Hz, 2H), 3.85 (bm, 4H), 1.39–1.26 (bm, 4H), 1.22–1.09 (bm, 4H), 1.09–0.71 (bm, 22H) ppm. 13C NMR (100 MHz, CDCl3, 25 °C): 154.9, 134.4, 129.5, 129.3, 128.1, 126.2, 125.9, 123.7, 120.9, 116.4, 70.2, 32.0, 29.7, 29.43, 29.42, 25.9, 23.0, 14.5 ppm.

Poly[6,6′-dibromo-2,2′-dioxycyclobenzoyl]-alt-2,5-bithiophene (BNpα2). Following the general procedure (R)-6,6′-dibromo-2,2′-bis(4-octyloxybenzoyl)-1,1′-binaphthyl (2) (60 mg, 0.071 mmol), 5,5′-bis(trimethylstannyl)-2,2′-5,5′-2′-terthiophene (59 mg, 0.071 mmol), PdCl2dppf (5 mg, 0.007 mmol), and KF (70 mg, 0.707 mmol) were reacted. Yield: 56%. 1H NMR (400 MHz, CD2Cl2, 80 °C): δ = 7.85 (d, 9.0 Hz, 2H), 7.78 (d, 8.1 Hz, 2H), 7.34 (d, 9.0 Hz, 2H), 7.24 (m, 2H), 7.12 (t, J = 7.4, 7.4 Hz, 2H), 7.07 (d, J = 8.1 Hz, 2H), 3.85 (bm, 4H), 1.39–1.26 (bm, 4H), 1.22–1.09 (bm, 4H), 1.09–0.71 (bm, 22H) ppm. 13C NMR (100 MHz, CDCl3, 25 °C): 154.9, 134.4, 129.5, 129.3, 128.1, 126.2, 125.9, 123.7, 120.9, 116.4, 70.2, 32.0, 29.7, 29.43, 29.42, 25.9, 23.0, 14.5 ppm.
where measured with excitation at 395 nm and the horizontally polarized output beam from the GWU (second harmonic) was first passed through a Thorlabs depolarizer (WDPOL-A) and after by a Glan-Thompson polarizer (Newport 10GT04) with vertical polarization. Emission at 90° geometry collected at magic angle polarization was detected through a double subtractive Oriel Cornerstone 260 monochromator by a Hamamatsu microchannel plate photomultiplier (R3809U-50). Signal acquisition and data processing was performed employing a Becker & Hickl SPC-630 TCSPC module. Fluorescence decays and the instrumental response function (IRF) where collected using 4096 channels in a 0.814 ps/channel scale, until 5 × 10^13 counts at maximum were reached. The full width at half-maximum (fwhm) of the IRF was about 22 ps and was highly reproducible with identical system parameters.

The experimental setup (Applied Photophysics laser flash photolysis apparatus pumped by a Spectra-Physics Nd:YAG laser) used to obtain triplet spectra and triplet yields has been described elsewhere.29,32 Triplet state absorption spectra were also characterized by pulse radiolysis (using the Free Radical Research Facility, Daresbury, UK) as has been described in detail elsewhere.33 Spectra are comparable with those obtained by laser flash photolysis, except at longer wavelengths (>700 nm), where those obtained by flash photolysis appear to be attenuated by the diminished response of the detector. First-order kinetics were observed in all cases for the decay of the lowest triplet state, and lifetime values were in the microsecond range.

Special care was taken in determining triplet yields to have optically matched dilute solutions (abs ≈ 0.2 in a 10 mm square cell) and low laser energy (≤ 2 mJ) to avoid multiphoton and T–T annihilation effects.

The triplet molar absorption coefficients were obtained by the singlet depletion and energy transfer methods. Details on the experimental procedures and data analysis used can be found in refs 29 and 32.

Room-temperature singlet oxygen phosphorescence was detected at 1270 nm with equipment and procedures elsewhere reported.29,32 From these signals the singlet oxygen quantum yields were obtained. Thin films from the compounds were obtained with a desktop precision spin-coating system, model P6700 series from Speedline Technologies. Solid-state thin film from the samples was obtained by deposition of a few drops from a solution of the compounds into a circular sapphire substrate (10 mm diameter) followed by spin-coating (2500 rpm) in a nitrogen-saturated atmosphere (2 psi). The solutions for spin-coating were prepared by adding 2 mg of the samples to 15 mg of Zeonex in 200 mL of toluene solution with stirring at 40 °C for 30 min. The fluorescence emission spectra of the thin films were obtained with a Horiba-Jobin-Yvon integrating sphere. The solid-state photoluminescence quantum yields in thin films were obtained as previously described.29,32

Molecular orbital calculations were carried out with Hyperchem 7.0 program, using AM1 for geometry optimization and AM1-CI, PM3-CI, and ZINDO/S-CI (129 CIs) for the calculation of the electronic states.33

Scheme 3 shows examples of optimized conformations of BNpt2 and BNpt3. In all cases, the dihedral angle between the naphthyl groups was close to 90°, which, if correct, implies that the biphenyl units act as efficient conjugation barriers. This gains particular relevance for the investigation of the relative contribution of conformational relaxation vs energy migration (transfer) within the copolymers chains.

### Results and Discussion

**Synthesis of the Polymers.** The synthesis of the oligothiophene monomers has been described elsewhere.34 The chemical structures of the desired biphenyl–oligothiophene copolymers and the synthetic protocols are schematically depicted in Scheme 1. The biphenyl monomers were prepared from the respective diols by alkylation with 1-bromooctane or 1-(bromomethyl)-4-octylbenzene using anhydrous potassium carbonate as base in acetonitrile to give 6,6′-dibromo-2,2′-bis(octyloxy)-1,1′-biphenyl-
Scheme 1. Preparative Procedure toward Binaphthyl–Oligothiophene Copolymers X; n = 2, 3

Scheme 2. Structures and Acronyms of the Binaphthyl–Oligothiophene Copolymers Investigated

Absorption and Fluorescence. The absorption spectra for the copolymers at room temperature (293 K), low temperature (77 K, given by the fluorescence excitation spectra), and in the solid state (thin films) are broad and devoid of vibrational structure (Figure 1). It is worth noting that the absorption band due to the (1,1′-bi)naphthyl chromophore, BNp (Figure 2), is not observed in the copolymers, indicating that part of this moiety is in conjugation with the oligothiophene units of the backbone.

A red shift of ca. ~20 nm in the absorption maximum was observed with the increase of size of the oligothiophene group (see Table 2). The calculated values for the lowest electronic transitions of the parent compounds of BNpa2 and BNpa3 (Scheme 3) are 457 nm (AM1) and 512 nm (ZINDO/S), in agreement with the experimental values (see Table 2), which gives support to the calculated molecular geometries.

Upon going from the spectra in solution at 293 K to those in glasses at 77 K, the absorption bands were red-shifted (with a minimum of ~8 nm in BNpa3 and a maximum of 17 nm in OPhBNpa3; see Table 2). Analogous behavior was observed at 293 K on going from solution to the solid-state films.

Figure 1 also presents the fluorescence emission spectra for the copolymers at 293 and 77 K and in thin films. A small red shift (1–3 nm) of the 0–0 vibronic is observed upon cooling from RT to LT (see Table 2).

In contrast to the absorption spectra, the RT fluorescence emission spectra show vibrational structure. Again, there is no evidence for (1,1′-bi)naphthyl-related emission bands, supporting the idea that π-electron delocalization involves both the (bi)-
naphthyl and the oligothiophene moieties (Figure 1). Moreover, and in contrast with the absence of any significant shift of the 0→0 vibronic, there is a significant change in the progression of the remaining vibronic bands at 77 K relative to 293 K. In view of the Franck–Condon forbidden nature of the first absorption band of these compounds and of the highly resolved nature of the emission band (at both 293 and 77 K) there is clearly a change in the potential energy curves related to the ground and first excited singlet states. In the case of the oligothiophenes20 (and naphthalene–oligothiophenes21), this has been attributed to the existence of an ensemble of different conformations in the ground state and to a planar quinoidal-like structure22 in S1. This more rigid, quinoidal-like structure in thiophene-like oligomers and polymers promotes a good resolution of the vibronic energetic levels, which is consistent with the observed well-resolved fluorescence spectra observed under all conditions. While detailed assignment of the vibronic modes involved is outside the scope of this study, it is clear by comparison of the fluorescence spectra of the binaphthyl–oligothiophenes copolymers in Figure 1 with those of the model thiophene oligomers (bithiophene and terthiophene) in Figure 2, that for the terthiophene copolymer derivatives (Bnπ3 and OPhBNπ3) the emission spectra present a similar (although red-shifted) vibronic progression to the analogous terthiophenes copolymers in Figure 1 with those of the model polymer building block of the copolymers (2,2’-octyloxy-1,1’-binaphthyl-BNP; Figure 2, structure as inset) in methylcyclohexane solution display much shorter absorption and emission wavelength regions of the emission spectra, depicted in Figure 3 for BNπ2 and OPhBNπ2.

Moreover the fluorescence quantum yield (φf) is also much greater than any of the values found for the polymers studied (see Table 3).

In the solid state (thin films) a loss of vibronic structure is seen with the concomitant broadening and red shift of the emission spectra. These results provide strong evidence for aggregation in the solid state. The broadening and loss of structure of the emission spectra in the solid state are generally attributed to intermolecular interactions giving rise to different types of packing between the polymer chains and different degrees of twisting in the polymer/oligomer backbone, leading to broad and unstructured emission bands. Moreover, the evidence for aggregates in thin films comes from the absorption spectra which show an increased absorbance at the longer wavelength region as compared to those in their corresponding solution states at 293 K (see Figure 1).

It is also interesting to observe the fluorescence excitation spectra obtained at 293 and 77 K and collected in different regions of the emission spectra, depicted in Figure 3 for BNπ2.

Table 2. Spectroscopic Parameters for the Binaphthyl–Oligothiophenes Copolymers in Toluene at Room Temperature (293 K), Low Temperature (77 K), and in Thin Films

<table>
<thead>
<tr>
<th>polymer</th>
<th>λ_{max} (nm) 293 K</th>
<th>λ_{max} (nm) 77 K</th>
<th>λ_{max} (nm) film</th>
<th>λ_{max} (nm) film</th>
<th>k_NR (ns) 293 K</th>
<th>k_NR (ns) 77 K</th>
<th>k_S (ns) 293 K</th>
<th>k_S (ns) 77 K</th>
<th>φ (°) 293 K</th>
<th>φ (°) film</th>
<th>ν (cm⁻¹) film</th>
<th>ε (M⁻¹ cm⁻¹) film</th>
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<tr>
<td>BNπ2</td>
<td>404</td>
<td>415</td>
<td>410</td>
<td>462, 511, 545</td>
<td>465, 500, 521</td>
<td>495, 560</td>
<td>660</td>
<td>4150</td>
<td>31.8</td>
<td>0.63</td>
<td>0.32</td>
<td>0.07</td>
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<tr>
<td>BNπ3</td>
<td>422</td>
<td>430</td>
<td>435</td>
<td>494, 534, 573</td>
<td>495, 535, 558</td>
<td>520, 570, 615</td>
<td>680</td>
<td>25000</td>
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<td>0.36</td>
<td>0.30</td>
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<tr>
<td>OPhBNπ2</td>
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<td>410</td>
<td>405</td>
<td>462, 513, 545</td>
<td>464, 500, 516</td>
<td>480, 580</td>
<td>660</td>
<td>28800</td>
<td>1.42</td>
<td>0.21</td>
<td>0.31</td>
<td>0.05</td>
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<tr>
<td>OPhBNπ3</td>
<td>423</td>
<td>440</td>
<td>435</td>
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<td>495, 535, 564</td>
<td>515, 570, 610</td>
<td>680</td>
<td>44385</td>
<td>1.37</td>
<td>0.16</td>
<td>0.43</td>
<td>0.04</td>
</tr>
</tbody>
</table>

* The italic wavelength is the band maximum.

What is observed for the bithiophene copolymers that should adopt a structure similar to that found for the mononaphthalene derivatives, the poly[2,6-(1,5-dioctylnaphthalene)]thiophenes.11

The dependence of the photophysical properties has also been seen with thiophene polymers and copolymers and is found to be strongly influenced by the substitution, degree of backbone twisting in the polymer/oligomer backbone, leading to broad and unstructured emission bands. Moreover, the evidence for aggregates in thin films comes from the absorption spectra which show an increased absorbance at the longer wavelength region as compared to those in their corresponding solution states at 293 K (see Figure 1).

The absorption and emission spectra for the binaphthyl building block of the copolymers (2,2’-octyloxy-1,1’-binaphthyl-BNP; Figure 2) are shown in Figure 3 for BNπ2 and OPhBNπ2 in methylcyclohexane solution display much shorter absorption and emission wavelength maxima than any of the copolymers investigated (Table 2). Moreover the fluorescence quantum yield (φf) is also much greater than any of the values found for the polymers studied (see Table 3).

In the solid state (thin films) a loss of vibronic structure is seen with the concomitant broadening and red shift of the emission spectra. These results provide strong evidence for aggregation in the solid state. The broadening and loss of structure of the emission spectra in the solid state are generally attributed to intermolecular interactions giving rise to different types of packing between the polymer chains and different degrees of twisting in the polymer/oligomer backbone, leading to broad and unstructured emission bands. Moreover, the evidence for aggregates in thin films comes from the absorption spectra which show an increased absorbance at the longer wavelength region as compared to those in their corresponding solution states at 293 K (see Figure 1).

It is also interesting to observe the fluorescence excitation spectra obtained at 293 and 77 K and collected in different regions of the emission spectra, depicted in Figure 3 for BNπ2.
Figure 3. Fluorescence excitation spectra for BNp3 in toluene at 293 and 77 K collected at different emission wavelengths.

Table 4. Fluorescence Decay Times (τ) and Preexponential Factors (a_i) for the Copolymers in Toluene at 293 K and with λ_{exc} = 380 nm.

<table>
<thead>
<tr>
<th>Polymer</th>
<th>λ_{em}</th>
<th>τ_1 (ps)</th>
<th>τ_2 (ps)</th>
<th>τ_3 (ps)</th>
<th>a_1</th>
<th>a_2</th>
<th>a_3</th>
<th>χ^2</th>
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<tr>
<td>BNp2</td>
<td>460</td>
<td>60</td>
<td>420</td>
<td>720</td>
<td>0.38</td>
<td>0.54</td>
<td>0.08</td>
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<td></td>
<td>510</td>
<td></td>
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<td>0.23</td>
<td>0.52</td>
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<tr>
<td></td>
<td>670</td>
<td></td>
<td></td>
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<td>-0.18</td>
<td>0.16</td>
<td>0.84</td>
<td>1.34</td>
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<tr>
<td>BNp3</td>
<td>495</td>
<td>49</td>
<td>434</td>
<td>700</td>
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<td>0.51</td>
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<tr>
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<td></td>
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<td>0.29</td>
<td>0.71</td>
<td>1.13</td>
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<tr>
<td>OPhBNp2</td>
<td>460</td>
<td>40</td>
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<td>-0.20</td>
<td>0.12</td>
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<td>OPhBNp3</td>
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<td>0.58</td>
<td>0.17</td>
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<td>0.58</td>
<td>1.05</td>
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</table>

The data presented were obtained by global analysis of the decays. Also presented are the chi-squared values for a better judgment of the quality of the fits.

A red shift in the wavelength maxima and the appearance of a new band at longer wavelengths are visible upon the increment of the emission wavelength. At low temperature, by collecting the emission at longer wavelengths, there is clear evidence for two distinct bands with maxima at ~440 and 523 nm (Figure 3). This points out to the existence of different segments (chromophoric units) within the polymer with different absorptions. While the assignment of this long wavelength absorption (around 520 nm) is not yet completely clear, we feel that it is unlikely to be due to either impurities or aggregation. The binaphthyl unit is expected to act as a conjugation barrier. However, the longer wavelength band at low temperatures may reflect some of the polymers having longer range conjugation, possibly due to this unit not completely blocking conjugation.

Time-Resolved Fluorescence. Fluorescence lifetimes (τ_f) were obtained in solution (with picosecond time resolution) and in the solid state (with nanosecond time resolution) (see Tables 3 and 4). The lifetimes in films (Table 3) are about 6-fold longer than those in room temperature solution and are of a similar order of magnitude to the radiative lifetimes calculated from room temperature fluorescence quantum yields and lifetimes (1/k_R). In contrast, the fluorescence quantum yields in films and in low temperature glasses are, in general, lower than in solution at room temperature (Table 3). This strongly suggests that the nonradiative deactivation pathways present different contribution in solution and the solid state, most probably resulting from intrachain processes in isolated molecules in solution and interchain interactions (or interactions with the matrix) in the solid state.

The fluorescence lifetimes in toluene solution presented in Table 3 are similar to the major decay components obtained with picosecond time resolution. The fluorescence decay times in solution were also collected as a function of the emission wavelength (Table 4 and Figure 4). In general, the fluorescence decay times do not show any significant change across the fluorescence spectra and can be considered to be independent of the emission wavelength (λ_{em}) over the whole range studied. However, the same is not true with the pre-exponential factors (a_i). This allows the global analysis of the temporal profiles of fluorescence decays. As can be seen from Table 4 and Figure 4, the fluorescence decays for the polymers studied are best fitted with triple-exponential decay laws according to the equation

\[ I_f(t) = a_1 e^{-t/\tau_1} + a_2 e^{-t/\tau_2} + a_3 e^{-t/\tau_3} \]  

where \( a_i \) (i = 1, 2, 3) are the pre-exponential factors and \( \tau_i \) are the decay times (with \( i = 1 \) for \( \lambda_{em} = 460 \) nm, \( i = 2 \) for \( \lambda_{em} = 510 \) nm, and \( i = 3 \) for \( \lambda_{em} = 670 \) nm for BNp2). In Table 4, the shortest lifetime component appears as a decay time at short wavelengths and as a risetime at longer wavelengths.

There is an ongoing controversy on the exact meaning of the fast (few picosecond) components in π-conjugated polymers fluorescence decays. These have been attributed to fast relaxation processes within the polymer backbone including both excitation energy migration in the polymer to the lowest energy conjugation segment and torsional motions leading to conformational reorganization of the excited state. While in general the dominant mechanism for the fast decay has been attributed to energy transfer along the chain, it has been pointed out that conformational relaxation is a general phenomenon occurring in nonrigid polymers and oligomers, which much coexist with other eventual processes such energy transfer. Further, the coexistence of these two processes has also been proposed.

With the present data and for the polymers studied here, differentiation between the two mechanisms is, a priori, uncertain, and thus the nature of the shortest decay component
can be associated with conformational relaxation and/or energy migration. Additional information is however available from the temperature dependence of the decay times shown in Figure 5. Upon increasing temperature, the shortest decay time remains constant within experimental error. Because conformational relaxation strongly depends on temperature and solvent viscosity, we conclude that it is not responsible for the shortest decay component. The absence of detectable relaxation results from the absence of side chains in the naphthyl—oligothiophene—naphthyl chromophoric unit, capable of slowing down the otherwise very fast torsions of the chain at normal viscosities (e.g., toluene). The torsional time has been empirically found to be linearly dependent on the side-chain volume. Extrapolation of this correlation to the van der Waals volume of the hydrogen atom predicts a relaxation time in the subpicosecond time range. Accordingly, dithiophene and terthiophene and other thiophene oligomers do not show detectable conformational relaxation in the picoseconds time range.

The alternative assignment of the fast decay component to energy transfer is not trivial because all segments have the same length and the spectral overlap is small. However, it may occur with slower rate, and also, it is likely that the Dexter exchange mechanism will contribute. This would lead to relatively long times for energy transfer, as observed.

The longest decay time is attributed to the emission of the relaxed species of the polymer, and its contribution increases with the emission wavelength, showing that the emission from the relaxed species is dominant at lower energies. The additional middle decay time (∼390–480 ps) could be related to the emission of an ensemble of conformers with shorter conjugation length (within the polymer) that predominantly emit at shorter wavelengths as seen by the pre-exponential factors and as will be discussed below.

Since there is strong decoupling between the oligothiophene and naphthalene chromophores, it is likely that the relaxed excited species would partially reflect this geometry and would therefore be similar to that present in Scheme 3.

It is again fruitful to compare the decay times of the copolymers studied with the more simple analogues of 1-naphthyl(oligo)thiophenes with two [Naph(α2)] and three [Naph(α3)] thiophene units. As reported elsewhere, the fluorescence lifetimes of these are 360 and 470 ps, respectively. From Table 4 it can be seen that these values are close to the intermediate decay times of the binaphthyl—oligothiophene copolymers investigated. This suggests that the intermediate component could result from the emission of conformers where one of the naphthyl groups, in the naphthyl—oligothiophene—naphthyl, is twisted (and hence loses conjugation). Given the nature of the binaphthyl link to the polymer chain, planarization could be too slow to allow the relaxation to the full planar chromophore.

Within its lifetime. In agreement with this is the observation that the emission spectra of Naph(α2) and Naph(α3) show residual emission above λem = 650 nm and emission maxima at ∼450 nm (435 nm for Naph(α2) and 470 nm for Naph(α3)). With PPV oligomers and polymers, an alternative explanation has been given involving possible trans→cis isomerization upon excitation, with the residual emission attributed to the cis isomer. Although this hypothesis cannot be discarded with these binaphthyl-α-thiophene copolymers, since trans and cis geometries can occur with both the bithiophene and terthiophene units, in the present case we are more in favor of the first hypothesis: isolated segments involving the conjugation of two naphthyls connected by an oligothiophene unit.

**Steady-State Fluorescence Anisotropy.** Additional information on the relaxation processes on these polymers can be obtained from fluorescence (excitation and emission) anisotropy measurements. For these studies, we selected BNpo3 as an illustrative example.

From Figure 6 it can be seen that the fluorescence anisotropy values with excitation at shorter wavelengths (390 nm) gradually decrease from 0.1 at λexc = 460 nm to 0.02 (above ∼575 nm). This can be related to energy transfer between different chromophores. However, with excitation at 420 nm the anisotropy values decrease from 0.1 to 0.05 (at ∼620 nm) and increases thereon. Excitation at longer wavelengths further increases the anisotropy values, and in the case of λexc = 500 nm the value that is reached is close to the maximum value predicted theoretically, consistent with the fact that segments absorbing at longer wavelengths suffer insignificant changes upon excitation. This is also consistent with the observation in Figure 7 where the anisotropy values in the excitation spectra increases sharply for λ > 500 nm.

All of the above indicates that excitation in the tail of the absorption spectra leads to emission of more relaxed conjugation segments which cannot further depolarize (or transfer excess energy) to other segments and consequently display higher anisotropy values.

**Figure 5.** Fluorescence decay times (τi) dependence with temperature for BNpo2 and BNpo3 in toluene.

**Figure 6.** Polarized fluorescence emission spectra and anisotropy values for BNpo3 in toluene at 293 K. The signal S was collected with horizontal (h) and vertical (v) orientation of the excitation (first subscript) and emission (second subscript) polarizers.
the triplet energy of the copolymers. For all the polymers the pulse radiolysis energy transfer method was used to determine the triplet energies. In the absence of phosphorescence we have observed following photolysis of aerated toluene solutions of the copolymers (Table 3). The values were determined by plotting the initial phosphorescence intensity at 1270 nm as a function of the laser dose and comparing the slope with that obtained with 1H-phenalen-1-one in toluene as the standard.

**Photophysical Properties.** From the photophysical parameters obtained in solution and in the solid state (Table 3) it can be seen that as a general trend the radiationless processes ($\phi_F + \phi_{ISC}$) involving in principle isolated polymer chains, are the dominant excited-state deactivation routes in solution for the polymers studied. In contrast, as discussed earlier, in the solid state interactions with the medium play a major role in reducing fluorescence quantum yields. However, from the rate constants presented in Table 3, it is also clear that in all cases radiationless pathways are dominant in the excited-state deactivation.

It is also worth noting that in methycyclohexane both the $\phi_T$ and $r_T$ values slightly decrease relative to toluene (Table 3). The only significant difference occurs with OPPhBNpt3 where the $\phi_T$ value in methycyclohexane decreases ca. 1/2 relative to toluene. This most likely may have its cause in the poor solubility of this polymer in methycyclohexane. Nevertheless, the radiative and radiationless rate constants are approximately identical in both solvents, thus showing that the same deactivation processes are operative independently of the solvent used. This independence of the solvent properties for the photophysical parameters is characteristic of thiophene-like oligomers and polymers.

In contrast to the behavior found for similar copolymers incorporating oligothiophene segments, in the present case we found that the fluorescence quantum yield ($\phi_F$) does not change with the number of thiophene units in the polymer. A similar observation has been found with quinoxaline/oligothiophene copolymers. It is also worth noting that the $\phi_F$ does not seem to be affected by the introduction of octylbenzoxyl side chains in the 1,1′-binaphthyl spacer instead of octyloxyl groups.

Nevertheless, the addition of thienyl units results in a moderate increase in the intersystem-crossing quantum yield ($\phi_{ISC}$) which is made at the expenses of a decrease in the internal conversion decay processes ($\phi_C = 1 - \phi_F - \phi_{ISC}$). Theoretical studies suggest that intersystem crossing in systems containing thiophene units may result from extensive spin–orbit coupling. At low temperature the $\phi_T$ values have in general lower values than those obtained at 293 K. The exception occurs with OPPhBNpt2 where no change in the $\phi_T$ value is observed upon going to 77 K. This could be attributed to the fact that the $S$ state probably has the same geometry as that found at 293 K. The decrease of the $\phi_T$ values upon lowering the temperature can be associated with an induced conformational change, which is also in agreement with the marked red shift of the absorption and emission spectra compared to 293 K. However, at this stage the possibility of aggregation at low temperature cannot be discarded. The decrease in the $\phi_T$ upon going to 77 K contrasts with the behavior found for thiophene oligomers (where an increase or constancy of this value was observed), which is in agreement with what has been found with poly[(2,6-(1,5-diocytlnaphthalene)]thiophenes.

From Table 3 and Figure 1 it can be noted that the $\phi_T$ values in thin films are lower than those obtained in toluene solution at 293 K. This is a general pattern found for $\pi$-conjugated polymers in which the low $\phi_T$ in the solid state are due to the increase of the nonradiative decay processes, which, as discussed.
earlier, results from intermolecular interactions. Again with the copolymers studied, aggregation may be significant as suggested by the red shift and the broadening of the absorption and emission spectra when going from solution to thin films. This result may be associated with increasing intrachain or interchain interaction generated in the solid state of the copolymers.

From Table 3 it can also be seen that singlet oxygen formation yields are very close to the quantum yields for triplet formation, providing support for the latter values and indicating that the efficiency of triplet energy transfer to produce singlet oxygen ($\Delta \phi_s = \phi_s/\phi_T$) is very close to unity. These results suggest that reaction with molecular oxygen may be an important pathway for the deactivation of the triplet state of these copolymers.

Conclusions

A complete spectral and photophysical investigation of four naphthalene–thiophene copolymers with bithiophene and terthiophene units was carried out in solution and in the solid state. Particular interest is focused on the role of the binaphtyl unit in introducing conjugation breaks into the system, which thus show similarities to naphthalene–thiophene copolymers. For all the polymers studied in solution at room temperature, the fluorescence decay is dominated by nonradiative processes on isolated chains. The data obtained follow previous findings with thiophene-based polymers showing the influence of these units in the emissive properties of the polymers and support the dominant role of the bithiophene moiety as a conjugation break.

However, in contrast with a general trend found where the photophysical parameters ($\phi$) increases with the incremental contribution of thiophene units, the substitution of a bithiophene (n = 2) with a terthiophene (n = 3) unit in the polymer does not change the photophysical properties of the polymers. This brings out new developments to the interpretation and full understanding of this behavior. In addition, at low temperatures, a new band is observed in the fluorescence excitation spectra around 520 nm, which may hint at there being some conjugation communication across the binaphtyl linkage. A detailed interpretation of the complex time-resolved fluorescence studies in the picosecond time domain (2 ps time resolution) has been made, indicating that only a decay law with three exponentials can properly fit the observed decays. Rise and decay components are associated with the shortest lifetime component at long and short emission wavelengths, respectively, of the polymer spectra and are attributed to intramolecular energy transfer. While conformational relaxation may exist, it is probably outside our observation time window. Steady-state fluorescence anisotropy data indicate that excitation along the absorption spectra of BNP3 (used as a representative polymer) leads to a progressive increase in the anisotropy values. This is indicative of the existence of a gradual contribution from more relaxed segments, absorbing and emitting at longer wavelengths.

In solution in the presence of oxygen, singlet oxygen formation is observed, with very similar yields to those of triplet formation. In films and low-temperature glasses, a major contribution to the decrease in fluorescence quantum yields comes from intermolecular interactions with the matrix.

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References and Notes
