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# Switching on H-tunneling through conformational control 

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Supporting Information Placeholder


#### Abstract

H-tunneling is a ubiquitous phenomenon, relevant to fields from biochemistry to materials science, but harnessing it for mastering the manipulation of chemical structures still remains nearly illusory. Here, we demonstrate how to switch on H-tunneling by conformational control using external radiation. This is outlined with a triplet 2-hydroxyphenylnitrene generated in an $\mathrm{N}_{2}$ matrix at 10 K by UV-irradiation of an azide precursor. The anti-orientation of its OH moiety was converted to syn by selective vibrational excitation at the $2 v(\mathrm{OH})$ frequency, thereby moving the H -atom closer to the vicinal nitrene center. This triggers spontaneous H -tunneling to a singlet 6-imino-2,4cyclohexadienone. Computations unravel that such fast H -tunneling occurs through crossing the triplet-to-singlet potential energy surfaces. Our experimental realization provides an exciting novel strategy to attain control over tunneling, opening new avenues for directing chemical transformations.


Molecular conformers can dictate the reactivity of chemical ${ }^{1,2}$ and biological systems, ${ }^{3,4}$ materials functions, ${ }^{5,6}$ and even control anticancer activity. ${ }^{7}$ However, it is still challenging to understand how different conformers can lead to different reactions paths or kinetics and what conditions allow conformational control of reactants. ${ }^{8-10}$ Conformational changes are difficult to assess at ambient conditions due to their typically observed fast interconversions. Nonetheless, a few conformer-specific reactions have been demonstrated under low-temperature conditions, paving the way for a better understanding of the conformational influence on the selectivity of chemical reactions and functions. ${ }^{11-16}$
Quantum mechanical tunneling (QMT) is very sensitive to the distances nuclei move in a reaction, ${ }^{17-20}$ therefore the conformation of a reactive moiety can become critical for the occurrence of QMT. ${ }^{21}$ Indeed, one of us recently demonstrated the first examples of conformer dependent H -atom QMT (Scheme 1a). ${ }^{22,23}$ After trapping hydroxycarbenes $\mathbf{1}$ in argon matrices at temperatures as low as 3 K , we observed that the trans-conformer $\mathbf{t - 1}$ rearranges by QMT to aldehyde 2, whereas the cis-conformer $\mathbf{c - 1}$ remained unreactive. We have also demonstrated that different conformations of an aldehyde moiety in the vicinity of a nitrene center give access to different QMT reactions (Scheme 1b). ${ }^{24,25}$ It was observed that the anti-conformer $\mathbf{a}^{-3} \mathbf{3}$, once generated in argon matrices, rearranges by H-atom QMT to imino-ketene 4 whereas the syn-conformer s- ${ }^{3} 5$ cyclizes by heavy-atom QMT to 2,1-benzisoxazole $\mathbf{6}$. Moreover, a rare example of conformer dependent heavy-atom QMT was reported by Sander et al. ${ }^{26,27}$ They observed by matrix isolation spectroscopy that syn-5-methoxy-

2 H -benzazirine $\mathbf{s}-7$ rearranges by QMT to ketenimine $\mathbf{a - 8}$, while the anti-conformer a-7 is too reactive to be detected (Scheme 1c). It has been also suggested that enzymatic $\mathrm{C}-\mathrm{H}$ transfer reactions rely on the dynamic sampling of enzyme-substrate conformations to tune the donor-acceptor distance conductive to QMT. ${ }^{28,29}$ Structural enzymatic modifications either near the bound substrate or at a remote residue has been argued to influence H -atom QMT through changes in conformational flexibility. ${ }^{30-33}$

Scheme 1. Examples of conformer-dependent QMT reactions directly observed by matrix isolation spectroscopy ${ }^{34}$
(a) H-atom QMT in hydroxycarbenes

(b) H- and heavy-atom QMT in 2-formylphenylnitrenes

(c) Heavy-atom QMT in 5-methoxy-2H-benzazirine


While these investigations have begun to decipher how QMT reactivity can significantly depend on conformations, even in moieties remotely located in relation to the reaction center, further breakthroughs can be envisaged through effective conformational control of reactants. Thus, we demonstrate here the experimental realization of how to switch on a QMT reaction through selective conversion of an unreactive to a reactive conformer using vibrational excitation (Scheme 2). ${ }^{35}$ Controlled manipulation of anti-OH triplet 3-fluoro-2-hydroxyphenylnitrene $\mathbf{a}^{-3} \mathbf{1 0}$ to the reactive syn-OH conformer s- ${ }^{\mathbf{3}} \mathbf{1 0}$ was achieved in an $\mathrm{N}_{2}$ matrix at 10 K by narrowband irradiation at the $2 v(\mathrm{OH})$ frequency, triggering fast H -atom QMT to singlet 2 -fluoro-6-imino-2,4cyclohexadienone $\mathbf{Z - 1 1}$. Computations unravel that such fast H-atom QMT occurs through crossing potential energy surfaces (PESs).

Scheme 2. Summary of the H-atom QMT of 2-hydroxyphenylnitrene $\mathrm{s}^{-3} 10$ to imino $\mathbf{Z}-\mathbf{1 1}$ switched on upon conformational isomerization of a- $\mathbf{-}^{\mathbf{1}} \mathbf{1 0}$ by vibrational excitation


The photolysis of parent 2-hydroxyphenylazide 9' under matrix isolation conditions was independently studied by Sander et al. ${ }^{36}$ and Tomioka et al. ${ }^{37}$ Both groups reported the formation of an $E / Z$ mixture of 6-imino-2,4-cyclohexadienone 11', presumably via [1,4]-H shift reaction of 2-hydroxyphenylnitrene 10'. However, neither group detected ground-state triplet nitrene ${ }^{3} \mathbf{1 0}$ ' by matrix isolation spectroscopy. Considering the short distance [ $2.28 \AA$ at B3LYP/6-311+G(2d,p)] between the nitrene center and the hydrogen atom of the OH moiety in syn-conformation, it is likely that the $[1,4]-\mathrm{H}$ shift occurs by fast H -atom QMT, making s- ${ }^{\mathbf{3}} \mathbf{1 0}{ }^{\text {' }}$ elusive. ${ }^{38}$ On the other hand, the longer distance [ $3.73 \AA$ at B3LYP/6-311+G(2d,p)] between the nitrene center and the hydrogen atom of the OH moiety in the anti-conformation makes H -atom QMT unlikely and the capture of $\mathrm{a}-{ }^{\mathbf{3} 10}$ ' conceivable under cryogenic temperatures. The lack of detection of $\mathbf{a -}{ }^{3} \mathbf{1 0}$ ' could be due to the preservation of the syn- OH conformation in the photolysis of $\mathbf{9}^{9},{ }^{39}$ as previously observed for the aldehyde conformation in the photolysis of 2-formylphenylazides (Scheme 1b). ${ }^{24,25}$ CBS-QB3 computations indicate that $9^{\prime}$ ' in gas-phase at room temperature (r.t.) exclusively in the most stable conformation as-9', with the OH moiety in syn-conformation relative to the $\mathrm{N}_{3}$ moiety (Table S1).
To significantly populate the anti-OH conformer of 2-hydroxyphenylazides, we considered the inclusion of a vicinal fluorine atom acting as hydrogen-bond acceptor for the OH moiety, following the approach used to stabilize the more energetic aldehyde conformer in 2-formylphenylazides. ${ }^{24}$ CBS-QB3 computations for the 3-fluoro-2-hydroxyphenylazide 9 predict the anti-OH conformers (sa-9 and aa-9) to be only $2-4 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy than the most stable syn-OH conformer (as-9), i.e., a stabilization of $\sim 10 \mathrm{~kJ} \mathrm{~mol}^{-1}$ relatively to $9^{\prime}$ (Table S1).

Accordingly, the predicted gas-phase equilibrium population at r.t. consists of $\sim 34 \%$ of the anti- OH and $\sim 66 \%$ of the syn- OH conformers, indicating that 9 would be a promising system to access a 2-hydroxyphenylnitrene having an anti-OH conformation (a- ${ }^{3} 10$ ).
3-Fluoro-2-hydroxyphenylazide 9 was synthesized and characterized as described in the SI. Matrices with monomers of 9 were prepared by sublimating the sample at r.t. and depositing it with a large excess of $\mathrm{N}_{2}$ gas onto a CsI optical window ( 10 K ). The IR spectrum of matrix-isolated 9 confirms the presence of both syn- OH and anti- OH conformers in a ratio that roughly mimics the predicted 66:34 ratio (Figures 1 and S 1 ). This is particularly evident in the mid- and near-IR regions corresponding to the $v(\mathrm{OH})$ and $2 v(\mathrm{OH})$ modes (Figure 1). The lower frequency bands centered at 3557 and $6938 \mathrm{~cm}^{-1}$ are assigned to the vibrational modes of syn-OH conformer as- 9 computed at $3573[v(\mathrm{OH})]$ and $6969 \mathrm{~cm}^{-1}[2 v(\mathrm{OH})]$. The slightly overlapping higher frequency bands centered at 3603 and $7039 \mathrm{~cm}^{-1}$ as well as at $\sim 3591$ and $\sim 7009 \mathrm{~cm}^{-1}$ are assigned to the vibrational modes of the anti-OH conformers, sa-9 and aa-9, computed at $3621[v(\mathrm{OH})]$ and 7074 $\mathrm{cm}^{-1}[2 v(\mathrm{OH})]$, and at $3608[v(\mathrm{OH})]$ and $7045 \mathrm{~cm}^{-1}[2 v(\mathrm{OH})]$, respectively. ${ }^{40}$


Figure 1. (a) Selected regions of the experimental mid-IR (right) and near-IR (left) spectrum collected immediately after the deposition of 3-fluoro-2-hydroxyphenylazide 9 in an $\mathrm{N}_{2}$ matrix at 10 K (back line) and the corresponding difference IR spectrum (red line) resulting from irradiation at $\lambda=310 \mathrm{~nm}(10 \mathrm{~min}, 10 \mathrm{~mW})$, subsequent to the full consumption of 9 with a set of irradiations at $\lambda=255 \mathrm{~nm}$. The $2 v(\mathrm{OH})$ band of $\mathbf{a}-{ }^{-} \mathbf{1 0}$ is inferred to be centered at $\sim 6976 \mathrm{~cm}^{-1}$. (b) B3LYP/6-311+G(2d,p) computed anharmonic wavenumbers and IR intensities for the $v(\mathrm{OH})$ (right) and $2 v(\mathrm{OH})$ (left) modes of as-9, aa-9 and sa-9 conformers (black triangles), and of $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ (red circles). The intensities of as- $\mathbf{9}$, aa- $\mathbf{9}$ and sa-9 were scaled considering the computed population ratio of 66:11:23, respectively. The intensity of $\mathbf{a}-{ }^{3} \mathbf{1 0}$ was scaled by an arbitrary factor.
The photolysis of 9 was then promoted by irradiations at $\lambda=255$ or 280 nm [UV-Vis of 9 in acetonitrile: $\lambda_{\text {max }}=\sim 250$ and 280 nm (Figure S3)] and found to lead to (Z)-2-fluoro-6-imino-2,4-cyclohexadienone $\mathbf{Z}-\mathbf{1 1}$, as the major product, and to the antiOH conformer of 3-fluoro-2-hydroxyphenylnitrene $\mathbf{a}^{\mathbf{3}} \mathbf{1 0}$, as the minor product (Figure S4). The most distinctive IR bands of Z-11 appear as a broad feature at $\sim 3216 \mathrm{~cm}^{-1}$ and as two strong signals at 1691 and $1648 \mathrm{~cm}^{-1}$, corresponding to the $v(\mathrm{NH}), v(\mathrm{C}=\mathrm{O})$, and $v(\mathrm{C}=\mathrm{C})_{\text {as }}$ modes, whereas those of $\mathbf{a}^{-}{ }^{\mathbf{3}} \mathbf{1 0}$ appear at $\sim 3576$ and $1571 \mathrm{~cm}^{-1}$, corresponding to the $v(\mathrm{OH})$ and $v(\mathrm{CC})$ modes. The detailed assignment of the IR spectra of both products is addressed in the next section. The successful generation of nitrene
a- ${ }^{\mathbf{3}} \mathbf{1 0}$ in a nitrogen matrix at 10 K confirms our working hypothesis regarding the existence of an "unreactive" anti-OH conformer ( $\mathbf{a}^{-3} \mathbf{1 0}$ ) besides of a "reactive" syn-OH conformer ( $\mathbf{s}-{ }^{3} \mathbf{1 0}$ ). ${ }^{41}$
It has been shown that energy can be selectively deposited in a vibrational mode of a molecular system and be transferred to a torsional reaction coordinate, via an intramolecular relaxation mechanism, for promoting conformational changes. ${ }^{42-46}$ Efficient energy transfer to the torsional reaction coordinate is usually achieved if the stretching bond of the conformational moiety is excited to its first overtone. ${ }^{46-49}$ To implement this approach and attain controlled switch of the "unreactive" $\mathbf{a}-{ }^{3} \mathbf{1 0}$ to the "reactive" $\mathbf{s}^{-} \mathbf{}^{\mathbf{3}} \mathbf{1 0}$, the $2 v(\mathrm{OH})$ frequency of $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ was examined. A dedicated experiment where 9 was totally consumed $(\lambda=255 \mathrm{~nm})$ and then the remaining $\mathbf{a}^{-}{ }^{\mathbf{3} 10}$ was converted to $\mathbf{1 1}(\lambda=310 \mathrm{~nm})$, allowed the measurement of the $v(\mathrm{OH})$ band of $\mathbf{a}-{ }^{-3} \mathbf{1 0}$ centered at $\sim 3576$ $\mathrm{cm}^{-1}$, between the consumed $v(\mathrm{OH})$ bands of as- $\mathbf{9}$ and aa-9 (Figure 1a). Unfortunately, the corresponding $2 v(\mathrm{OH})$ band was not detected, due to its broad and low absorption features alongside with the very small amount of $\mathbf{a}-{ }^{3} \mathbf{1 0}$ present in the matrix. Nevertheless, we can infer that the $2 v(\mathrm{OH})$ band of $\mathbf{a}^{-3} \mathbf{1 0}$ should be centered at $\sim 6976 \mathrm{~cm}^{-1}$, between the $2 v(\mathrm{OH})$ bands of as- 9 $\left(\sim 6938 \mathrm{~cm}^{-1}\right)$ and aa-9 $\left(\sim 7009 \mathrm{~cm}^{-1}\right)$, mirroring the pattern observed for the corresponding $v(\mathrm{OH})$ bands (Figure 1a). This rationalization is supported by B3LYP/6-311+G(2d,p) anharmonic computations, which estimated the $v(\mathrm{OH})$ and $2 v(\mathrm{OH})$ modes of $\mathbf{a}^{-}{ }^{\mathbf{3}} \mathbf{1 0}$ at frequencies between those of as-9 and aa-9 (Figure 1b).
The effect of selective vibrational excitation of $\mathbf{a}^{-3} \mathbf{1 0}$ at its $2 v(\mathrm{OH})$ frequency was then explored. The irradiation of freshly generated $\mathbf{a}^{-3} \mathbf{1 0}$ in an $\mathrm{N}_{2}$ matrix at 10 K with narrowband light tuned to $6976 \mathrm{~cm}^{-1}$ induces its transformation to $\mathbf{Z} \mathbf{- 1 1} .^{50,51}$ The process was monitored by mid-IR spectroscopy using a long-pass filter blocking IR light above $1580 \mathrm{~cm}^{-1}$ to exclude the absorption of photons above the OH -rotamerization barrier $\left[\mathbf{a}-{ }^{-3} \mathbf{1 0} \rightarrow \mathbf{s -}{ }^{\mathbf{3} 10}\right.$; $\sim 1700 \mathrm{~cm}^{-1}$ at B3LYP/6-311+G(2d,p)]. The negative and positive bands of the experimental difference spectrum are well reproduced by the B3LYP/6-311+G(2d,p) computed spectra of a- ${ }^{-10}$ and Z-11, respectively (Figure 2). Some distinctive depleted bands are observed at $1538,1444,1303,768 / 765$, and $696 \mathrm{~cm}^{-1}$, matching the characteristic vibration modes of a- ${ }^{\mathbf{3}} \mathbf{1 0}$ computed at $1536[v(\mathrm{CC})], 1433[\delta(\mathrm{CH})], 1298[v(\mathrm{CN})], 754[\gamma(\mathrm{CH})]$, and 686 $[v(C C)] \mathrm{cm}^{-1}$. Representative produced bands are observed at $1256,1164,1114,836$, and $727 \mathrm{~cm}^{-1}$, matching the characteristic vibration modes of $\mathbf{Z}-\mathbf{1 1}$ computed at $1241[v(\mathrm{CF})], 1163$ $[\delta(\mathrm{CH})], 1098[v(\mathrm{C}-\mathrm{C})], 828[\delta($ ring $)]$, and $723[\gamma(\mathrm{CH})] \mathrm{cm}^{-1} . \mathrm{A}$ detailed assignment of the IR spectra of $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z} \mathbf{- 1 1}$ is given in Tables S4 and S5.


Figure 2. (a) B3LYP/6-311+G(2d,p) computed IR spectrum of triplet anti-3-fluoro-2-hydroxyphenylnitrene $\mathbf{a}^{-3} \mathbf{1 0}$. (b) Experimental difference IR spectrum showing changes resulting from irradiation at $6976 \mathrm{~cm}^{-1}(60 \mathrm{~mW}, 1 \mathrm{~h})$, subsequent to irradiation of 9 at $\lambda=255 \mathrm{~nm}$ (Figure S4). The downward bands are due to consumed species assigned to $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$. The upward bands are due to the produced species assigned to $\mathbf{Z}-\mathbf{1 1}$. Asterisk (*) indicate minor changes in the population of conformers of 9. (c) B3LYP/6$311+\mathrm{G}(2 \mathrm{~d}, \mathrm{p})$ computed IR spectrum of (Z)-2-fluoro-6-imino-2,4cyclohexadienone $\mathbf{Z}-\mathbf{1 1}$.

To better understand the "reactive" nature of syn-OH nitrene $\mathbf{s}-{ }^{\mathbf{3}} \mathbf{1 0}$, the PES around this species was computed (Figure 3). A conceivable mechanism for the reaction of $\mathbf{s}_{-}{ }^{3} \mathbf{1 0}$ is the occurrence of $\mathbf{H}$-atom QMT on the triplet surface to $\mathbf{Z}-{ }^{-3} \mathbf{1 1}$, followed by intersystem crossing to $\mathbf{Z}-\mathbf{1 1} .{ }^{52}$ However, the height ( $\sim 98.2 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ) of the barrier from $\mathbf{s -}^{-3} \mathbf{1 0}$ to $\mathbf{Z -}{ }^{-3} \mathbf{1 1}$ seems incompatible with the occurrence of a very fast QMT reaction justifying the elusive nature of $\mathbf{s}^{3} \mathbf{1 0}$; its half-life time would be $1.14 \times 10^{3} \mathrm{~s}$ based on the Wentzel-Kramers-Brillouin approximation (see SI). An alternative mechanism for the reaction of $\mathbf{s}^{-3} \mathbf{1 0}$ is the occurrence of H -atom QMT to $\mathbf{Z - 1 1}$ through crossing surfaces from triplet to singlet. To explore this possibility, we searched for a minimumenergy crossing point (MECP) connecting the triplet $\mathbf{s}^{-3} \mathbf{1 0}$ and singlet $\mathbf{Z - 1 1}$ surfaces, first at $\operatorname{CASSCF}(8,8) / 6-31 \mathrm{G}(\mathrm{d})$ and then at B3LYP/6-311+G(2d,p) (see SI). Starting with a geometry where the two PESs cross as a function of $\mathrm{r}(\mathrm{O}-\mathrm{H})$ distance, a global optimization algorithm was employed and an MECP was found with a relative energy of $\sim 39.3 \mathrm{~kJ} \mathrm{~mol}^{-1}$ at B3LYP/6-311+G(2d,p) + ZPVE (Figure 3). The spin-forbidden H -atom QMT rate was then calculated by non-adiabatic transition state theory (NA-TST) using weak coupling (WC) formulation (see SI). ${ }^{53}$ Considering the MECP energy and all additional data from the B3LYP/6-311+G(2d,p) computations, the calculated WC rate constant at 10 K yields for $\mathbf{s}^{-3} \mathbf{1 0}$ an half-life of $\sim 2.60 \times 10^{-2}$ s. This indicates that $\mathbf{s}^{-} \mathbf{- 1 0} \mathbf{1 0}$ reacts promptly by H -atom QMT through crossing triplet to singlet PESs to $\mathbf{Z - 1 1}$, deciphering its reactive and elusive nature.


Intrinsic Reaction Coordinates / Bohr
Figure 3. B3LYP/6-311+G(2d,p) computed intrinsic reaction coordinate (IRC) profile for the OH -rotamerization of ${ }^{\mathbf{3} 10}$, for the H -shift of $\mathbf{s -}^{\mathbf{3}} \mathbf{1 0}$ to $\mathbf{Z -}{ }^{\mathbf{3}} \mathbf{1 1}$, and from the MECP connecting $\mathbf{s}^{\mathbf{3}}{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z}$-11. For visualization proposes, the IRC curve from the MECP was superimposed on the IRC curve for H -shift and shifted along the horizontal axis so the point corresponding to s- ${ }^{\mathbf{3}} \mathbf{1 0}$ has the same reaction coordinate in both curves; all the other reaction coordinates represent different structures for each IRC profile. The horizontal bars at each single point (minima, transition states and MECP) show the corresponding computed ZPVE-corrected energy relative to $\mathbf{a}-{ }^{3} \mathbf{1 0}$.

In conclusion, we demonstrate how to switch on a QMT reaction by conformational control using external radiation. Upon trapping triplet 2-hydroxyphenylnitrene $\mathbf{a}-\mathbf{3} \mathbf{1 0}$ in an $\mathrm{N}_{2}$ matrix at 10 K , its OH moiety was manipulated (anti $\rightarrow$ syn) using vibrational excitation at the $2 v(\mathrm{OH})$ frequency, moving the H -atom closer to the vicinal nitrene center. Consequently, fast H -atom QMT from s- ${ }^{\mathbf{3}} \mathbf{1 0}$ to singlet $\mathbf{Z - 1 1}$ was triggered. NA-TST computations unravel that such a reaction occurs through crossing PESs, establishing in this way a rare model for a nonadiabatic H -atom QMT. Overall, our work provides an important conceptual strategy to harness the control of QMT, which will inspire new advances that can extend from enzymatic catalysis ${ }^{54}$ to quantum switches. ${ }^{55}$

## ASSOCIATED CONTENT

## Supporting Information

The Supporting Information is available free of charge on the ACS Publications website.
Experimental and computational procedures, additional experimental results, vibrational assignments, MECP and tunneling rate computations, and computational data (PDF).

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## Notes

The authors declare no competing financial interests.

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Supporting Information (SI)

# Switching on H-tunneling through conformational control 

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## 1. Experimental and Computational Methods

## Synthesis of 3-fluoro-2-hydroxyphenylazide 9

General: Commercial reagents were used as purchased. ${ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}$ Nuclear Magnetic Resonance (NMR) spectra were recorded on an NMR spectrometer Bruker Avance III operating at 400 and 100 MHz , respectively. Chemical shifts are referred to the residual signal of DMSO- $\mathrm{d}_{6},{ }^{1} \mathrm{H}$ NMR $\left(\delta_{\mathrm{H}}=2.50\right)$ and ${ }^{13} \mathrm{C}$ NMR $\left(\delta_{\mathrm{C}}=39.5\right)$, or to the internal standard tetramethylsilane (TMS). Chemical shifts are given in parts per million ( ppm ) relative to TMS and coupling constants $J$ are given in Hertz. Thin-layer chromatography (TLC) was carried out on silica gel $60 \mathrm{~F}_{254}$ plates (AL TLC $20 \times 20$ ). Column chromatography was performed on Silica Gel $60(0.04-0.063 \mathrm{~mm})$. Highresolution mass spectra (HRMS) were obtained with a Waters Micromass VG Autospec M ESI-TOF spectrometer.

Synthesis: The synthesis of 3-fluoro-2-hydroxyphenylazide 9 (Scheme S1) was carried out based on a procedure described in the literature for the preparation of other azidophenol derivatives. ${ }^{1}$


Scheme S1. Synthesis of 3-fluoro-2-hydroxyphenylazide 9.

Concentrated $\mathrm{HCl}(37 \% \mathrm{v} / \mathrm{v}, 1.1 \mathrm{~mL})$ was added dropwise to a suspension of 2-amino-6-fluorophenol ( $0.5 \mathrm{~g}, 3.9 \mathrm{mmol}$ ) in water $(12.5 \mathrm{~mL})$, at $0{ }^{\circ} \mathrm{C}$, followed by the addition of a solution of $\mathrm{NaNO}_{2}(0.27 \mathrm{~g}, 3.9 \mathrm{mmol}, 1 \mathrm{eq})$ in water $(1.2 \mathrm{~mL})$. After 10 min , the cooling bath was removed and the reaction mixture was allowed to warm to room temperature. Sodium azide ( $0.31 \mathrm{~g}, 4.8$ mmol and 1.2 eq ) was added portionwise and, after stirring at room temperature for 1 h , the reaction mixture was extracted with EtOAc ( 25 mL ), the organic layer washed with brine ( 25 mL ), water ( 25 mL ) and dried over anhydrous sodium sulphate. The solvent was removed by evaporation followed by purification of the crude product by flash chromatography [ethyl acetate/hexane (1:2)]. The azide 9 was obtained as a light brown solid in $82 \%$ yield ( $0.49 \mathrm{~g}, 3.2$ mmol ).

Characterization of 9: ${ }^{1} \mathrm{H}$ NMR ( 400 MHz, DMSO-d ${ }_{6}$ ): $\delta(\mathrm{ppm}) 10.21$ (bs, $1 \mathrm{H}, \mathrm{OH}$ ), 7.04-6.97 (m, 1H, Ph-H), 6.86-6.79 (m, 2H, Ph-H). ${ }^{13}$ C NMR (100 MHz, DMSO-d ${ }_{6}$ ): $\delta(\mathrm{ppm}) 151.9$ (d, $\left.J_{\mathrm{CF}}=239.5 \mathrm{~Hz}\right), 138.1\left(\mathrm{~d}, J_{\mathrm{CF}}=16.7 \mathrm{~Hz}\right), 128.6\left(\mathrm{~d}, J_{\mathrm{CF}}=4.8 \mathrm{~Hz}\right), 119.4\left(\mathrm{~d}, J_{\mathrm{CH}}=8.4 \mathrm{~Hz}\right), 116.3$ $\left(\mathrm{d}, J_{\mathrm{CH}}=2.7 \mathrm{~Hz}\right), 112.6\left(\mathrm{~d}, J_{\mathrm{CH}}=18.8 \mathrm{~Hz}\right) . \mathrm{HRMS}\left(\mathrm{ESI}^{+}, m / z\right)$ calcd for $\mathrm{C}_{6} \mathrm{H}_{4} \mathrm{FN}_{3} \mathrm{ONa}(\mathrm{M}+\mathrm{Na})^{+}$ $=176.0231$, found 176.0232 .

Matrix Isolation IR Spectroscopy: A sample of 3-fluoro-2-hydroxyphenylazide 9 was placed in a glass tube, which was then connected to a closed-cycle helium cryostat (Advanced Research Systems DE-202) through a stainless-steel needle valve (SS-4BMRG NUPRO). Possible traces of volatile impurities were removed by pumping the sample through the cryostat at room temperature. Monomeric matrices were prepared by co-deposition of vapors of 9 at room temperature and a large excess of nitrogen gas (N50, Air Liquide) onto a cold (10 K) CsI optical substrate. The temperature of the CsI window was directly measured by a silicon diode, connected to a digital controller (LakeShore 311), and stabilized with an accuracy of 0.1 K .

Infrared spectra were recorded using a Thermo Nicolet 6700 Fourier transform infrared (FTIR) spectrometer, equipped with a Mercury Cadmium Telluride (MCT-B) detector (cooled with liquid nitrogen) and a KBr beam splitter, for the mid-IR range ( $4000-400 \mathrm{~cm}^{-1}$ ), or a Indium Gallium Arsenide (InGaAs) detector and a $\mathrm{CaF}_{2}$ beam splitter, for the near-IR range (7500-4000 $\mathrm{cm}^{-1}$ ). When there was the need to protect the sample from part of the infrared light, the mid-IR spectra were collected in the range $1580-400 \mathrm{~cm}^{-1}$ by using a long-pass filter (Spectrogon LP6300 nm , transmission cut-off value $\sim 1587 \mathrm{~cm}^{-1}$ ) placed between the spectrometer light source and the sample. The infrared spectra were recorded with resolution $0.5 \mathrm{or}^{2 \mathrm{~cm}^{-1}}$ in the mid-IR and near-IR ranges, respectively. To reduce the interference from atmospheric $\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{CO}_{2}$, a stream of dry and $\mathrm{CO}_{2}$ filtered-off air was used to continuously purge the optical path of the spectrometer.

UV and Near-IR Irradiation Experiments: The matrix-isolated species were irradiated through the outer KBr window of the cryostat using a frequency-tunable narrowband light [full width at half-maximum (fwhm) of $\left.\sim 0.2 \mathrm{~cm}^{-1}\right]$ generated by a frequency-doubled signal (UV range) or an idler (near-IR range) beam of an optical parametric oscillator (Spectra Physics Quanta-Ray MOPO-SL) pumped with a pulsed Nd:YAG laser (Spectra-Physics PRO-230: output power $\sim 4.4 \mathrm{~W}$; wavelength $=355 \mathrm{~nm}$; duration $=10 \mathrm{~ns}$; repetition rate $=10 \mathrm{~Hz}$ ).

IR Spectra and Thermochemistry Computations: The geometry optimizations and harmonic frequency computations were performed for the most relevant stationary points of $\mathbf{9}, \mathbf{9}^{\mathbf{\prime}}, \mathbf{1 0}$, and 11 at the B3LYP/6-311+G(2d,p) level of theory, ${ }^{2-5}$ using the Gaussian 16 software package. ${ }^{6}$ All computations were performed using tight optimization criteria. The nature of each stationary point was inspected by the analysis of the Hessian matrices. To take into consideration the neglected anharmonic effects, the limitations of the implemented methods and the limitations of the basis set, the harmonic vibrational frequencies were scaled by a factor of 0.979 or 0.955 for wavenumbers above or below $2500 \mathrm{~cm}^{-1}$, respectively. ${ }^{7}$ The scaled harmonic vibrational frequencies and respective absolute intensities were used to simulate the IR spectra by convoluting each peak with a Lorentzian function having an fwhm of $2 \mathrm{~cm}^{-1}$. Anharmonic frequency computations for $\mathbf{9}$ and $\mathbf{a}-{ }^{-3} \mathbf{1 0}$ were carried out at the B3LYP/6-311+G(2d,p) level of theory, using the fully automated generalized second-order vibrational perturbative theory (GVPT2), as implemented in Gaussian 16. ${ }^{8-10}$ Thermochemistry computations for $\mathbf{9}$ and $\mathbf{9}^{\prime}$ were carried out at the B3LYP/6-311+G(2d,p) and CBS-QB3 ${ }^{11,12}$ levels of theory. The computed data were then used to estimate the relative Gibbs energy at $298.15 \mathrm{~K}\left(\Delta \mathrm{G}_{298 \mathrm{~K}}\right)$ for the different conformers and their room temperature Boltzmann populations.

Normal mode analysis: The theoretical normal modes of $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z}-\mathbf{1 1}$ were analyzed by carrying out potential energy distribution (PED) calculations. The calculated force constants with respect to Cartesian coordinates, obtained from the Gaussian 16 computations, were transformed into the force constants with respect to internal coordinates, which allowed the PED analysis to be carried out as described elsewhere. ${ }^{13}$ The set of internal coordinates used for $\mathbf{a}-{ }^{3} \mathbf{1 0}$ and $\mathbf{Z}-\mathbf{1 1}$ were defined as recommended by Pulay et al. ${ }^{14}$ and are given in Tables S2 and S3, respectively. The atom numberings of $\mathbf{a}-{ }^{-10}$ and $\mathbf{Z}-\mathbf{1 1}$, used for the definition of the internal coordinates, are shown in Figure S5. The vibrational assignment of $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z} \mathbf{- 1 1}$ and the obtained PED matrices are presented in Tables S4 and S5, respectively.

MECP Computations: Due to the open-shell nature of the nitrene species, the potential energy surface (PES) connecting $\mathbf{a -}{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z}-\mathbf{1 1}$ was preliminarily screened at the multiconfigurational $\operatorname{CASSCF}(8,8) / 6-31 \mathrm{G}(\mathrm{d})$ level of theory using the Gaussian 16 software package. ${ }^{6}$ An eight electron- eight orbital $(8,8)$ complete active space (CAS) was used for all species. For nitrene 10,
the defined CAS consists of seven out-of-plane $\pi / \pi^{*}$ orbitals (A") and one in-plane p atomic orbital (A') (Figure S6). ${ }^{15}$ We computed the minimum corresponding to two conformers of the triplet ground-state (TGS) nitrene ( $\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{s}^{\mathbf{3}} \mathbf{1 0}$ ), two conformers of the open-shell singlet (OSS) nitrene $\left(\mathbf{a}_{-}{ }^{\mathrm{OSS}} \mathbf{1 0}\right.$ and $\left.\mathbf{s}-{ }^{\mathrm{OSS}} \mathbf{1 0}\right)$, and one conformer of the closed-shell singlet (CSS) nitrene $\left(\mathbf{a}^{-}{ }^{\mathrm{CSS}} \mathbf{1 0}\right)$. The putative $\mathbf{s}^{\mathrm{CSS}} \mathbf{1 0}$ conformer, with a OH group facing the reactive nitrene center could not be located as a minimum. For the imino $\mathbf{Z} \mathbf{- 1 1}$, the defined CAS consists of eight out-of-plane $\pi / \pi^{*}$ orbitals (A") (Figure S6). We computed the minimum corresponding to the singlet groundstate imino ( $\mathbf{Z}-\mathbf{1 1}$ ) and to triplet state imino $\left(\mathbf{Z}-{ }^{3} \mathbf{1 1}\right)$.

The nitrene $\mathbf{s - ~}^{\mathbf{3}} \mathbf{1 0}$ and imino $\mathbf{Z} \mathbf{- 1 1}$ have different ground state multiplicities (triplet and singlet) and, therefore, the existence of a minimum-energy crossing point (MECP) connecting both species was investigated. An initial MECP search was performed using the so-called partial optimization method. In this case, by running relaxed PES scans as a function of the $\mathrm{r}(\mathrm{OH})$ coordinate (i.e., optimizing all internal coordinates while incrementally fixing the OH bond distance), starting from the minima found before; $\mathbf{s}_{-}{ }^{\mathbf{1}} \mathbf{1 0}, \mathbf{s}_{-}{ }^{\text {Oss }} \mathbf{1 0}, \mathbf{Z}-\mathbf{1 1}$, and $\mathbf{Z -}{ }_{-}^{\mathbf{3}} \mathbf{1 1}$ (Figure S7). The increment used between each scan step was $0.02 \AA$ (regarding the OH bond distance), a small value to avoid discontinuities in the active space. The partial optimization method is a modest approximation for the search of a MECP because all the coordinates except the $\mathrm{r}(\mathrm{OH})$ are optimized separately in each surface, rather than converging together. ${ }^{16,17}$ Nonetheless, this approach provides useful insights into the PES investigated. First, it shows the dissociative character of the $\mathbf{s}^{\mathrm{CSS}} \mathbf{1 0}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ surface along the $\mathrm{r}(\mathrm{OH})$ coordinate as connecting to singlet $\mathbf{Z}-\mathbf{1 1}$, thus explaining the lack of success in finding the s- ${ }^{\mathrm{CSS}} \mathbf{1 0}$ minimum. Second, it shows that $\mathbf{s}-{ }^{\mathrm{OSS}} \mathbf{1 0}$ $\left({ }^{1} \mathrm{~A}^{\prime \prime}\right)$ surface is parallel to $\mathbf{s}^{-3} \mathbf{1 0}\left({ }^{3} \mathrm{~A}\right.$ ") surface, and therefore, the existence of a MECP between these surfaces is unlikely. Finally, it shows that it is likely that there exists an MECP at the crossing between the triplet $\mathrm{A}^{\prime \prime}$ and CSS A' surfaces, directly connecting $\mathbf{s}^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z} \mathbf{- 1 1}$. A subsequent full MECP search was performed using a global optimization algorithm employing the EasyMECP software package,,$^{18}$ which is a python script developed based on the work of Harvey et al. ${ }^{19}$ This software was used as default, requiring solely the inclusion of an energy-parse code to allow reading the data of CASSCF computations performed in Gaussian 16. The structure on the singlet surface previously obtained by the partial optimization method, near the crossing between triplet $\mathrm{A}^{\prime \prime}$ and CSS A' surfaces, was used as the input geometry. The initial guess active space orbitals of the input structure at triplet and singlet multiplicities were those of the active space orbitals of
$\mathbf{a}^{-}{ }^{\mathbf{3}} \mathbf{1 0}$ and Z-11, respectively. Using this method, one MECP was found. Vibrational analysis on the seam of the MECP was carried out using the GLOWfreq software package. ${ }^{20}$ Here, two Hessian matrices obtained on each of the diabatic surfaces are combined in one effective Hessian. Then the motion orthogonal to the crossing seam is projected out and the effective Hessian is finally diagonalized to provide the effective vibrational frequencies of the MECP. Finally, intrinsic reaction coordinate (IRC) computations with each multiplicity were performed by following the steepest descendent path from the optimized MECP structure, which confirms that the found MECP connects $\mathbf{a}^{-3} \mathbf{1 0}$ and $\mathbf{Z}-\mathbf{1 1}$.

To properly account for dynamic electron correlation and to obtain more reliable energies, CASSCF computations are usually augmented by single point computations using a method such as CASPT2. However, the geometry of an MECP can be extremely dependent on the level of theory. Because the MECP depends on the relative positions of two surfaces, changing the level of theory from CASSCF to CASPT2 for performing single point calculations can shift the surfaces relative to one another and lead to unreliable energy values. ${ }^{19,21}$ Alternatively, one can perform the MECP search using a time affordable computational method that already includes dynamic electron correlation such as DFT, as long as the method chosen is adequate for the system being studied (e.g., the species to be computed do not possess multiconfigurational character). The need to use multiconfigurational methods for a correct description of most nitrenes species is mainly because of their OSS states. ${ }^{22}$ Triplet ground-state arylnitrenes, such as the $\mathbf{s}^{-} \mathbf{1 0}$, are properly and reliably described by DFT computations. The product $\mathbf{Z} \mathbf{- 1 1}$ is a closed-shell singlet system and there is no need for using multiconfigurational computations. Since the MECP found during the CASSCF search is at the crossing between a triplet and a singlet surface, connecting the $\mathbf{s}^{-} \mathbf{1 0} \mathbf{1 0}$ and $\mathbf{Z - 1 1}$, it is reasonable to assume that this MECP can be correctly described by DFT methods. Thus, we performed B3LYP/6-311+G(2d,p) relaxed PES scans along the $\mathrm{r}(\mathrm{OH})$ coordinate and obtained qualitatively identical results to those performed at the CASSCF level (compare Figures S7 and S8). Accordingly, the partial optimization method carried out at the B3LYP level also suggests the existence of an MECP at a crossing between the triplet $\mathrm{A}^{\prime \prime}$ and CSS A' surfaces, connecting the $\mathbf{s}^{-}{ }^{\mathbf{1} 10}$ and $\mathbf{Z - 1 1}$ species. A subsequent full MECP search was performed at the B3LYP/6-311+G(2d,p) level, according to the method described above for the similar search performed at the $\operatorname{CASSCF}(8,8) / 6-31 \mathrm{G}(\mathrm{d})$ level. Additional details are provided in the main text.

Tunneling Rate Computations: The transformation of triplet $\mathbf{s}-^{\mathbf{3}} \mathbf{1 0}$ to singlet $\mathbf{Z} \mathbf{- 1 1}$ is a spinforbidden reaction involving H -atom tunneling. Two different mechanisms can be envisioned: (i) A two-step mechanism involving an H -atom QMT on the triplet surface from $\mathbf{s}^{-3} \mathbf{1 0}$ to $\mathbf{Z}-{ }^{\mathbf{3}} \mathbf{1 1}$ followed by intersystem crossing (ISC) to singlet ground state $\mathbf{Z}-\mathbf{1 1}$; (ii) A single-step mechanism involving an H -atom QMT through crossing triplet and single surfaces from ${ }^{3} \mathbf{1 0}$ to $\mathbf{Z} \mathbf{- 1 1}$.
(i) The H -atom QMT rate on the triplet surface (first hypothesis) was computed by the Wentzel-Kramers-Brillouin (WKB) approximation. ${ }^{23-25}$ According to this model, the probability $P(E)$ of tunneling is given by:

$$
\begin{equation*}
P(E)=\mathrm{e}^{-\pi^{2} w \sqrt{2 m\left(V_{0}-E\right)} / h} \tag{1.1}
\end{equation*}
$$

where m is the mass of the particle tunneling through a barrier with height $V_{0}$ and width $w,\left(V_{0}-E\right)$ is the energy difference of the particle E with respect to the top of the barrier, and $h$ is Planck's constant. Tunneling rates are given by the product of the tunneling probability and the frequency of attempts (i.e., the frequency of the normal mode associated to the reaction coordinate). To obtain the parameters of the barrier, intrinsic reaction coordinate (IRC) profiles connecting s- ${ }^{\mathbf{3}} \mathbf{1 0}$ to $\mathbf{Z}-{ }^{\mathbf{3}} \mathbf{1 1}$ were computed at the B3LYP/6-311+G(2d,p). The strictly electronic energy profiles were then corrected by the zero-point vibrational energy of the stationary points. After this procedure, the relative barrier height ( $V_{0}-E$ ) was computed to be $98.2 \mathrm{~kJ} \mathrm{~mol}^{-1}$ and the barrier width $w$ to be 2.10 bohr (Figure S9). Using equation 1.1, the tunneling probability was estimated to be $1.49 \times 10^{-17}$. Considering the computed $\delta(\mathrm{OH})=1361 \mathrm{~cm}^{-1}$ of $\mathbf{s}^{\mathbf{3}} \mathbf{1 0}$ as the frequency of attempts, the tunneling rate was estimated to be $6.07 \times 10^{-4} \mathrm{~s}^{-1}$ (half-life of $1.14 \times 10^{3} \mathrm{~s}$ ).
(ii) The H -atom QMT rate through crossing triplet and singlet surfaces (second hypothesis) was computed with the weak coupling (WC) $)^{26-28}$ formulation of non-adiabatic transition state theory (NA-TST), ${ }^{26,29-34}$ which we have employed very recently for the calculation of spin-forbidden QMT rate constants. ${ }^{21,35}$ The reader is referred to those papers for a detailed description of the procedure. By running the GLOWfreq code with data obtained from B3LYP/6-311+G(2d,p) computations (see above), we obtain the quantities required by the NA-TST equation such as the norm of the difference of the gradients on the two surfaces at the MECP ( $1.26 \times 10^{-1}$ hartree bohr $^{-1}$ ), the geometric mean of the norms of the mentioned gradients ( $4.05 \times 10^{-2}$ hartree bohr $^{-1}$ ), the zero-point energy at the MECP ( 0.0840 hartrees, which leads to an MECP barrier of $\sim 42.1 \mathrm{~kJ}$ $\mathrm{mol}^{-1}$ relative to $\mathbf{s}^{\mathbf{-}} \mathbf{1 0}\left[\sim 39.3 \mathrm{~kJ} \mathrm{~mol}^{-1}\right.$ relative to $\left.\left.\mathbf{a}-{ }^{\mathbf{3}} \mathbf{1 0}\right]\right)$ and the reduced mass of the mode orthogonal to the crossing seam surface ( 11.2 amu ). However, as suggested by Harvey, ${ }^{29}$ we can
expect that in the deep tunneling regime of this particular reaction, the process of breaking an $\mathrm{O}-\mathrm{H}$ bond and forming an $\mathrm{N}-\mathrm{H}$ bond will be dominated by the motion of the hydrogen atom alone. Therefore, instead of using 11.2 amu for the reduced mass, we have used a value of 1 amu . The spin-orbit coupling, which also enters the NA-TST equation, was calculated at the B3LYP/6$311+G(2 d, p)$ level with MolSOC ${ }^{36-38}$ and found to be $21.17 \mathrm{~cm}^{-1}$. With these data, the NA-TST rate constant at 10 K yields $26.7 \mathrm{~s}^{-1}$ (half-life of 0.026 s ).

## 2. Figures



Figure S1. (a) Experimental IR spectrum of 3-fluoro-2-hydroxyphenylazide 9 in a nitrogen matrix at 10 K . The asterisk symbol indicates contribution from monomeric water. (b) IR spectrum of 9 computed at the B3LYP/6-311+G(2d,p) level of theory considering the population of as-9, aa-9 and sa-9 conformers in a ratio of 66:11:23 (see Table S1).

(b)
(a)

Figure S2. (a) Selected regions showing the mid-IR spectrum collected immediately after the deposition of 3-fluoro-2-hydroxyphenylazide 9 in a nitrogen matrix at 10 K (back line) and the corresponding difference IR spectrum resulting from irradiation at $6938 \mathrm{~cm}^{-1}(60 \mathrm{~mW}, 1 \mathrm{~h})$ (gray line). Negative bands indicate the consumption of as-9 and positive bands the production of aa-9. The conformer sa-9 remains unchanged. (b) Anharmonic wavenumbers and IR intensities computed at the B3LYP/6-311+G(2d,p) level for the as-9, aa-9 and sa-9 conformers, considering the population ratio of 66:11:23, respectively.


Figure S3. UV-Vis spectrum of 3-fluoro-2-hydroxyphenylazide (9) in acetonitrile (ACN). The spectrum was recorded in a Shimadzu UV2501-PC at room temperature.


Figure S4. Experimental difference IR spectrum showing changes after irradiation of 9 at $\lambda=255 \mathrm{~nm}(1 \mathrm{~min}, 20 \mathrm{~mW})$. Negative bands are due to the consumption of 9 and the positive bands are due to the production of $\mathbf{a - ~}^{-3} \mathbf{1 0}(\bullet$, closed red circles) and $\mathbf{Z - 1 1}$ ( $\square$, open blue squares). Further details regarding the identification of $\mathbf{a -}{ }^{\mathbf{3}} \mathbf{1 0}$ and $\mathbf{Z} \mathbf{- 1 1}$ are discussed in the main text.


Figure S5. Geometry and atom numbering of triplet anti-3-fluoro-2-hydroxyphenylnitrene a- ${ }^{\mathbf{3}} \mathbf{1 0}$ and (Z)-2-fluoro-6-iminocyclohexa-2,4-dienone $\mathbf{Z - 1 1}$ used for the definition of internal coordinates. Color codes: blue - nitrogen, red - oxygen, green - fluor, gray - carbon, white hydrogen.


Figure S6. Active space orbitals used in the $\operatorname{CASSCF}(8,8)$ computations for nitrene $\mathbf{1 0}$ and imino Z-11 species. Orbital occupation of each specific state can be found in 4 - Computational Section.


Figure S7. $\operatorname{CASSCF}(8,8) / 6-31 \mathrm{G}(\mathrm{d})$ relaxed PES scans as a function of the OH distance for nitrene $\mathbf{1 0}$ and imino $\mathbf{Z - 1 1}$, considering the difference in the multiplicity [singlet (OSS and CSS) shown by blue circles and triplet by red squares] and symmetry [ A ' symmetry shown in open shape and $\mathrm{A}^{\prime \prime}$ in closed shape]. The arrows indicate the starting point and direction of the scans.


Figure S8. B3LYP/6-311+G(2d,p) relaxed PES scans as a function of the OH distance computed for nitrene $\mathbf{1 0}$ and imino $\mathbf{Z - 1 1}$, considering the difference in multiplicity [closed-shell singlet (CSS) shown in blue circles and triplet in red squares] and symmetry [A' symmetry shown in open shape and A" in closed shape]. The arrows indicate the starting point and direction of the scans.


Figure S9. B3LYP/6-311+G(2d,p) intrinsic reaction coordinate (IRC) profile for the H -shift of $\mathbf{s}^{\mathbf{3}} \mathbf{1 0}$ to $\mathbf{Z}-{ }^{\mathbf{3}} \mathbf{1 1}$. The relative electronic energy is given in relation to the energy of $\mathbf{s -}{ }^{\mathbf{3}} \mathbf{1 0}$. The vertical arrow establishes the zero-point corrected energy of the reactant $\mathbf{s}^{\mathbf{3}} \mathbf{1 0}$ relative to the transitionstate ( $\mathrm{h}=98.2 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ). The horizontal arrow establishes the barrier width considering the zeropoint corrected energy values of the stationary points superimposed with the pure electronic IRC energy profile ( $\mathrm{w}=2.10$ Bohr).


Figure S10. ${ }^{1} \mathrm{H}$ (top) and ${ }^{13} \mathrm{C}$ (bottom) NMR spectra of 3-fluoro-2-hydroxyphenylazide (9) in $\mathrm{DMSO}^{-} \mathrm{d}_{6}$.


Figure S11. Electrospray ionization mass spectrum in positive-ion mode (HRMS-ESI ${ }^{+}$) of 3-fluoro-2-hydroxyphenylazide (9).

## 3. Tables

Table S1. Relative Gibbs energy at $298 \mathrm{~K}\left(\Delta G_{298 \mathrm{~K}}\right.$ in $\left.\mathrm{kJ} \mathrm{mol}{ }^{-1}\right)$ computed at the B3LYP/6-311+G(2d,p) and CBS-QB3 levels of theory for conformers of 2-hydroxyphenylazide $\mathbf{9}^{\prime}$ and 3-fluoro-2-hydroxyphenylazide 9 and their equilibrium populations at 298 K (Pop P988 in $\%$ ). ${ }^{\text {a }}$

| Structures $\begin{aligned} & \mathbf{9}^{\prime}, 9 \\ & \mathrm{R}=\mathrm{H}, \mathrm{~F} \end{aligned}$ |  |  |  |
| :---: | :---: | :---: | :---: |
| Name | as-9' ${ }^{\text {as-9 }}$ | aa-9' ${ }^{\text {aa-9 }}$ | sa-9' ${ }^{\text {d }}$ sa-9 |
| $\Delta G_{298 \mathrm{~K}}$ (B3LYP) | 0.0\|0.0 | 14.8 \| 3.5 | 10.6\|0.8 |
| $\Delta G_{298 \mathrm{~K}}$ (CBS-QB3) | $0.0 \mid 0.0$ | $15.5 \mid 4.5$ | 12.8 \|2.6 |
| Pop 298k | $99.2 \mid 66.1$ | $0.2 \mid 10.7$ | $\mathbf{0 . 6}$ \| 23.1 |

${ }^{\text {a }}$ Equilibrium populations were estimated from Boltzmann distribution based on the $\Delta G_{298 \mathrm{~K}}$ value computed at the CBS-QB3 level of theory. In the names of structures, a stands for anti and $\mathbf{s}$ stands for syn, which corresponds to the orientation of the azido (first letter) and hydroxy (second letter) moieties, relative to one another.

Table S2. Definition of internal coordinates used in the normal mode analysis of a- ${ }^{3} \mathbf{1 0}$. ${ }^{a}$

| Coordinate | Definition | Approximate description |
| :---: | :---: | :---: |
| $\mathrm{S}_{1}$ | $\mathrm{r}_{8,13}$ | $v(\mathrm{OH})$ |
| $\mathrm{S}_{2}$ | $\left(3^{-1 / 2}\right)\left(\mathrm{r}_{6,12}+\mathrm{r}_{5,11}+\mathrm{r}_{4,10}\right)$ | $v_{\mathrm{a}}(\mathrm{CH})$ |
| $\mathrm{S}_{3}$ | $\left(2^{-1 / 2}\right)\left(r_{6,12}-r_{4,10}\right)$ | $v_{b}(\mathrm{CH})$ |
| $\mathrm{S}_{4}$ | $\left(6^{-1 / 2}\right)\left(-r_{6,12}+2 r_{5,11}-r_{4,10}\right)$ | $v_{c}(\mathrm{CH})$ |
| $\mathrm{S}_{5}$ | $\left(12^{-1 / 2}\right)\left(2 r_{3,4}-r_{4,5}-r_{5,6}+2 r_{6,1}-r_{1,2}-r_{2,3}\right)$ | $v_{a}(\mathrm{CC})$ |
| $\mathrm{S}_{6}$ | $\left(4^{-1 / 2}\right)\left(\mathrm{r}_{4,5}-\mathrm{r}_{5,6}+\mathrm{r}_{1,2}-\mathrm{r}_{2,3}\right)$ | $v_{b}(\mathrm{CC})$ |
| $\mathrm{S}_{7}$ | $\left(4^{-1 / 2}\right)\left(r_{4,5}+r_{5,6}-r_{1,2}-r_{2,3}\right)$ | $v_{c}(\mathrm{CC})$ |
| $\mathrm{S}_{8}$ | $\left(12^{-1 / 2}\right)\left(2 r_{3,4}-r_{4,5}+r_{5,6}-2 r_{6,1}+r_{1,2}-r_{2,3}\right)$ | $v_{\text {d }}(\mathrm{CC})$ |
| S9 | $\left(6^{-1 / 2}\right)\left(r_{3,4}-r_{4,5}+r_{5,6}-r_{6,1}+r_{1,2}-r_{2,3}\right)$ | $v_{\mathrm{e}}(\mathrm{CC})$ |
| $\mathrm{S}_{10}$ | $\left(6^{-1 / 2}\right)\left(r_{3,4}+r_{4,5}+r_{5,6}+r_{6,1}+r_{1,2}+r_{2,3}\right)$ | $v_{f}(\mathrm{CC})$ |
| $\mathrm{S}_{11}$ | $\mathrm{r}_{1,7}$ | $v(\mathrm{C}-\mathrm{N})$ |
| $\mathrm{S}_{12}$ | $\mathrm{r}_{2,8}$ | $v(\mathrm{C}-\mathrm{O})$ |
| $\mathrm{S}_{13}$ | $\mathrm{r}_{3,9}$ | $v(\mathrm{C}-\mathrm{F})$ |
| $\mathrm{S}_{14}$ | $\left(6^{-1 / 2}\right)\left(\beta_{1,12,6}-\beta_{5,12,6}+\beta_{6,11,5}-\beta_{4,11,5}+\beta_{5,10,4}-\beta_{3,10,4}\right)$ | $\delta_{\mathrm{a}}(\mathrm{CH})$ |
| $\mathrm{S}_{15}$ | $\left(4^{-1 / 2}\right)\left(\beta_{1,12,6}-\beta_{5,12,6}-\beta_{5,10,4}+\beta_{3,10,4}\right)$ | $\delta_{\mathrm{b}}(\mathrm{CH})$ |
| $\mathrm{S}_{16}$ | $\left(12^{-1 / 2}\right)\left(-\beta_{1,12,6}+\beta_{5,12,6}+2 \beta_{6,11,5}-2 \beta_{4,11,5}-\beta_{5,10,4}+\beta_{3,10,4}\right)$ | $\delta_{\mathrm{c}}(\mathrm{CH})$ |
| $\mathrm{S}_{17}$ | $\left(2^{-1 / 2}\right)\left(\beta_{6,7,1}-\beta_{2,7,1}\right)$ | $\delta(\mathrm{C}-\mathrm{N})$ |
| $\mathrm{S}_{18}$ | $\left(2^{-1 / 2}\right)\left(\beta_{4,9,3}-\beta_{2,9,3}\right)$ | $\delta(\mathrm{C}-\mathrm{F})$ |
| $\mathrm{S}_{19}$ | $\left(2^{-1 / 2}\right)\left(\beta_{1,8,2}-\beta_{3,8,2}\right)$ | $\delta(\mathrm{C}-\mathrm{O})$ |
| $\mathrm{S}_{20}$ | $\beta_{2,13,8}$ | $\delta(\mathrm{OH})$ |
| $\mathrm{S}_{21}$ | $\left(6^{-1 / 2}\right)\left(\beta_{3,1,2}-\beta_{2,6,1}+\beta_{1,5,6}-\beta_{6,4,5}+\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\mathrm{a}}$ (ring) |
| $\mathrm{S}_{22}$ | $\left(12^{-1 / 2}\right)\left(2 \beta_{3,1,2}-\beta_{2,6,1}-\beta_{1,5,6}+2 \beta_{6,4,5}-\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\mathrm{b}}$ (ring) |
| $\mathrm{S}_{23}$ | $\left(4^{-1 / 2}\right)\left(\beta_{2,6,1}-\beta_{1,5,6}+\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\text {c }}$ (ring) |
| $\mathrm{S}_{24}$ | $\left(6^{-1 / 2}\right)\left(\tau_{3,2,1,6}-\tau_{2,1,6,5}+\tau_{1,6,5,4}-\tau_{6,5,4,3}+\tau_{5,4,3,2}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{a}}$ (ring) |
| $\mathrm{S}_{25}$ | $\left(12^{-1 / 2}\right)\left(-\tau_{32,1,6}+2 \tau_{2,1,6,5}-\tau_{1,6,5,4}-\tau_{6,5,4,3}+2 \tau_{5,4,3,2}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{b}}$ (ring) |
| $\mathrm{S}_{26}$ | $\left(4^{-1 / 2}\right)\left(\tau_{3,2,1,6}-\tau_{1,6,5,4}+\tau_{6,5,4,3}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{c}}$ (ring) |
| $\mathrm{S}_{27}$ | $\left(2^{-1 / 2}\right)\left(\tau_{13,8,2,1}-\tau_{13,8,2,3}\right)$ | $\tau(\mathrm{OH})$ |
| $\mathrm{S}_{28}$ | $\gamma_{7,2,1,6}$ | $\gamma(\mathrm{C}-\mathrm{N})$ |
| $\mathrm{S}_{29}$ | $\gamma_{8,3,2,1}$ | $\gamma$ (C-O) |
| $\mathrm{S}_{30}$ | $\gamma_{9,4,3,2}$ | $\gamma$ (C-F) |
| $\mathrm{S}_{31}$ | $\left(3^{-1 / 2}\right)\left(\gamma_{12,1,6,5}+\gamma_{11,6,5,4}+\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{a}}(\mathrm{CH})$ |
| $\mathrm{S}_{32}$ | $\left(6^{-1 / 2}\right)\left(-\gamma_{12,1,6,5}+2 \gamma_{11,6,5,4}-\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{b}}(\mathrm{CH})$ |
| $\mathrm{S}_{33}$ | $\left(2^{-1 / 2}\right)\left(\gamma_{12,1,6,5}-\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{c}}(\mathrm{CH})$ |

${ }^{a}$ Abbreviations: $v=$ stretching, $\delta=$ in-plane bending, $\gamma=$ out-of-plane bending, $\tau=$ torsion, ring $=6$-member ring (benzene). See Figure $S 5$ for the atom numbering scheme; $r_{i, j}$ is the distance between atoms $A_{i}$ and $A_{j} ; \beta_{i, j, k}$ is the angle between vectors $A_{k} A_{i}$ and $A_{k} A_{j}$; $\tau_{i, j, k, l}$ is the dihedral angle between the plane defined by $A_{i}, A_{j}, A_{k}$ and the plane defined by $A_{j}, A_{k}$ and $A_{l}$ atoms; $\gamma_{i, j, k, l}$ is the angle between the vector $A_{k} A_{i}$ and the plane defined by atoms $A_{j}, A_{k}, A_{l}$. The combinations $[(+),(+)]$ and $[(+),(-)]$ denote in-phase and in-opposite-phase couplings between coordinates of different types.

Table S3. Definition of internal coordinates used in the normal mode analysis of $\mathbf{Z} \mathbf{- 1 1}$. ${ }^{a}$

| Coordinate | Definition | Approximate description |
| :---: | :---: | :---: |
| $\mathrm{S}_{1}$ | $\mathrm{r}_{7,13}$ | $v$ (NH) |
| $\mathrm{S}_{2}$ | $\left(3^{-1 / 2}\right)\left(\mathrm{r}_{6,12}+\mathrm{r}_{5,11}+\mathrm{r}_{4,10}\right)$ | $\mathrm{v}_{\mathrm{a}}(\mathrm{CH})$ |
| $\mathrm{S}_{3}$ | $\left(2^{-1 / 2}\right)\left(r_{6,12}-r_{4,10}\right)$ | $v_{\mathrm{b}}(\mathrm{CH})$ |
| S 4 | $\left(6^{-1 / 2}\right)\left(-r_{6,12}+2 r_{5,11}-r_{4,10}\right)$ | $\nu_{c}(\mathrm{CH})$ |
| $\mathrm{S}_{5}$ | $\mathrm{r}_{2,8}$ | $v(\mathrm{C}=\mathrm{O})$ |
| $\mathrm{S}_{6}$ | $\mathrm{r}_{1,7}$ | $v(\mathrm{C}=\mathrm{N})$ |
| $\mathrm{S}_{7}$ | $\left(2^{-1 / 2}\right)\left(r_{4,3}-r_{6,5}\right)$ | $v(\mathrm{C}=\mathrm{C})_{\mathrm{as}}$ |
| $\mathrm{S}_{8}$ | $\left(2^{-1 / 2}\right)\left(r_{4,3}+r_{6,5}\right)$ | $v(\mathrm{C}=\mathrm{C})_{s}$ |
| S9 | $\left(2^{-1 / 2}\right)\left(r_{2,3}-r_{4,5}\right)$ | $\mathrm{v}_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\text {as }}$ |
| $\mathrm{S}_{10}$ | $\left(2^{-1 / 2}\right)\left(\mathrm{r}_{2,3}+\mathrm{r}_{4,5}\right)$ | $\mathrm{v}_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}$ |
| $\mathrm{S}_{11}$ | $\left(2^{-1 / 2}\right)\left(\mathrm{r}_{2,1}-\mathrm{r}_{1,6}\right)$ | $\mathrm{v}_{\mathrm{b}}(\mathrm{C}-\mathrm{C})_{\mathrm{as}}$ |
| $\mathrm{S}_{12}$ | $\left(2^{-1 / 2}\right)\left(r_{2,1}+r_{1,6}\right)$ | $v_{b}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}$ |
| $\mathrm{S}_{13}$ | $\mathrm{r}_{3,9}$ | $v(\mathrm{C}-\mathrm{F})$ |
| $\mathrm{S}_{14}$ | $\left(6^{-1 / 2}\right)\left(\beta_{1,12,6}-\beta_{5,12,6}+\beta_{6,11,5}-\beta_{4,11,5}+\beta_{5,10,4}-\beta_{3,10,4}\right)$ | $\delta_{\text {a }}(\mathrm{CH})$ |
| $\mathrm{S}_{15}$ | $\left(4^{-1 / 2}\right)\left(\beta_{1,12,6}-\beta_{5,12,6}-\beta_{5,10,4}+\beta_{3,10,4}\right)$ | $\delta_{\mathrm{b}}(\mathrm{CH})$ |
| $\mathrm{S}_{16}$ | $\left(12^{-1 / 2}\right)\left(-\beta_{1,12,6}+\beta_{5,12,6}+2 \beta_{6,11,5}-2 \beta_{4,11,5}-\beta_{5,10,4}+\beta_{3,10,4}\right)$ | $\delta_{\mathrm{c}}(\mathrm{CH})$ |
| $\mathrm{S}_{17}$ | $\left(2^{-1 / 2}\right)\left(\beta_{6,7,1}-\beta_{2,7,1}\right)$ | $\delta(\mathrm{C}=\mathrm{N})$ |
| $\mathrm{S}_{18}$ | $\left(2^{-1 / 2}\right)\left(\beta_{4,9,3}-\beta_{2,9,3}\right)$ | $\delta(\mathrm{C}-\mathrm{F})$ |
| $\mathrm{S}_{19}$ | $\left(2^{-1 / 2}\right)\left(\beta_{1,8,2}-\beta_{3,8,2}\right)$ | $\delta(\mathrm{C}=\mathrm{O})$ |
| $\mathrm{S}_{20}$ | $\beta_{1,13,7}$ | $\delta(\mathrm{NH})$ |
| $\mathrm{S}_{21}$ | $\left(6^{-1 / 2}\right)\left(\beta_{3,1,2}-\beta_{2,6,1}+\beta_{1,5,6}-\beta_{6,4,5}+\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\text {a }}$ (ring) |
| $\mathrm{S}_{22}$ | $\left(12^{-1 / 2}\right)\left(2 \beta_{3,1,2}-\beta_{2,6,1}-\beta_{1,5,6}+2 \beta_{6,4,5}-\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\mathrm{b}}$ (ring) |
| $\mathrm{S}_{23}$ | $\left(4^{-1 / 2}\right)\left(\beta_{2,6,1}-\beta_{1,5,6}+\beta_{5,3,4}-\beta_{4,2,3}\right)$ | $\delta_{\mathrm{c}}$ (ring) |
| $\mathrm{S}_{24}$ | $\left(6^{-1 / 2}\right)\left(\tau_{3,2,1,6}-\tau_{2,1,6,5}+\tau_{1,6,5,4}-\tau_{6,5,4,3}+\tau_{5,4,3,2}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{a}}$ (ring) |
| $\mathrm{S}_{25}$ | $\left(12^{-1 / 2}\right)\left(-\tau_{32,1,6}+2 \tau_{2,1,6,5}-\tau_{1,6,5,4}-\tau_{6,5,4,3}+2 \tau_{5,4,3,2}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{b}}$ (ring) |
| $\mathrm{S}_{26}$ | $\left(4^{-1 / 2}\right)\left(\tau_{3,2,1,6}-\tau_{1,6,5,4}+\tau_{6,5,4,3}-\tau_{4,3,2,1}\right)$ | $\tau_{\mathrm{c}}$ (ring) |
| $\mathrm{S}_{27}$ | $\left(2^{-1 / 2}\right)\left(\tau_{13,7,1,2}-\tau_{13,7,1,6}\right)$ | $\tau(\mathrm{NH})$ |
| $\mathrm{S}_{28}$ | $\gamma_{7,2,1,6}$ | $\gamma(\mathrm{C}=\mathrm{N})$ |
| $\mathrm{S}_{29}$ | $\gamma_{8,3,2,1}$ | $\gamma(\mathrm{C}=\mathrm{O})$ |
| $\mathrm{S}_{30}$ | $\gamma_{9,4,3,2}$ | $\gamma(\mathrm{C}-\mathrm{F})$ |
| $\mathrm{S}_{31}$ | $\left(3^{-1 / 2}\right)\left(\gamma_{12,1,6,5}+\gamma_{11,6,5,4}+\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{a}}(\mathrm{CH})$ |
| $\mathrm{S}_{32}$ | $\left(6^{-1 / 2}\right)\left(-\gamma_{12,1,6,5}+2 \gamma_{11,6,5,4}-\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{b}}(\mathrm{CH})$ |
| $\mathrm{S}_{33}$ | $\left(2^{-1 / 2}\right)\left(\gamma_{12,1,6,5}-\gamma_{10,5,4,3}\right)$ | $\gamma_{\mathrm{c}}(\mathrm{CH})$ |

[^0]Table S4. Experimental IR spectral data (nitrogen matrix at 10 K ), B3LYP/6-311+G(2d,p) computed vibrational frequencies ( $\tilde{v}, \mathrm{~cm}^{-1}$ ) and infrared intensities $\left(A^{\text {th }}, \mathrm{km} \mathrm{mol}^{-1}\right)$, and vibrational assignment of triplet anti-3-fluoro-2-hydroxyphenylnitrene $\mathbf{a -}{ }^{3} \mathbf{1 0}$.

| Experimental ${ }^{a}$ |  | Computed ${ }^{\text {b }}$ |  | PED ${ }^{c}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\tilde{v}$ | I | $\tilde{v}$ | $A^{\text {th }}$ |  |
| 3576 | br | 3596 | 130.2 | $100.0[\mathrm{v}(\mathrm{OH})$ ] |
| 1571 | S | 1564 | 55.8 | $35.2\left[\mathrm{v}_{\mathrm{b}}(\mathrm{CC})\right] ; 29.4\left[\mathrm{v}_{\mathrm{a}}(\mathrm{CC})\right]$ |
| 1538 | S | 1536 | 75.9 | $21.8\left[v_{\mathrm{b}}(\mathrm{CC})\right] ; 20.1\left[\mathrm{v}_{\mathrm{a}}(\mathrm{CC})\right]: 17.2\left[\delta_{\mathrm{b}}(\mathrm{CH})\right]$ |
| 1444 | vs | 1433 | 141.8 | $26.0\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] ; 17.6[v(\mathrm{C}-\mathrm{O})] ; 12.3\left[\mathrm{v}_{\mathrm{e}}(\mathrm{CC})\right] ; 10.1\left[\mathrm{v}_{\mathrm{a}}(\mathrm{CC})\right]$ |
|  |  | 1420 | 8.4 | $20.9\left[\mathrm{v}_{\mathrm{c}}(\mathrm{CC})\right] ; 20.9\left[\delta_{\mathrm{b}}(\mathrm{CH})\right] ; 17.4\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] \quad 14.6[8(\mathrm{OH})]$ |
| 1339 | m | 1344 | 37.2 | $40.5\left[\mathrm{v}_{\mathrm{e}}(\mathrm{CC})\right]+34.4[8(\mathrm{OH})]$ |
| 1303 | w | 1298 | 26.5 | $50.7[v(\mathrm{C}-\mathrm{N})] ; 16.7\left[\delta_{\mathrm{a}}(\right.$ ring $\left.)\right] ; 10.2[v(\mathrm{C}-\mathrm{O})]$ |
| 1249 | w | 1244 | 79.5 | $29.7\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] ; 12.6\left[\mathrm{v}_{\mathrm{d}}(\mathrm{CC})\right] ; 12.6\left[\mathrm{v}_{\mathrm{f}}(\mathrm{CC})\right] ; 12.2[v(\mathrm{C}-\mathrm{F})]$ |
| 1230 | m | 1220 | 64.7 | $53.6\left[\mathrm{v}_{\mathrm{e}}(\mathrm{CC})\right]-26.1[\delta(\mathrm{OH})]$ |
| 1191/1187 | w | 1175 | 24.3 | $18.3[v(\mathrm{C}-\mathrm{O})] ; 16.6\left[\mathrm{v}_{\mathrm{d}}(\mathrm{CC})\right] ; 16.3[v(\mathrm{C}-\mathrm{F})] ; 14.1\left[\mathrm{v}_{\mathrm{a}}(\mathrm{CC})\right] ; 11.4[v(\mathrm{C}-\mathrm{N})]$ |
| 1144 | w | 1139 | 20.8 | $59.9\left[\delta_{\mathrm{c}}(\mathrm{CH})\right] ; 18.9\left[v_{\mathrm{b}}(\mathrm{CC})\right]$ |
| 1056 | w | 1051 | 15.0 | $35.6\left[v_{\mathrm{c}}(\mathrm{CC})\right] ; 29.4\left[\delta_{\mathrm{b}}(\mathrm{CH})\right] ; 16.6\left[\mathrm{v}_{\mathrm{f}}(\mathrm{CC})\right]$ |
| 996/990 | S | 978 | 103.1 | $25.2[v(\mathrm{C}-\mathrm{F})] ; 24.0\left[v_{\mathrm{d}}(\mathrm{CC})\right]$ |
|  |  | 938 | 0.2 | $116.0\left[\gamma_{\mathrm{b}}(\mathrm{CH})\right]$ |
|  |  | 851 | 0.1 | 99.5[ $\left.\gamma_{\mathrm{c}}(\mathrm{CH})\right]$ |
|  |  | 823 | 7.8 | $49.8\left[\delta_{\mathrm{a}}(\mathrm{ring})\right] ; 15.7[v(\mathrm{C}-\mathrm{O})] ; 12.3\left[\delta_{\mathrm{b}}(\right.$ ring $\left.)\right]$ |
| 768/765 | m | 754 | 47.1 | $67.5\left[\gamma_{\mathrm{a}}(\mathrm{CH})\right] ; 13.9[\gamma(\mathrm{C}-\mathrm{N})] \quad 11.5[\gamma(\mathrm{C}-\mathrm{F})]$ |
| 696 | w | 686 | 17.6 | $52.6\left[\mathrm{v}_{\mathrm{f}}(\mathrm{CC})\right] ; 13.1[v(\mathrm{C}-\mathrm{F})]$ |
| 693 | w | 681 | 21.7 | $30.4\left[\tau_{\mathrm{a}}(\mathrm{ring})\right] ; 30.1\left[\gamma_{\mathrm{a}}(\mathrm{CH})\right] ; 29.6[\gamma(\mathrm{C}-\mathrm{O})] ; 13.2[\gamma(\mathrm{C}-\mathrm{N})]$ |

${ }^{\text {a }}$ Experimental intensities (I) are given in qualitative terms: vs = very strong $\mathrm{s}=$ strong; $\mathrm{m}=$ medium; $\mathrm{w}=$ weak; $\mathrm{br}=$ broad. The IR spectrum reported corresponds to the region between 1750 and $600 \mathrm{~cm}^{-1}$ and the $v(\mathrm{OH})$ mode. ${ }^{b}$ Computed harmonic wavenumbers $\left(\mathrm{cm}^{-1}\right)$ were multiplied by a 0.979 factor, except the $v(\mathrm{OH})$ mode which multiplied by a 0.950 factor. ${ }^{c}$ Abbreviations: $v=$ stretching, $\delta=$ in-plane bending, $\gamma=$ out-of-plane bending, $\tau=$ torsion, ring $=$ six-membered ring (benzene). PEDs are expressed in \%, and the PED matrices lower than $10 \%$ are not included. Definition of internal coordinates is given in Table S2.

Table S5. Experimental IR spectral data (nitrogen matrix at 10 K ), B3LYP/6-311+G(2d,p) computed vibrational frequencies $\left(\tilde{v}, \mathrm{~cm}^{-1}\right)$ and infrared intensities $\left(A^{t h}, \mathrm{~km} \mathrm{~mol}^{-1}\right)$, and vibrational assignment of ( $Z$ )-2-fluoro-6-iminocyclohexa-2,4-dienone $\mathbf{Z}$ - $\mathbf{1 1}$.

| Experimental ${ }^{\text {a }}$ |  | Calculated ${ }^{b}$ |  | $\mathrm{PED}^{c}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\tilde{v}$ | I | $\tilde{v}$ | $A^{\text {th }}$ |  |
| 3216 | br | 3218 | 14 | 99.9 [ $\mathrm{v}(\mathrm{NH})$ ] |
| 1691 | s | 1701 | 170.3 | 81.3[v(C=O)] |
| 1648 | S | 1659 | 53.3 | $62.5\left[v(\mathrm{C}=\mathrm{C})_{\mathrm{as}}\right] ; 10.2\left[\delta_{\mathrm{b}}(\mathrm{CH})\right]$ |
|  |  | 1604 | 38.9 | $75.5[v(\mathrm{C}=\mathrm{N})]$ |
| 1576 | w | 1576 | 8.9 | $69.3\left[v(\mathrm{C}=\mathrm{C})_{\mathrm{s}}\right]$ |
| 1417/1407 | w | 1410 | 5.9 | $21.0\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] ; 14.8\left[\mathrm{v}_{\mathrm{b}}(\mathrm{C}-\mathrm{C})_{\mathrm{as}}\right] ; 14.4\left[\delta_{\mathrm{b}}(\mathrm{CH})\right] ; 12.9\left[\delta_{\mathrm{c}}(\mathrm{CH})\right] ; 11.7[8(\mathrm{NH})]$ |
| 1358 | m | 1347 | 30.0 | $46.9[\delta(\mathrm{NH})] ; 20.8\left[\delta_{\mathrm{b}}(\mathrm{CH})\right]$ |
| 1347 | m | 1340 | 61.7 | $39.5\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] ; 13.2\left[\mathrm{v}_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{as}}\right] ; 11.9\left[v_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right]$ |
| 1256 | m | 1241 | 67.8 | $34.3[v(\mathrm{C}-\mathrm{F})] ; 16.3\left[\delta_{\mathrm{b}}(\mathrm{CH})\right] ; 15.2\left[\delta_{\mathrm{a}}(\mathrm{CH})\right] ; 13.7\left[\mathrm{v}_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right]$ |
| 1164 | m | 1163 | 40.9 | 57.2[ $\delta_{\mathrm{c}}(\mathrm{CH})$ ] |
| 1114 | s | 1098 | 104.9 | $37.7\left[v_{\mathrm{b}}(\mathrm{C}-\mathrm{C})_{\mathrm{as}}\right] ; 18.8[\delta(\mathrm{NH})] ; 15.5\left[\delta_{\mathrm{c}}(\mathrm{CH})\right]$ |
| 1018 | w | 1012 | 20.2 | $31.0\left[v_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{as}}\right] \quad 14.5\left[\delta_{\mathrm{b}}(\mathrm{CH})\right]$ |
| 993 | S | 979 | 35.5 | $21.7\left[v_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right] ; 19.5[\delta(\mathrm{C}=\mathrm{O})] ; 18.2\left[\mathrm{v}_{\mathrm{b}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right] \quad 13.1[v(\mathrm{C}-\mathrm{F})]$ |
|  |  | 978 | 2.1 | 92.1[ $\left.\gamma_{\mathrm{b}}(\mathrm{CH})\right] ; 14.0\left[\gamma_{\mathrm{c}}(\mathrm{CH})\right] ; 10.5[\tau(\mathrm{NH})]$ |
| 937 | w | 953 | 28.0 | 78.8[ $\tau(\mathrm{NH})] ; 16.7\left[\gamma_{\mathrm{b}}(\mathrm{CH})\right]$ |
|  |  | 911 | 4.9 | 81.2[ $\left.\gamma_{\mathrm{c}}(\mathrm{CH})\right] ; 12.2\left[\gamma_{\mathrm{a}}(\mathrm{CH})\right]$ |
| 836 | m | 828 | 34.4 | 64.2[ $\mathrm{C}_{\mathrm{a}}$ (ring] $]$ |
|  |  | 823 | 3.8 | 29.3[ $\left.\gamma_{\mathrm{a}}(\mathrm{CH})\right] ; 27.9[\gamma(\mathrm{C}=\mathrm{O})] ; 26.9[\gamma(\mathrm{C}=\mathrm{N})]$; 9.7[ $\left.\gamma(\mathrm{C}-\mathrm{F})\right]$ |
| 727 | m | 723 | 88.4 | $52.9\left[\gamma_{\mathrm{a}}(\mathrm{CH})\right] ; 19.8[\gamma(\mathrm{C}=\mathrm{O})] ; 17.2\left[\tau_{\mathrm{a}}(\right.$ ring $\left.)\right]$ |
|  |  | 670 | 3.9 | $29.9\left[v_{\mathrm{b}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right] ; 21.8\left[\mathrm{v}_{\mathrm{a}}(\mathrm{C}-\mathrm{C})_{\mathrm{s}}\right] ; 11.5\left[v(\mathrm{C}=\mathrm{C})_{\mathrm{s}}\right]$ |

${ }^{a}$ Experimental intensities (I) are given in qualitative terms: $\mathrm{s}=$ strong; $\mathrm{m}=$ medium; $\mathrm{w}=$ weak; $\mathrm{br}=$ broad. The IR spectrum reported corresponds to the region between 1750 and $600 \mathrm{~cm}^{-1}$ and the $v(\mathrm{NH})$ mode. ${ }^{b}$ Computed harmonic wavenumbers $\left(\mathrm{cm}^{-1}\right)$ were multiplied by a 0.979 factor, except for the $v(\mathrm{NH})$ mode which multiplied by a 0.950 factor. ${ }^{c}$ Abbreviations: $v=$ stretching, $\delta=$ in-plane bending, $\gamma=$ out-of-plane bending, $\tau=$ torsion, ring $=$ six-membered ring. PEDs are expressed in \%, and the PED matrices lower than $10 \%$ are not included. Definition of internal coordinates is given in Table S3.

## 4. Computational Section

Optimized geometries (Cartesian coordinates, $\AA$ ), electronic energies ( $E, \mathrm{E}_{\mathrm{h}}$ ) and zero-point vibrational energy (ZPVE, $\mathrm{E}_{\mathrm{h}}$ ) computed at the B3LYP/6-311+G(2d,p), CBS-QB3 and $\operatorname{CASSCF}(8,8) / 6-31 \mathrm{G}(\mathrm{d})$ levels of theory.

| as-9, |  |  |  |  |
| :--- | ---: | ---: | ---: | :---: |
| B3LYP | $(E=-471.205545 ;$ |  |  |  |
|  | ZPVE $=0.107459)$ |  |  |  |
| C | 0.876096 | 0.944931 | 0.000006 |  |
| C | -0.199146 | 0.044211 | 0.000032 |  |
| C | 0.035767 | -1.328335 | 0.000032 |  |
| C | 1.339474 | -1.808128 | 0.000002 |  |
| C | 2.406391 | -0.913722 | -0.000024 |  |
| C | 2.177104 | 0.456687 | -0.000022 |  |
| O | 0.672545 | 2.288987 | 0.000008 |  |
| H | -0.281055 | 2.456711 | 0.000023 |  |
| N | -1.483298 | 0.656112 | 0.000069 |  |
| N | -2.485112 | -0.058578 | -0.000016 |  |
| N | -3.467545 | -0.616059 | -0.000082 |  |
| H | -0.799461 | -2.019441 | 0.000055 |  |
| H | 1.518499 | -2.875673 | 0.000002 |  |
| H | 3.424187 | -1.283691 | -0.000046 |  |
| H | 2.995041 | 1.166002 | -0.000043 |  |

CBS-QB3 ( $E=-470.325318 ;$ ZPVE $=0.106642$ )
aa-9'
B3LYP $(E=-471.199342 ; \mathrm{ZPVE}=0.107022)$

| C | 0.899628 | 0.921892 | 0.000007 |
| :--- | ---: | ---: | ---: |
| C | -0.220214 | 0.075129 | -0.000010 |
| C | -0.033114 | -1.306463 | -0.000008 |
| C | 1.245460 | -1.850789 | -0.000002 |
| C | 2.352087 | -1.010718 | -0.000001 |
| C | 2.175582 | 0.368733 | 0.000014 |
| O | 0.676141 | 2.266451 | -0.000014 |
| H | 1.520616 | 2.731262 | -0.000004 |
| N | -1.487416 | 0.707021 | 0.000008 |
| N | -2.494134 | 0.000847 | 0.000008 |
| N | -3.491413 | -0.532165 | 0.000000 |
| H | -0.897547 | -1.960434 | -0.000012 |
| H | 1.371603 | -2.925905 | -0.000008 |
| H | 3.353492 | -1.422543 | 0.000004 |
| H | 3.036877 | 1.029381 | 0.000024 |

CBS-QB3 ( $E=-470.319212 ;$ ZPVE $=0.106292$ )

| -0.898903 | 0.926762 | 0.000000 |
| ---: | ---: | ---: |
| 0.220847 | 0.073494 | 0.000000 |
| 0.029305 | -1.309353 | 0.000000 |
| -1.252839 | -1.851842 | 0.000000 |
| -2.358471 | -1.007591 | 0.000000 |
| -2.177963 | 0.374007 | 0.000000 |
| -0.668384 | 2.267754 | 0.000000 |
| -1.512515 | 2.730650 | 0.000000 |
| 1.489050 | 0.704134 | 0.000000 |
| 2.495506 | -0.002960 | 0.000000 |
| 3.497656 | -0.534654 | 0.000000 |
| 0.893532 | -1.964603 | 0.000000 |
| -1.381270 | -2.927304 | 0.000000 |
| -3.361784 | -1.416412 | 0.000000 |
| -3.038234 | 1.037130 | 0.000000 |

## sa-9 ${ }^{\prime}$

B3LYP $(E=-471.199578 ; \mathrm{ZPVE}=0.106834)$

| C | 0.260703 | 0.820334 | -0.000003 |
| :--- | ---: | ---: | ---: |
| C | -0.027894 | -0.553835 | -0.000005 |
| C | 1.031836 | -1.461959 | -0.000002 |
| C | 2.349569 | -1.027861 | 0.000004 |
| C | 2.627538 | 0.334417 | 0.000007 |
| C | 1.582528 | 1.250801 | 0.000003 |
| O | -0.797272 | 1.688065 | -0.000009 |
| H | -0.476549 | 2.596750 | -0.000001 |
| N | -1.320112 | -1.118918 | -0.000012 |
| N | -2.352830 | -0.443208 | 0.000003 |
| N | -3.395304 | -0.010403 | 0.000015 |
| H | 0.792748 | -2.517861 | -0.000004 |
| H | 3.154271 | -1.751900 | 0.000006 |
| H | 3.650878 | 0.688059 | 0.000012 |
| H | 1.788865 | 2.316746 | 0.000004 |

## as-9

B3LYP $(E=-570.470642 ; \mathrm{ZPVE}=0.099307)$

| C | -0.651871 | -0.729403 | 0.000000 |
| :--- | ---: | ---: | ---: |
| C | 0.536322 | 0.014850 | 0.000006 |
| C | 0.507597 | 1.408121 | 0.000004 |
| C | -0.711976 | 2.071246 | -0.000003 |
| C | -1.902879 | 1.350965 | -0.000009 |
| C | -1.855172 | -0.030087 | -0.000007 |
| O | -0.664322 | -2.083115 | 0.000002 |
| H | 0.252699 | -2.394651 | 0.000008 |
| N | 1.712276 | -0.781975 | 0.000011 |
| N | 2.809740 | -0.222951 | 0.000008 |
| N | 3.863963 | 0.180351 | 0.000004 |
| H | 1.434340 | 1.969106 | 0.000008 |
| H | -0.736670 | 3.152991 | -0.000005 |
| H | -2.867140 | 1.841886 | -0.000015 |
| F | -3.000292 | -0.738503 | -0.000013 |

CBS-QB3 $(E=-470.319660 ;$ ZPVE $=0.106031)$

| -0.256950 | 0.821906 | 0.000001 |
| ---: | ---: | ---: |
| 0.027017 | -0.557375 | 0.000000 |
| -1.037015 | -1.463244 | -0.000001 |
| -2.355851 | -1.024137 | -0.000002 |
| -2.629743 | 0.340827 | -0.000001 |
| -1.580195 | 1.256134 | 0.000000 |
| 0.806273 | 1.679946 | 0.000002 |
| 0.489432 | 2.588616 | 0.000003 |
| 1.319040 | -1.122492 | 0.000000 |
| 2.350422 | -0.444597 | 0.000002 |
| 3.395414 | -0.008407 | -0.000002 |
| -0.797915 | -2.519557 | -0.000002 |
| -3.163208 | -1.746153 | -0.000003 |
| -3.652501 | 0.698037 | -0.000001 |
| -1.783701 | 2.323281 | 0.000001 |

CBS-QB3 $(E=-569.487174 ;$ ZPVE $=0.098506)$

| 0.650089 | -0.730728 | 0.000001 |
| ---: | ---: | ---: |
| -0.535992 | 0.021342 | -0.000002 |
| -0.503397 | 1.417256 | 0.000000 |
| 0.720691 | 2.075841 | 0.000005 |
| 1.909882 | 1.347518 | 0.000008 |
| 1.859961 | -0.035703 | 0.000006 |
| 0.653687 | -2.083220 | 0.000000 |
| -0.267411 | -2.379142 | -0.000004 |
| -1.712540 | -0.776172 | -0.000008 |
| -2.812256 | -0.222284 | -0.000010 |
| -3.872081 | 0.177125 | -0.000012 |
| -1.429628 | 1.980130 | -0.000003 |
| 0.751140 | 3.157983 | 0.000006 |
| 2.877590 | 1.832658 | 0.000012 |
| 3.001060 | -0.749968 | 0.000010 |

## aa-9

B3LYP $(E=-570.469178 ;$ ZPVE $=0.099174)$

| C | 0.660174 | -0.718145 | 0.000000 |
| :--- | ---: | ---: | ---: |
| C | -0.559333 | -0.027028 | 0.000002 |
| C | -0.561099 | 1.369645 | 0.000002 |
| C | 0.631294 | 2.081276 | 0.000000 |
| C | 1.850353 | 1.410504 | 0.000000 |
| C | 1.830196 | 0.031745 | 0.000001 |
| O | 0.674044 | -2.073894 | 0.000000 |
| H | 1.592456 | -2.374724 | 0.000002 |
| N | -1.722968 | -0.829382 | 0.000004 |
| N | -2.818985 | -0.268991 | -0.000001 |
| N | -3.880423 | 0.118154 | -0.000005 |
| H | -1.505123 | 1.900912 | 0.000004 |
| H | 0.611659 | 3.163129 | -0.000002 |
| H | 2.797567 | 1.933034 | -0.000001 |
| F | 2.995356 | -0.673072 | -0.000003 |

## sa-9

B3LYP $(E=-570.469343 ; \mathrm{ZPVE}=0.099036)$

| C | -0.173711 | -0.521867 | -0.000001 |
| :--- | ---: | ---: | ---: |
| C | 0.388485 | 0.762175 | -0.000006 |
| C | -0.459980 | 1.873008 | 0.000001 |
| C | -1.838240 | 1.722496 | 0.000003 |
| C | -2.408325 | 0.451832 | -0.000001 |
| C | -1.558671 | -0.632819 | 0.000000 |
| O | 0.632427 | -1.618747 | 0.000005 |
| H | 0.083471 | -2.413953 | -0.000005 |
| N | 1.769248 | 1.038169 | -0.000001 |
| N | 2.628585 | 0.149709 | 0.000000 |
| N | 3.548906 | -0.502162 | 0.000001 |
| H | -0.007358 | 2.855785 | 0.000003 |
| H | -2.474921 | 2.597373 | 0.000006 |
| H | -3.478795 | 0.296168 | 0.000000 |
| F | -2.056260 | -1.901595 | -0.000001 |

CBS-QB3 $(E=-569.529129 ;$ ZPVE $=0.098414)$

| 0.660198 | -0.719481 | 0.000000 |
| ---: | ---: | ---: |
| -0.560703 | -0.025767 | 0.000000 |
| -0.560731 | 1.373717 | 0.000000 |
| 0.633761 | 2.085822 | 0.000000 |
| 1.853932 | 1.411991 | 0.000000 |
| 1.833789 | 0.031027 | 0.000000 |
| 0.674182 | -2.074247 | 0.000000 |
| 1.594874 | -2.364212 | 0.000000 |
| -1.724509 | -0.828413 | 0.000001 |
| -2.821159 | -0.268921 | 0.000000 |
| -3.887840 | 0.114813 | -0.000001 |
| -1.505739 | 1.904421 | 0.000000 |
| 0.614975 | 3.168234 | 0.000000 |
| 2.803213 | 1.931632 | -0.000001 |
| 2.996923 | -0.679147 | 0.000000 |

CBS-QB3 $(E=-569.529865 ; \mathrm{ZPVE}=0.098204)$

| -0.170769 | -0.521289 | 0.000000 |
| ---: | ---: | ---: |
| 0.388123 | 0.767052 | 0.000001 |
| -0.465574 | 1.877318 | 0.000002 |
| -1.845452 | 1.722015 | 0.000001 |
| -2.411687 | 0.446886 | 0.000000 |
| -1.558393 | -0.637872 | 0.000000 |
| 0.639356 | -1.614449 | 0.000000 |
| 0.087005 | -2.405785 | -0.000003 |
| 1.768613 | 1.043642 | 0.000002 |
| 2.625597 | 0.152929 | -0.000001 |
| 3.547183 | -0.503476 | -0.000002 |
| -0.012770 | 2.860444 | 0.000003 |
| -2.485622 | 2.595041 | 0.000002 |
| -3.481743 | 0.285599 | 0.000000 |
| -2.047662 | -1.910671 | -0.000001 |

B3LYP $(E=-460.906470 ; \mathrm{ZPVE}=0.087242) \quad \operatorname{CASSCF}(8,8)(E=-458.294809 ; \mathrm{ZPVE}=0.092045)$
Orbital occupation: 29-1.95; 30-1.89; 31-1.89; 32-1.00; 33-1.00; 34-0.11; 35-0.12; 36-0.04

| C | -0.169030 | -0.742130 | 0.000000 | 0.000000 | 0.759581 | 0.000000 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| C | 1.238506 | -0.458095 | -0.000001 | 1.294995 | 0.170564 | 0.000000 |
| C | 1.644889 | 0.917825 | 0.000000 | 1.398225 | -1.252334 | 0.000000 |
| C | 0.716908 | 1.931897 | 0.000000 | 0.264378 | -2.046989 | 0.000000 |
| C | -0.656394 | 1.643440 | 0.000000 | -1.011878 | -1.459829 | 0.000000 |
| C | -1.055344 | 0.320885 | 0.000000 | -1.104372 | -0.074421 | 0.000000 |
| O | -0.576821 | -2.027133 | 0.000000 | -0.108218 | 2.100257 | 0.000000 |
| H | -1.543810 | -2.052492 | 0.000000 | -1.023261 | 2.350722 | 0.000000 |
| N | 2.131575 | -1.427827 | 0.000000 | 2.398981 | 0.928436 | 0.000000 |
| H | 2.706529 | 1.125413 | 0.000000 | 2.379580 | -1.687916 | 0.000000 |
| H | 1.042872 | 2.964029 | 0.000000 | 0.356416 | -3.116776 | 0.000000 |
| H | -1.405271 | 2.424316 | 0.000000 | -1.909574 | -2.047609 | 0.000000 |
| F | -2.380443 | 0.007518 | 0.000000 | -2.308707 | 0.513448 | 0.000000 |

s- ${ }^{3} 10$
B3LYP $(E=-460.907658 ; \mathrm{ZPVE}=0.087321) \quad \operatorname{CASSCF}(8,8)(E=-458.293127 ; \mathrm{ZPVE}=0.091971)$
Orbital occupation: 29-1.89; 30-1.95; 31-1.89; 32-1.00;
33-1.00; 34-0.11; 35-0.11; 36-0.04

| C | -0.055429 | -0.758352 | 0.000000 | -2.023770 | 1.248816 | 0.083925 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| C | 1.281564 | -0.222297 | 0.000000 | -2.258265 | 0.555554 | 1.306119 |
| C | 1.453497 | 1.201343 | 0.000000 | -2.728437 | -0.792268 | 1.267415 |
| C | 0.358494 | 2.028438 | 0.000000 | -2.952792 | -1.415352 | 0.055546 |
| C | -0.941994 | 1.496619 | 0.000000 | -2.719815 | -0.726584 | -1.149517 |
| C | -1.123387 | 0.125100 | 0.000000 | -2.262661 | 0.584014 | -1.110542 |
| O | -0.283584 | -2.085212 | 0.000000 | -1.581113 | 2.519317 | 0.045204 |
| H | 0.565656 | -2.551163 | 0.000000 | -1.456710 | 2.861398 | 0.920615 |
| N | 2.296283 | -1.058997 | 0.000000 | -2.035153 | 1.175649 | 2.468399 |
| H | 2.462380 | 1.591107 | 0.000000 | -2.899708 | -1.299414 | 2.198092 |
| H | 0.491751 | 3.102657 | 0.000000 | -3.306277 | -2.428899 | 0.027616 |
| H | -1.814574 | 2.136626 | 0.000000 | -2.886454 | -1.188326 | -2.103738 |
| F | -2.371888 | -0.378851 | 0.000000 | -2.044830 | 1.228435 | -2.250261 |

a- ${ }^{\text {OSS }} \mathbf{1 0}$
$\operatorname{CASSCF}(8,8)(E=-458.265910 ; \mathrm{ZPVE}=0.091230)$
Orbital occupation: 29-1.94; 30-1.91; 31-1.89; 32-1.00; 33-1.00; 34-0.10; 35-0.11; 36-0.05

| C | -0.225262 | 0.741431 | 0.000009 |
| :--- | ---: | ---: | ---: |
| C | 1.230674 | 0.527714 | -0.000001 |
| C | 1.700209 | -0.870447 | 0.000007 |
| C | 0.813642 | -1.908819 | -0.000003 |
| C | -0.594720 | -1.675665 | -0.000005 |
| C | -1.052297 | -0.343743 | 0.000003 |
| O | -0.658156 | 2.012529 | 0.000006 |
| H | -1.607289 | 2.030892 | 0.000000 |
| N | 2.053491 | 1.502238 | -0.000019 |
| H | 2.760651 | -1.032758 | 0.000011 |
| H | 1.172625 | -2.921137 | 0.000001 |
| H | -1.305158 | -2.478040 | -0.000007 |
| F | -2.373721 | -0.115297 | 0.000002 |

$\mathbf{s}^{-\mathrm{OSS}} \mathbf{1 0}$
$\operatorname{CASSCF}(8,8)(E=-458.264936 ; \mathrm{ZPVE}=0.090940)$
Orbital occupation: 29-1.89; 30-1.94; 31-1.91; 32-1.00; 33-1.00; 34-0.09; 35-0.11; 36-0.05

| C | -2.017179 | 1.268339 | 0.056366 |
| :--- | ---: | ---: | ---: |
| C | -2.252901 | 0.570573 | 1.332106 |
| C | -2.737267 | -0.817931 | 1.275932 |
| C | -2.954616 | -1.420909 | 0.071706 |
| C | -2.719695 | -0.726132 | -1.153011 |
| C | -2.255879 | 0.603649 | -1.113984 |
| O | -1.574821 | 2.537277 | 0.049019 |
| H | -1.458456 | 2.855870 | 0.934913 |
| N | -2.036743 | 1.171973 | 2.436872 |
| H | -2.907713 | -1.322696 | 2.207002 |
| H | -3.308171 | -2.434455 | 0.034914 |
| H | -2.887056 | -1.189956 | -2.104701 |
| F | -2.045490 | 1.226738 | -2.268262 |

$\mathbf{a}^{-{ }^{\mathrm{CSS}} 10}$
$\operatorname{CASSCF}(8,8)(E=-458.228879 ; \mathrm{ZPVE}=0.093165)$
Orbital occupation: 29-1.90; 30-1.96; 31-1.91; 32-1.62; 33-0.40; 34-0.03; 35-0.08; 36-0.09

| C | -0.130632 | 0.751156 | -0.000003 |
| :--- | ---: | ---: | ---: |
| C | 1.242914 | 0.409400 | 0.000015 |
| C | 1.599289 | -0.967194 | -0.000017 |
| C | 0.640251 | -1.964749 | -0.000006 |
| C | -0.719281 | -1.613022 | 0.000002 |
| C | -1.071128 | -0.280218 | -0.000004 |
| O | -0.512148 | 2.028759 | -0.000023 |
| H | -1.459474 | 2.085920 | -0.000032 |
| N | 2.197094 | 1.364490 | 0.000039 |
| H | 2.646880 | -1.201081 | -0.000036 |
| H | 0.922972 | -3.000303 | -0.000012 |
| H | -1.495140 | -2.355617 | 0.000012 |
| F | -2.362910 | 0.075260 | 0.000008 |

## TS-OH ${ }_{\text {rot }}$

B3LYP $(E=-460.897541 ; \mathrm{ZPVE}=0.085957)$

| C | 0.128349 | -0.753984 | 0.007462 |
| :--- | ---: | ---: | ---: |
| C | -1.255103 | -0.364306 | 0.002683 |
| C | -1.575211 | 1.032461 | -0.002343 |
| C | -0.574532 | 1.977152 | -0.003988 |
| C | 0.770717 | 1.592148 | -0.004227 |
| C | 1.089744 | 0.239395 | 0.002361 |
| O | 0.472770 | -2.071934 | 0.107190 |
| H | 0.611667 | -2.467635 | -0.762091 |
| N | -2.215902 | -1.269358 | -0.007977 |
| H | -2.619685 | 1.314171 | -0.003358 |
| H | -0.825803 | 3.030352 | -0.005880 |
| H | 1.571711 | 2.319718 | -0.007814 |
| F | 2.387498 | -0.119203 | -0.003802 |

## Z- ${ }^{\mathbf{3}} 11$

B3LYP $(E=-460.908591 ; \mathrm{ZPVE}=0.086885) \quad \operatorname{CASSCF}(8,8)(E=-458.290950 ; \mathrm{ZPVE}=0.091681)$
Orbital occupation: 29-1.94; 30-1.91; 31-1.03; 32-1.88;
33-0.97; 34-0.10; 35-0.12; 36-0.05

| C | -0.076310 | -0.839142 | -0.000006 | -0.028622 | -0.831968 | -0.055221 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| C | -1.304504 | 0.000546 | 0.000000 | 1.298462 | -0.159732 | -0.028015 |
| C | -1.161148 | 1.394963 | 0.000000 | 1.342238 | 1.237190 | 0.015609 |
| C | 0.118120 | 2.029329 | 0.000000 | 0.136402 | 2.040726 | 0.035119 |
| C | 1.269461 | 1.290108 | 0.000001 | -1.096202 | 1.446205 | 0.010729 |
| C | 1.166740 | -0.119670 | -0.000002 | -1.174506 | 0.027287 | -0.034101 |
| O | -0.128483 | -2.080594 | 0.000003 | -0.128065 | -2.057766 | -0.093291 |
| H | -2.406085 | -1.563292 | -0.000002 | 2.205023 | -1.848094 | -0.076559 |
| N | -2.516086 | -0.542967 | 0.000001 | 2.427940 | -0.865764 | -0.044745 |
| H | -2.063817 | 1.993542 | 0.000001 | 2.304710 | 1.710929 | 0.034383 |
| H | 0.170507 | 3.110609 | 0.000001 | 0.222920 | 3.109906 | 0.069181 |
| H | 2.251794 | 1.744115 | 0.000003 | -2.009102 | 2.010171 | 0.024246 |
| F | 2.290434 | -0.819584 | 0.000001 | -2.365191 | -0.532218 | -0.057329 |

## Z-11

B3LYP $(E=-460.953305 ; \mathrm{ZPVE}=0.089351) \quad \operatorname{CASSCF}(8,8)(E=-458.356915 ; \mathrm{ZPVE}=0.094747)$
Orbital occupation: 29-1.92; 30-1.94; 31-1.85; 32-1.9;
33-0.17; 34-0.05; 35-0.09; 36-0.08

| C | -0.046141 | -0.854440 | -0.000004 | -0.039481 | -0.843419 | -0.000146 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| C | 1.320381 | -0.152972 | -0.000002 | 1.311954 | -0.168488 | -0.000033 |
| C | 1.356654 | 1.305579 | -0.000001 | 1.351273 | 1.298539 | -0.000004 |
| C | 0.215125 | 2.013844 | 0.000000 | 0.212530 | 2.023581 | -0.000003 |
| C | -1.091474 | 1.374276 | 0.000002 | -1.108158 | 1.384223 | -0.000013 |
| C | -1.212893 | 0.038017 | 0.000000 | -1.214840 | 0.045911 | -0.000046 |
| O | -0.129200 | -2.064857 | -0.000001 | -0.140246 | -2.048775 | 0.000188 |
| H | 2.156313 | -1.839169 | 0.000008 | 2.185982 | -1.836474 | -0.000042 |
| N | 2.395264 | -0.839753 | 0.000004 | 2.395416 | -0.850218 | 0.000020 |
| H | 2.331017 | 1.776602 | -0.000004 | 2.321543 | 1.757095 | 0.000029 |
| H | 0.241249 | 3.096408 | -0.000003 | 0.248109 | 3.096776 | 0.000034 |
| H | -1.986353 | 1.985322 | 0.000005 | -1.996716 | 1.986834 | 0.000031 |
| F | -2.413932 | -0.551985 | 0.000001 | -2.391356 | -0.558713 | -0.000010 |

## TS-H shift

B3LYP $(E=-460.865278 ; \mathrm{ZPVE}=0.082345)$

| C | 0.154803 | -0.728945 | -0.016784 |
| :--- | ---: | ---: | ---: |
| C | 1.248903 | 0.211022 | 0.017665 |
| C | 1.021546 | 1.605581 | -0.018370 |
| C | -0.288619 | 2.028310 | -0.033104 |
| C | -1.372528 | 1.116837 | -0.009296 |
| C | -1.153840 | -0.244700 | 0.014701 |
| O | 0.582398 | -1.968614 | -0.089607 |
| H | 1.898004 | -1.594612 | -0.092381 |
| N | 2.360979 | -0.506004 | 0.095364 |
| H | 1.849175 | 2.301878 | -0.013018 |
| H | -0.509690 | 3.088116 | -0.037736 |
| H | -2.392698 | 1.478968 | 0.001647 |
| F | -2.188046 | -1.101337 | 0.051326 |

## MECP

B3LYP $(E=-460.888293 ; \mathrm{ZPVE}=0.0840)$
$\operatorname{CASSCF}(8,8)(E=-458.250949 ; \mathrm{ZPVE}=0.089486)$
Triplet Occupation: 29-1.89; 30-1.95; 31-1.91; 32-1.00;
33-1.00; 34-0.10; 35-0.10; 36-0.04
Singlet Occupation: 29-1.94; 30-1.97; 31-1.88; 32-1.91;

|  |  |  | $\mathbf{3 3 - 0 . 1 3 ; ~ 3 4 - 0 . 0 3 ; \mathbf { 3 5 } - 0 . 0 7 ; \mathbf { 3 6 } - 0 . 0 6}$ |  |  |  |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | :---: |
|  |  |  |  |  |  |  |  |
| C | 0.097463 | -0.758423 | 0.007280 | 0.114275 | -0.750916 | 0.008232 |  |
| C | 1.412020 | -0.053653 | 0.003939 | 1.430873 | -0.062799 | 0.005876 |  |
| C | 1.395259 | 1.396964 | 0.004944 | 1.401490 | 1.401969 | 0.005957 |  |
| C | 0.208741 | 2.059198 | 0.013092 | 0.216953 | 2.060141 | 0.013471 |  |
| C | -1.030607 | 1.346228 | 0.008278 | -1.051991 | 1.341828 | 0.009987 |  |
| C | -1.093131 | -0.019729 | 0.001573 | -1.107538 | -0.010947 | 0.004802 |  |
| O | 0.195628 | -2.041699 | 0.013844 | 0.193450 | -2.003310 | 0.011251 |  |
| H | 1.257062 | -2.097765 | 0.015605 | 1.236337 | -2.089038 | 0.011595 |  |
| N | 2.421414 | -0.842784 | 0.003710 | 2.412800 | -0.873229 | 0.005467 |  |
| H | 2.346971 | 1.912342 | 0.001600 | 2.341534 | 1.919145 | 0.000316 |  |
| H | 0.172542 | 3.140690 | 0.019116 | 0.176318 | 3.132976 | 0.017612 |  |
| H | -1.965439 | 1.895893 | 0.009212 | -1.974550 | 1.892237 | 0.006859 |  |
| F | -2.281917 | -0.650391 | 0.006190 | -2.253943 | -0.671188 | -0.005413 |  |

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[^0]:    ${ }^{a}$ Abbreviations: $v=$ stretching, $\delta=$ in-plane bending, $\gamma=$ out-of-plane bending, $\tau=$ torsion, ring $=$ six-membered ring. See Figure $S 5$ for the atom numbering scheme; $\mathrm{r}_{\mathrm{i}, \mathrm{j}}$ is the distance between atoms $\mathrm{A}_{\mathrm{i}}$ and $\mathrm{A}_{\mathrm{j}} ; \beta_{\mathrm{i}, \mathrm{j}, \mathrm{k}}$ is the angle between vectors $A_{k} A_{i}$ and $A_{k} A_{j} ; \tau_{i, j, k, l}$ is the dihedral angle between the plane defined by $A_{i}, A_{j}, A_{k}$ and the plane defined by $A_{j}$, $A_{k}$ and $A_{l}$ atoms; $\gamma_{i, j, k, l}$ is the angle between the vector $A_{k} A_{i}$ and the plane defined by atoms $A_{j}, A_{k}, A_{l}$. The combinations $[(+),(+)]$ and $[(+),(-)]$ denote in-phase and in-opposite-phase couplings between coordinates of different types.

